

11-2010


The Positive Solutions of the Matukuma Equation and the Problem of Finite Radius and Finite Mass

Jurgen Batt

Yi Li

Wright State University - Main Campus, yi.li@csun.edu

Follow this and additional works at: <https://corescholar.libraries.wright.edu/math>

 Part of the [Applied Mathematics Commons](#), [Applied Statistics Commons](#), and the [Mathematics Commons](#)

Repository Citation

Batt, J., & Li, Y. (2010). The Positive Solutions of the Matukuma Equation and the Problem of Finite Radius and Finite Mass. *Archive for Rational Mechanics and Analysis*, 198 (2), 613-675.
<https://corescholar.libraries.wright.edu/math/68>

This Article is brought to you for free and open access by the Mathematics and Statistics department at CORE Scholar. It has been accepted for inclusion in Mathematics and Statistics Faculty Publications by an authorized administrator of CORE Scholar. For more information, please contact corescholar@www.libraries.wright.edu, library-corescholar@wright.edu.

The Positive Solutions of the Matukuma Equation and the Problem of Finite Radius and Finite Mass

JÜRGEN BATT & YI LI

Abstract

This work is an extensive study of the 3 different types of positive solutions of the Matukuma equation $\frac{1}{r^2} (r^2 \phi')' = -\frac{r^{\lambda-2}}{(1+r^2)^{\lambda/2}} \phi^p, p > 1, \lambda > 0$: the E -solutions (regular at $r = 0$), the M -solutions (singular at $r = 0$) and the F -solutions (whose existence begins away from $r = 0$). An essential tool is a transformation of the equation into a 2-dimensional asymptotically autonomous system, whose limit sets (by a theorem of H.R. Thieme) are the limit sets of Emden-Fowler systems, and serve as a characterization of the different solutions. The emphasis lies on the study of the M -solutions. The asymptotic expansions obtained make it possible to apply the results to the important question of stellar dynamics which solutions lead to galactic models (stationary solutions of the Vlasov-Poisson system) of finite radius and/or finite mass for different p, λ .

Contents

1. Introduction	2
2. Preliminaries	6
2.1. The different types of solutions of the Matukuma equation	6
2.2. Transformation to Lotka–Volterra systems	8
2.3. Asymptotically autonomous systems	9
2.4. The linearization of $(EFS_{p,q})$ at the stationary points	10
2.5. The linearization of $(EFS_{p,1})$ at the stationary points	11
2.6. The flow of $(MS_{p,\lambda})$	11
3. F -solutions	13
4. E -solutions	14
5. M -solutions	19
5.1. The case $q > p$	19
5.1.1 The case $q > p + 1$	21
5.1.2 The case $q = p + 1$	24

5.1.3	The case $p < q < p + 1$	26
5.2.	The case $q = p$	32
5.3.	The case $q < p, q \neq \frac{p+1}{2}$	35
5.4.	The case $q = \frac{p+1}{2}$	37
6.	Solutions with $R < \infty$	43
7.	P_2 -solutions (solutions of fast decay)	45
8.	P_3^+ -solutions (solutions of slow decay)	51
9.	Global behaviour of the solutions and their radius and mass	52
9.1.	The case $q \leq \frac{p+1}{2}$	53
9.2.	M -solutions for $q > \frac{p+1}{2}$	54
9.3.	F -solutions	56
9.4.	Known result on E -solutions	59

Appendix

1. Introduction

In the present paper we study the positive solutions $\phi = \phi(r)$ of the differential equation

$$\frac{1}{r^2} (r^2 \phi')' = -\frac{r^{\lambda-2}}{(1+r^2)^{\lambda/2}} \phi^p, \quad p > 1, \lambda > 0, \quad (\text{M}_{p,\lambda})$$

on their maximal intervals (R_-, R) of existence on the positive real line \mathbb{R}^+ . For $\lambda = 2$, these solutions are the radial solutions of the semilinear elliptic equation

$$\Delta \phi = -\frac{1}{1+x^2} \phi^p,$$

on \mathbb{R}^3 , which T. Matukuma proposed in 1930 for the description of certain stellar globular clusters in a steady state [22]. In this paper, we attach the name "Matukuma equation" to the more general form $(\text{M}_{p,\lambda})$. Earlier, with the appearance of R. Emden's book "Gaskugeln" in 1907, certain stellar dynamic models called polytropes had been introduced; their study has lead to the Emden-Fowler equation

$$\frac{1}{r^2} (r^2 \phi')' = -r^{q-3} \phi^p, \quad p > 1, q > 1. \quad (\text{EF}_{p,q})$$

As was shown in [2], all these models are particular (classical) examples of a more general theory of time-independent spherically symmetric solutions of the Vlasov-Poisson system

$$\left. \begin{aligned} \partial_t f + v \partial_x f - \partial_x U(t, x) \partial_v f &= 0 \\ \Delta U(t, x) &= 4\pi \rho(t, x) \\ \rho(t, x) &:= \int f(t, x, v) dv, \end{aligned} \right\} \quad (\text{VPS})$$

$x, v \in \mathbb{R}^3$. Here $f = f(t, x, v) \geq 0$ is the distribution function of the considered system of gravitating mass in the space-velocity space $\mathbb{R}^3 \times \mathbb{R}^3$, where

$t \geq 0$ is the time, $U = U(t, x)$ the Newtonian potential and $\rho = \rho(t, x)$ the local density. Jeans' theorem says that the time-independent solutions f in the case of spherical symmetry have the form $f = \Psi(E, F)$, where

$$\begin{aligned} E &:= \frac{v^2}{2} + U(r) = \frac{w^2}{2} + \frac{1}{2} \frac{F}{r^2} + U(r), \\ F &:= x^2 v^2 - (xv)^2 \end{aligned}$$

are the local energy and the angular momentum in the coordinates

$$r := |x|, \quad w := \frac{x \cdot v}{|x|}, \quad F := x^2 v^2 - (xv)^2.$$

Given $\Psi \geq 0$, a solution U of Poisson's equation

$$\frac{1}{r^2} (r^2 U'(r))' = 4\pi \rho(r), \quad (1.1)$$

where

$$4\pi \rho(r) = \frac{8\pi^2}{r^2} \iint_{\substack{w > 0 \\ F > 0}} \Psi \left(\frac{w^2}{2} + \frac{1}{2} \frac{F}{r^2} + U(r), F \right) d(w, F) =: h_\Psi(r, U(r)), \quad (1.2)$$

together with $f := \Psi \left(\frac{w^2}{2} + \frac{1}{2} \frac{F}{r^2} + U(r), F \right)$ and ρ given by (1.2) leads to a solution of the (VPS) and hence to a spherically symmetric stationary stellar dynamic model (U, f, ρ) .

If we make the choice

$$\Psi(E, F) := (E_0 - E)_+^{n-3/2} F^k, \quad (\)_+ = \text{positive part},$$

where $E_0 > 0$, $n > 1/2$, $k > -1$, such that $n + k > 1$, (1.1) becomes

$$\frac{1}{r^2} (r^2 U'(r))' = cr^{2k} (E_0 - u)_+^{n+k}, \quad (1.3)$$

where $c := 2^{k+7/2} \pi^2 B\left(\frac{1}{2}, k+1\right) B\left(n - \frac{1}{2}, k + \frac{3}{2}\right)$ (B is the Beta function). The positive solutions $U < E_0$ of (1.3) correspond to the positive solutions ϕ of (EF_p, q) with $p := n + k$, $q := 2k + 3$ via the substitution $\phi(r) = E_0 - U\left(c^{-\frac{1}{2k+2}} r\right)$.

Camm's choice [6]

$$\Psi(E, F) := (E_0 - E - aF)_+^{n-3/2} F^k,$$

where $a > 0$, leads to (1.1) in the form

$$\frac{1}{r^2} (r^2 U'(r))' = \frac{cr^{2k}}{(1 + 2ar^2)^{k+1}} (E_0 - U(r))_+^{n+k}.$$

The positive solutions $U < E_0$ correspond to the positive solutions ϕ of $(\text{M}_{p,\lambda})$ with $\lambda := 2k + 2$ via the substitution

$$\phi(r) = \left(\frac{c}{(2a)^k + 1} \right)^{\frac{1}{n+k-1}} \left(E_0 - U \left(\frac{r}{\sqrt{2a}} \right) \right)$$

(see [2]).

Two physically most relevant problems are associated with the equations (1.1), (1.2), namely, to find those Ψ and those initial conditions for (1.1), for which

- A) the local density $\rho(r)$ has bounded support,
- B) the total mass $M := 4\pi \int \rho(r) r^2 dr$ is finite.

These problems had been completely solved for all positive solutions of the Emden-Fowler equation (EF $_{p,q}$) with $p, q > 1$ in [4] and investigated in the relativistic case in [1]; the results have been applied in the stability theory for the (VPS) in [3] and [10]. The problem in full generality seems to be extremely difficult, as pointed out in [2]. Further partial answers have been given by G. Rein and A.D. Rendall [27–29, 31, 32] for the (VPS) and the Vlasov-Einstein system, where this problem also occurs, and by M. Heinzle, A.D. Rendall and C. Uggle [11]. In [32] it is proven that the induced f has compact support and finite mass if Ψ is of the form

$$\Psi(E, F) = \Psi_1(E) \cdot F^l,$$

where $\Psi_1(E) = c \cdot (E_0 - E)_+^k + O\left((E_0 - E)_+^{k+\delta}\right)$ ($E \rightarrow E_0$), with $k > -1$, $l > -1/2$, $k + l + 1/2 > 0$, $k < l + 3/2$ and $c, \delta > 0$ (also the case of the Vlasov-Einstein equations is considered). The result is generalized in [11] to nonnegative functions $\Psi_1 \in C^1(0, \infty)$ such that Ψ_1' is locally bounded in $(0, \infty)$ and $\Psi_1'(E) \leq \text{const} \cdot E^{k'}$ for some $k' > -2$ near $E = 0$. The papers [4, 11] reduce the question to a phase space analysis for solutions of systems of ordinary differential equations.

As for the Matukuma equation and related equations, almost the entire work has been done for the E -solutions, beginning with the study of Matukuma's own conjectures about these solutions [24, p. 2]; among those, whose work is more closely related to the present paper, we mention W. Kawano, M. Kwong, Y. Li, W.-M. Ni, Y. Santanilla, E. Yanagida, S. Yotsutani [12–20, 23, 24, 35–38]. In Section 9.4, we refer to their work in more detail.

One aim of this paper is to develop a comprehensive theory of **all** positive solutions ϕ of the Matukuma equation. The three different types of solutions known from the Emden-Fowler equation also exist for the Matukuma equation, namely the E -solutions (which are regular at $r = 0$), the M -solutions (singular at $r = 0$) and the F -solutions (with $R_- > 0$). Particularly the M -solutions have an extremely rich structure. They correspond to stellar dynamic models with a singular, but integrable local density at $r = 0$. In this context we point out that the initial value problem for the (VPS) has

been satisfactorily solved for bounded (smooth) initial data, but for singular data a proof for global existence is still open: the singular stationary solutions, which we now know, are not contained in the class of initial conditions of any theorem yielding global existence of classical solutions [30]. To fill this gap is certainly a promising new line of research for the (VPS). Let us also mention the following result pertaining to the Vlasov-Einstein system [33]: If a spherically symmetric solution develops a singularity at all then the first singularity must appear at the center of the symmetry.

For the purpose of the latter applications, we derive asymptotic expansions at $r = 0$ and at $r = R = \infty$, even more accurate ones than those known for the Emden-Fowler equation [4, 5]. A first and powerful tool is the method which was invented for a precise asymptotic study of (not necessarily radial) solutions of certain semilinear equations at infinity [14, 16]. This method begins with rough estimates and improves the accuracy step by step to the desired extent. We show the applicability of this method at $r = 0$ for the E - and M -solutions. In the latter case, we found in Section 5.1 that for $q > p$ every M -solution ϕ admits a splitting of the form $\phi = S + \Theta$, where S is a singular and Θ is a regular part, which is an "E-solution" of a certain (singular) second order equation. For $p < q < p + 1$, S has the form of a "Laurent expansion" with respect to r^{q-p} with finitely many singular terms.

The aforementioned paper [4] was based on the substitution

$$u(t) = r^{q-2} \frac{\phi^p(r)}{-\phi'(r)}, \quad v(t) = r \frac{-\phi'(r)}{\phi(r)}, \quad r = e^t,$$

which maps the solutions of (EF_p, q) onto the solutions $\varphi = (u, v)$ of the Lotka-Volterra system

$$\begin{aligned} \dot{u} &= u(q - u - pv), \\ \dot{v} &= v(-1 + u + v). \end{aligned} \quad (\text{EFS}_p, q)$$

The F -, E - and M -solutions and those with $R < \infty$ could be characterized and analyzed by their limit sets $L^+(\varphi), L^-(\varphi)$ in the quadrant $\mathbb{R}^+ \times \mathbb{R}^+$, which can be found by means of the Poincaré-Bendixson theorem. If one applies this substitution to solutions of the $(\text{M}_{p,\lambda})$, one gets

$$\begin{aligned} \dot{u} &= u(q(t) - u - pv), \\ \dot{v} &= v(-1 + u + v). \end{aligned} \quad (\text{MS}_{p,\lambda})$$

The coefficient $q(t)$ is time-dependent, but the limits $\lim_{t \rightarrow -\infty} q(t) = 1 + \lambda, \lim_{t \rightarrow +\infty} q(t) = 1$ exist, and this fact makes this system an asymptotically autonomous one. We can make use of a theorem of H.R. Thieme [34], who has proved a structure theorem for the limit sets of asymptotically autonomous systems and we profit from the fact that the limiting systems in our case are just particular systems of Emden-Fowler type (EFS_p, q) , which had been completely studied in [4].

We have given an outlook of what the reader can expect in Sections 3, 4 and 5 which are devoted to the F -, E - and M -solutions, after the necessary preparations in Section 2. Sections 6, 7 and 8 are devoted to the solutions with $R < \infty$, and with rapid and slow decay at $R = \infty$. To derive the asymptotic expansions at $R = \infty$, we present another method, which makes use of the visualization of the solution by the phase space analysis, the representation of the $(MS_{p,\lambda})$ as a time-dependent perturbation of the (EF_p, q) and elements of the proof the Hartman-Grobman theorem. Section 9 addresses the questions A) and B) of finite R (i.e., the existence of a zero of ϕ) and finite M . The formula

$$M = \bar{c} \lim_{r \downarrow R_-} r^2 \phi'(r) - \bar{c} \lim_{r \uparrow R} r^2 \phi'(r), \quad (1.4)$$

where $\bar{c} := \left[\frac{(2a)^{\frac{q-p}{2}}}{c} \right]^{\frac{1}{p-1}}$ shows that B) is a question of the asymptotic expansions at $r = 0$ and $R = \infty$. Though there are solutions whose graph intersects itself, we have numerous sets of initial conditions for which the F -, E - and M -solutions have finite or infinite R and M (see also the appendix with graphs which illustrate our results). Finally, we remark that it follows from Section 2.2 that the (EFS_p, q) plays the key role of being the limit system for many other equations. But we refrain from extending our results to more general settings.

2. Preliminaries

2.1. The different types of solutions of the Matukuma equation

Let K be a positive function in $C^1(\mathbb{R}^+)$ with $r^2 K(r)$ bounded away from zero for $r \rightarrow \infty$, and $p > 1$. Let $\phi: (R_-, R) \rightarrow (0, \infty)$ be a maximal solution of the more general equation

$$\frac{1}{r^2} (r^2 \phi')' = -K(r) \phi^p \quad \text{on } \mathbb{R}^+ := (0, \infty). \quad (2.1.1)$$

Then $0 \leq R_- < R \leq \infty$. Let $r_0 \in (R_-, R)$ and

$$H(r) := \phi'(r_0) r_0^2 - \int_{r_0}^r s^2 K(s) \phi(s)^p ds \quad \text{on } (R_-, R). \quad (2.1.2)$$

Then $\phi'(r) = \frac{H(r)}{r^2}$ and $H'(r) = -r^2 K(r) \phi(r)^p < 0$ on (R_-, R) . Hence the limit

$$H_0 := \lim_{r \downarrow R_-} H(r) \in (-\infty, +\infty]$$

exists. If $H_0 > 0$, then $R_- > 0$. In fact, if $R_- = 0$, then there exists $r_1 \in (0, r_0)$ with $H(r_1) > 0$ and $H(r) > H(r_1)$ on $(0, r_1]$, hence

$$\begin{aligned} \phi(r) &= \phi(r_0) - \int_r^{r_0} \phi'(s) ds = \phi(r_0) - \int_r^{r_1} \frac{H(s)}{s^2} ds - \int_{r_1}^{r_0} \phi'(s) ds \\ &\leq \phi(r_0) - H(r_1) \int_r^{r_1} \frac{ds}{s^2} - \int_{r_1}^{r_0} \phi'(s) ds \longrightarrow -\infty \quad (r \rightarrow 0), \end{aligned}$$

which is a contradiction. Hence $R_- > 0$. In this case there exists an $R_0 \in (R_-, R)$ with $H(R_0) = 0 = \phi'(R_0)$: Otherwise we would have $H > 0$ on (R_-, R) , H and ϕ' are bounded and $\phi' > 0$. This implies $R = \infty$ and ϕ is increasing, so that $H(r) \rightarrow -\infty$ by the assumption on K . This is a contradiction, and $\phi' < 0$ on (R_0, R) and $H > 0$ and $\phi' > 0$ on (R_-, R_0) by letting $r_0 = R_0$ in (2.1.2). Then $\lim_{r \rightarrow R_-} \phi'(r)$ is finite and $\lim_{r \rightarrow R_-} \phi(r) = 0$. If $H_0 \leq 0$, then $R_- = 0$. Assume $R_- > 0$. Then since $H' < 0$ on (R_-, R_0) there exists $r_2 \in (R_-, r_1)$ so that $H(r_2) < 0$, $\phi'(r_2) < 0$, H and ϕ' are bounded on (R_-, r_2) and ϕ could be extended beyond R_- . Hence $R_- = 0$, $\phi' < 0$ on $(0, R)$ and the limit

$$\lim_{r \downarrow 0} \phi(r) \in (0, \infty]$$

exists. In this case, we define $R_0 := 0$ and have

$$R_0 = \inf \{ r \in (R_-, R) : \phi'(r) < 0 \}$$

for all solutions.

Now we classify the solutions as follows. We have shown: $H_0 > 0 \iff R_- > 0 \iff R_0 > R_-$. In this case we call ϕ an F -solution. Furthermore: $H_0 \leq 0 \iff R_- = 0 \iff R_0 = 0$. In this case, if $\lim_{r \downarrow 0} \phi(r) < \infty$, we call ϕ an E -solution, if $\lim_{r \downarrow 0} \phi(r) = \infty$, an M -solution.

Particular examples of (2.1.1) are the Emden-Fowler equation

$$\frac{1}{r^2} (r^2 \phi')' = -r^{q-3} \phi^p, \quad p > 1, q > 1, \quad (\text{EF}_{p,q})$$

and the Matukuma equation

$$\frac{1}{r^2} (r^2 \phi')' = -\frac{r^{\lambda-2}}{(1+r^2)^{\lambda/2}} \phi^p, \quad p > 1, \lambda > 0. \quad (\text{M}_p, \lambda)$$

We will show that the three types of solutions also exist for this equation. As for the E -solutions, we will frequently make use of the following lemma which is part of the more general Theorem 5.6 in [2, p. 172], see also [24].

Lemma 2.1. *Let $\alpha \in \mathbb{R}$ and let $h : (0, \infty) \times \mathbb{R} \rightarrow \mathbb{R}$ satisfy the following conditions:*

- (i) $h \in C^1((0, \infty) \times \mathbb{R})$,
- (ii) $sh(\cdot, \alpha) \in L^1_{loc}[0, \infty)$,
- (iii) there exists a number $\delta > 0$ and a function

$$L_{0\alpha} : (0, \delta) \rightarrow [0, \infty] \text{ with } rL_{0\alpha} \in L^1[0, \delta],$$

such that for all $r \in (0, \delta)$ and $u_1, u_2 \in [\alpha - \delta, \alpha + \delta]$ we have

$$|h(r, u_1) - h(r, u_2)| \leq L_{0\alpha}(r)|u_1 - u_2|.$$

Then the initial value problem

$$\frac{1}{r^2} (r^2 \phi')' = h(r, \phi), \quad \phi(0) = \alpha$$

has a unique (maximal) solution ϕ on $(0, R)$ with $\phi(0) := \lim_{r \rightarrow 0} \phi(r) = \alpha$.

If, for some $\gamma \geq 0$, we have

$$\frac{1}{r^{\gamma+1}} \int_0^r s^2 L_{0\alpha}(s) ds = O(1) \quad (r \rightarrow 0),$$

then

$$r^{1-\gamma} \phi'(r) - \frac{1}{r^{\gamma+1}} \int_0^r s^2 h(s, \alpha) ds \rightarrow 0 \quad (r \rightarrow 0).$$

2.2. Transformation to Lotka–Volterra systems

In this section we consider solutions ϕ of (2.1.1) on their intervals $J_\phi := (R_0, R)$. We define

$$u(t) := rK(r) \frac{\phi^p(r)}{-\phi'(r)}, \quad v(t) := r \frac{-\phi'(r)}{\phi(r)}, \quad r := e^t.$$

Then $\varphi := (u, v) : I_\varphi \rightarrow \mathbb{R}^+ \times \mathbb{R}^+$, $I_\varphi := \ln J_\phi$ is a maximal solution of the system

$$\begin{aligned} \dot{u} &= u(q(t) - u - pv), \\ \dot{v} &= v(-1 + u + v), \quad \text{with } q(t) := 3 + \frac{rK'(r)}{K(r)}. \end{aligned} \quad (2.2.1)$$

$\mathbb{R}^+ \times \mathbb{R}^+$ is an invariant set of this system (the positive u - and v -axes are invariant). The inverse is

$$\phi(r) = \left(\frac{u(\ln r) v(\ln r)}{r^2 K(r)} \right)^{\frac{1}{p-1}}. \quad (2.2.2)$$

For (EF_p, q) we have $K(r) = r^{q-3}$ and thus obtain

$$\begin{aligned} \dot{u} &= u(q - u - pv), \\ \dot{v} &= v(-1 + u + v). \end{aligned} \quad (\text{EFS}_p, q)$$

For (M_p, λ) we have

$$K(r) = \frac{r^{\lambda-2}}{(1+r^2)^{\lambda/2}}, \quad (2.2.3)$$

and

$$u(t) = \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\phi^p(r)}{-\phi'(r)}, \quad v(t) = r \frac{-\phi'(r)}{\phi(r)} \quad (2.2.4)$$

satisfy

$$\begin{aligned} \dot{u} &= u(q(t) - u - pv), \\ \dot{v} &= v(-1 + u + v), \end{aligned} \quad (\text{MS}_p, \lambda)$$

with

$$q(t) = 3 + \frac{\lambda - 2 - 2r^2}{1 + r^2} = 1 + \frac{\lambda}{1 + r^2}.$$

The inverse is

$$\begin{aligned}\phi(r) &= \left((1+r^2)^{\lambda/2} r^{-\lambda} u(\ln r) v(\ln r) \right)^{\frac{1}{p-1}}, \\ \phi'(r) &= -\frac{1}{r} \phi(r) \cdot v(\ln r),\end{aligned}\quad (2.2.5)$$

and the limits

$$\lim_{t \rightarrow -\infty} q(t) = 1 + \lambda, \quad \lim_{t \rightarrow +\infty} q(t) = 1$$

exist. We define $q := 1 + \lambda$. Then $(MS_{p,\lambda})$ is asymptotically autonomous with respect to $(EFS_{p,q})$ for $t \rightarrow -\infty$ and to $(EFS_{p,1})$ for $t \rightarrow +\infty$ in the sense of the next section. The following decompositions:

$$q(t) = q - (q-1) \frac{e^{2t}}{1+e^{2t}} = q - (q-1) e^{2t} + (q-1) \frac{e^{4t}}{1+e^{2t}} \quad (2.2.6)$$

$$= 1 + (q-1) \frac{e^{-2t}}{1+e^{-2t}} = 1 + (q-1) e^{-2t} - (q-1) \frac{e^{-4t}}{1+e^{-2t}}, \quad (2.2.7)$$

are useful for $t \rightarrow -\infty$ and $t \rightarrow +\infty$ respectively.

In the sequel unless otherwise stated, ϕ denotes a solution of $(M_{p,\lambda})$ on (R_0, R) with $0 \leq R_0 < R \leq \infty$ and $\varphi = (u, v)$ is the associated solution of $(MS_{p,\lambda})$ on (T_0, T) with $-\infty \leq T_0 = \ln R_0 < T = \ln R \leq \infty$. And it is always assumed that $p, q > 1$, $\lambda = q - 1 > 0$.

2.3. Asymptotically autonomous systems

A system of ordinary differential equations in the plane \mathbb{R}^2

$$\dot{x} = f(t, x) \quad (\text{NA})$$

is called asymptotically autonomous with respect to the autonomous system

$$\dot{x} = g(x) \quad (\text{A})$$

for $t \rightarrow +\infty$, if $f(t, x) \rightarrow g(x)$ as $t \rightarrow +\infty$ uniformly on compact subsets of \mathbb{R}^2 . As usual we use $L^+(\varphi)$, $L^-(\varphi)$ for the positive, negative limit set and $C^+(\varphi)$, $C^-(\varphi)$ for the positive, negative half trajectory of an orbit φ .

For asymptotically autonomous systems H.R. Thieme [34] has proved the following theorem.

Theorem 2.2. (H.R. Thieme) *Let φ on (τ, ∞) be a solution of (NA) and assume $L^+(\varphi) \subset K$ for a compact subset K of \mathbb{R}^2 . If K only contains finitely many stationary points of (A), then exactly one of the following (distinct) 3 cases occurs:*

- (A) $L^+(\varphi)$ consists of a stationary point of (A),
- (B) $L^+(\varphi)$ is the union of periodic orbits of (A) and possibly center points of (A), which are encircled by periodic orbits of (A) lying in $L^+(\varphi)$,

© $L^+(\varphi)$ contains stationary points of (A) which are cyclically chained by orbits of (A) lying in $L^+(\varphi)$.

A similar statement holds for $t \rightarrow -\infty$.

As is clear from the last paragraph, we can apply Thieme's theorem to an analysis of $(MS_{p,\lambda})$ with its autonomous limit systems $(EFS_{p,q})$ and $(EFS_{p,1})$ being in the good position that the limit sets of $(EFS_{p,q})$ are completely known from [4].

2.4. The linearization of $(EFS_{p,q})$ at the stationary points

The system $(EFS_{p,q})$, $p, q > 1$, has the following stationary points: $P_1 = (0, 0)$, $P_2 = (0, 1)$, $P_3^- = (q, 0)$; for $q < p$: $P_4 = \left(\frac{p-q}{p-1}, \frac{q-1}{p-1}\right)$. If the stationary point is $P = (a, b)$, and if we introduce the coordinates $\bar{u} := u - a, \bar{v} := v - b$, then we can write $(MS_{p,\lambda})$ as a time-dependent perturbation of $(EFS_{p,q})$:

$$\begin{pmatrix} \dot{\bar{u}} \\ \dot{\bar{v}} \end{pmatrix} = A \begin{pmatrix} \bar{u} \\ \bar{v} \end{pmatrix} + \begin{pmatrix} -\bar{u}^2 - p\bar{u}\bar{v} \\ \bar{u}\bar{v} + \bar{v}^2 \end{pmatrix} + \begin{pmatrix} -(q-1)(\bar{u}+a) \frac{e^{2t}}{1+e^{2t}} \\ 0 \end{pmatrix}, \quad (2.4.1)$$

with

$$A := \begin{pmatrix} q - 2a - pb & -pa \\ b & a + 2b - 1 \end{pmatrix}. \quad (2.4.2)$$

For the following, we consider the linear part of (2.4.1).

For P_1 we have $A = \begin{pmatrix} q & 0 \\ 0 & -1 \end{pmatrix}$ with eigenvalues $\lambda_1 = q$, $\lambda_2 = -1$ and eigenvectors $\xi_1 = (0, 1)$, $\xi_2 = (1, 0)$, and P_1 is a saddle.

For P_2 we have $A = \begin{pmatrix} q - p & 0 \\ 1 & 1 \end{pmatrix}$ with $\lambda_1 = q - p$, $\lambda_2 = 1$ and $\xi_1 = (q - p - 1, 1)$, $\xi_2 = (0, 1)$. That is, for $q > p$ P_2 is an unstable improper node. In particular, for $q > p + 1$: $0 < \lambda_2 < \lambda_1$ and P_2 is a 2-tangential improper node with main direction ξ_2 ; for $q = p + 1$: $0 < \lambda_2 = \lambda_1 = 1$ and P_2 is a 1-tangential node with direction $(0, 1)$; for $p < q < p + 1$: $0 < \lambda_1 < \lambda_2$ and P_2 is a 2-tangential node with main direction $-\xi_1$; for $q = p$: $\lambda_1 = 0$ and P_2 is an unstable 2-tangential node with main direction $(1, -1)$; for $q < p$: P_2 is a saddle.

For P_3^- we have $A = \begin{pmatrix} -q & -pq \\ 0 & q - 1 \end{pmatrix}$ with $\lambda_1 = -q$, $\lambda_2 = q - 1$ and $\xi_1 = (1, 0)$, $\xi_2 = \left(-1, \left(2 - \frac{1}{q}\right) \frac{1}{p}\right)$, and P_3^- is a saddle.

For

$$q < p : P_4 = \left(\frac{p-q}{p-1}, \frac{q-1}{p-1}\right) =: (u_4, v_4) \quad (u_4 + v_4 = 1),$$

$$A = \begin{pmatrix} v_4 - 1 & p(v_4 - 1) \\ v_4 & v_4 \end{pmatrix}, \quad \lambda_{1,2} = v_4 - \frac{1}{2} \pm \frac{1}{2} \sqrt{\Theta(v_4)},$$

where $\Theta(s) = 1 - 4ps(1 - s)$ with two distinct roots $\eta_{1,2} = \frac{1}{2} \left(1 \pm \sqrt{1 - \frac{1}{p}}\right)$. For $v_4 > \frac{1}{2} \Leftrightarrow q > \frac{p+1}{2}$ we have $\operatorname{Re} \lambda_i > 0, i = 1, 2$ and P_4 is unstable. In particular, for $\eta_1 \leq v_4 < 1$, P_4 is an improper node and for $\frac{1}{2} < v_4 < \eta_1$ a spiral point. For $v_4 = \frac{1}{2} \Leftrightarrow q = \frac{p+1}{2}$ we have $\lambda_{1,2} = \pm \frac{1}{2} \sqrt{p-1}i$ and P_4 is a center. For $0 < v_4 < \frac{1}{2} \Leftrightarrow q < \frac{p+1}{2}$ we have $\operatorname{Re} \lambda_i < 0, i = 1, 2$ and P_4 is stable. In particular, for $\eta_2 < v_4 < \frac{1}{2}$ P_4 is a spiral point and for $0 < v_4 \leq \eta_2$ P_4 is an improper node.

2.5. The linearization of $(\text{EFS}_{p,1})$ at the stationary points

The system $(\text{EFS}_{p,1})$, $p > 1$, has the following stationary points: $P_1 = (0, 0)$, $P_2 = (0, 1)$, $P_3^+ = (1, 0)$. With the notation of Section 2.4 (here we have $q = 1$)

$$\begin{pmatrix} \dot{\bar{u}} \\ \dot{\bar{v}} \end{pmatrix} = A \begin{pmatrix} \bar{u} \\ \bar{v} \end{pmatrix} + \begin{pmatrix} -\bar{u}^2 - p\bar{u}\bar{v} \\ \bar{u}\bar{v} + \bar{v}^2 \end{pmatrix}, \text{ where } A := \begin{pmatrix} -1 - 2a - pb & -pa \\ b & -1 + a + 2b \end{pmatrix}.$$

For P_1 we have $A = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ and P_1 is a saddle.

For P_2 we have $A = \begin{pmatrix} 1 - p & 0 \\ 1 & 1 \end{pmatrix}$ with eigenvalues $\lambda_1 = 1 - p < 0$, $\lambda_2 = 1$ and eigenvectors $\xi_1 = (p, -1)$, $\xi_2 = (0, 1)$ and P_2 is a saddle.

For P_3^+ we have $A = \begin{pmatrix} -1 & -p \\ 0 & 0 \end{pmatrix}$ with $\lambda_1 = -1$, $\lambda_2 = 0$ and $\xi_1 = (1, 0)$, $\xi_2 = (-p, 1)$ and P_3^+ is a stable 2-tangential improper node with main direction ξ_2 .

2.6. The flow of $(\text{MS}_{p,\lambda})$

From (2.2.6), (2.2.7) we have $\dot{q}(t) < 0$, so that $q(\cdot)$ strictly decreases from $q(-\infty) = q$ to $q(+\infty) = 1$ on $\bar{\mathbb{R}} := \{-\infty \leq t \leq +\infty\}$. We define

$$\begin{aligned} S(t, u, v) &:= q(t) - u - pv, & t \in \bar{\mathbb{R}}, \\ W(u, v) &:= -1 + u + v, \end{aligned}$$

and have $(\text{MS}_{p,\lambda})$ in the form

$$\begin{aligned} \dot{u} &= uS(t, u, v), \\ \dot{v} &= vW(u, v). \end{aligned}$$

For functions $F: \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}$, we let

$$F_0 := F^{-1}\{0\}, \quad F_+ := F^{-1}\{\mathbb{R}^+\}, \quad F_- := F^{-1}\{\mathbb{R}^-\}.$$

$S_0(t)$, $t \in \bar{\mathbb{R}}$, is the straight line $v = -\frac{1}{p}u + \frac{q(t)}{p}$ with intercepts $(q(t), 0)$ and $(0, \frac{q(t)}{p})$. It moves downward in the closed strip Z between $S_0(-\infty)$

and $S_0(+\infty)$ when t runs from $-\infty$ to $+\infty$. W_0 is the fixed straight line $v = 1 - u$. Furthermore, we define

$$G(u, v) := u + \frac{p+1}{2}(v-1).$$

G_0 is the straight line $v = -\frac{2}{p+1}u + 1$ with intercepts $(\frac{p+1}{2}, 0)$ and $(0, 1)$. If $q > p$, $P_2 \in Z$. If $q > \frac{p+1}{2}$, we define $T^* \in \mathbb{R}$ by $q(T^*) := \frac{p+1}{2}$ and $T^* =: -\infty$ for $q \leq \frac{p+1}{2}$. Then $S_0(t) \subset G_-$ for $t \geq T^*$.

Lemma 2.3. *Let φ be a solution of $(MS_{p,\lambda})$ on (T_0, T) . Let $q > \frac{p+1}{2}$.*

If $\varphi(t_0) \in G_+$ for $t_0 < T^$, then $\varphi(t) \in G_+$ for $t_0 \leq t \leq T^*$.*

If $\varphi(t_0) \in G_- \cup G_0$ for $t_0 \leq T^$, then $\varphi(t) \in G_-$ for $t < t_0$.*

If $\varphi(t_0) \in G_- \cup G_0$ for $t_0 \geq T^$, then $\varphi(t) \in G_-$ for $t > t_0$.*

If $\varphi(t_0) \in G_+ \cup G_0$ for $t_0 > T^$, then $\varphi(t) \in G_+$ for $T^* \leq t < t_0$.*

For $q \leq \frac{p+1}{2}$ the last two statements are true for all t_0 .

Proof. We have

$$\begin{aligned} \frac{d}{dt}G(\varphi(t)) &= \dot{u} + \frac{p+1}{2}\dot{v} = u(q(t) - u - pv) + \frac{p+1}{2}v(-1 + u + v) \\ &= u\left(q(t) - u + \frac{-p+1}{2}v\right) + \frac{p+1}{2}v(-1 + v). \end{aligned}$$

If $\varphi(t) \in G_0$, then $-u = \frac{p+1}{2}(v-1)$, hence

$$\begin{aligned} \left. \frac{d}{dt}G(\varphi(t)) \right|_{\varphi(t) \in G_0} &= u\left(q(t) + \frac{p+1}{2}(v-1) + \frac{-(p+1)}{2}v - v\right) \\ &= u\left(q(t) - \frac{p+1}{2}\right) \\ &= \begin{cases} u(q(t) - q(T^*)) > 0 & \text{for } q > \frac{p+1}{2} \text{ and } t < T^*, \\ u(q(t) - q(T^*)) < 0 & \text{for } q > \frac{p+1}{2} \text{ and } t > T^*, \\ < 0 & \text{for } q \leq \frac{p+1}{2} \text{ and all } t. \end{cases} \end{aligned}$$

This proves the assertions for $t_0 \neq T^*$. If $q > \frac{p+1}{2}$, then from the above and an easy calculation we have

$$\left. \begin{aligned} \frac{d}{dt}G(\varphi(t)) \Big|_{\{\varphi(T^*) \in G_0\}} &= 0, \\ \frac{d^2}{dt^2}G(\varphi(t)) \Big|_{\{\varphi(T^*) \in G_0\}} &= \frac{(q - \frac{p+1}{2})(p-1)(p+1)}{q-1}(v(T^*) - 1) < 0, \end{aligned} \right\} \quad (2.6.1)$$

which shows the claim also for $t_0 = T^* > -\infty$.

Remark 1. Essentially, the foregoing proof consists of the calculation

$$uS(t, u, v) + \frac{p+1}{2}vW(u, v) \Big|_{(u,v) \in G_0} = u\left(q(t) - \frac{p+1}{2}\right).$$

The left-hand side is equal to $\langle n, (\dot{u}, \dot{v}) \rangle|_{(u,v) \in G_0}$, where $n := (1, \frac{p+1}{2})$ is a vector normal to G_0 . The geometric interpretation of the result is then clear: the flow is directed outward of G_- for $q > \frac{p+1}{2}, t < T^*$, and inward for $t > T^*$, and for $q \leq \frac{p+1}{2}$ for all t .

Lemma 2.4. *No solution φ of $(MS_{p,\lambda})$ can converge to P_1 for $t \rightarrow -\infty$ nor for $t \rightarrow +\infty$.*

Proof. We have $P_1 \subset S_+(t) \cap W_-$, where $\dot{u}(t) > 0$ and $\dot{v}(t) < 0$ for all t . \square

3. F -solutions

For $p, q > 1$ we consider solutions ϕ of $(M_{p,\lambda})$ on (R_0, R) , $0 \leq R_0 < R \leq \infty$, and their corresponding solutions φ of $(MS_{p,\lambda})$ on (T_0, T) . The following characterization of F -solutions implies their existence.

Theorem 3.1. *The following conditions are equivalent:*

- a) ϕ is an F -solution, that is, $R_0 > 0$.
- b) $T_0 > -\infty$.
- c) $C^-(\varphi)$ is unbounded as $t \downarrow T_0$.
- d) There exists $t_0 \in (T_0, T)$ such that $u(t_0) \geq q$.
- e) $u(t) \uparrow \infty$ as $t \downarrow T_0$.

In this case, $v(t) \downarrow 0$ as $t \downarrow T_0$.

Proof. a) \leftrightarrow b) follows from the definitions for T_0 and R_0 .

b) \rightarrow c): If $C^-(\varphi)$ were bounded, the solution φ could be extended to a solution $\bar{\varphi}: [T_0, T) \rightarrow \mathbb{R}_0^+ \times \mathbb{R}_0^+$. But $\bar{\varphi}(T_0)$ is not a point on the axes, hence $\bar{\varphi}(T_0) \in \mathbb{R}^+ \times \mathbb{R}^+$ and φ could again be extended beyond T_0 . This is a contradiction.

c) \rightarrow d): If d) were not true, u would be bounded and v must be unbounded as $t \downarrow T_0$. Hence there exists $t_0 \in (T_0, T)$ with $v_0 := v(t_0) > 1$. We have $\dot{v} > v(-1 + v)$. The solution of the initial-value problem $\dot{z} = z(-1 + z)$, $z(t_0) = v_0$ is

$$z(t) = \frac{1}{1 - ce^t} \quad \text{with} \quad c := \frac{v_0 - 1}{v_0} e^{-t_0} \quad \text{on} \quad \left(-\infty, \ln \frac{1}{c}\right)$$

and we have $v(t) \leq z(t)$, $T_0 < t \leq t_0$. Because $z(t) \rightarrow 1$ ($t \rightarrow -\infty$) v is bounded on $(T_0, t_0]$. This is a contradiction.

d) \rightarrow b), e): Let $u_0 := u(t_0)$. We have $\dot{u} < u(q(t) - u) < u(q - u)$. The solution of the initial-value problem $\dot{z} = z(q - z)$, $z(t_0) = u_0$ is

$$z(t) = \frac{q}{1 + ce^{-qt}} \quad \text{with} \quad c := \frac{q - u_0}{u_0} e^{qt_0} < 0 \quad \text{on} \quad \left(-\frac{1}{q} \ln \left(-\frac{1}{c}\right), t_0\right]$$

and we have $u(t) \geq z(t)$ as long as u and z exist. This implies b) and e).

e) \rightarrow c) is trivial.

Now we prove that $v(t) \downarrow 0$ for $t \downarrow T_0$. With t_0 as in d) it is clear that $\varphi(t) \in W_+ \cap S_-(t)$ (see Section 2.6) and hence $\dot{v} > 0$ on $(T_0, t_0]$. Assume $\lim_{t \rightarrow T_0} v(t) = \bar{v} > 0$. Then

$$\dot{v}(t) = v(t)(-1 + u(t) + v(t)) > v(t)(q - 1 + v(t)) > \bar{v}(q - 1 + \bar{v}) > 0$$

for $(T_0, t_0]$ and this contradicts $\bar{v} > 0$. \square

4. E -solutions

For $p, q > 1$ we consider solutions ϕ of $(M_{p,\lambda})$ on $(0, R)$, $0 < R \leq \infty$, and their corresponding solutions φ of $(MS_{p,\lambda})$ on $(-\infty, T)$, $T := \ln R$. Let K be defined by (2.2.3). Also, let $\hat{v}(u) := v(t(u))$, where $t = t(u)$ is the inverse function of $u = u(t)$ in a neighborhood of P_3 (where $\dot{u} < 0$). In the following theorem we characterize the E -solutions, from which also their existence follows.

Theorem 4.1. *Let $p, q > 1$. The following conditions are equivalent:*

- a) ϕ is an E -solution.
b) There exists some $\alpha > 0$ such that

$$\left. \begin{aligned} \phi(r) &= \alpha - \frac{\alpha^p}{\lambda(\lambda+1)} r^\lambda [1 + o(1)], \\ \phi'(r) &= -\frac{\alpha^p}{\lambda+1} r^{\lambda-1} [1 + o(1)] \quad (r \rightarrow 0). \end{aligned} \right\} \quad (4.1)$$

- c) $r^{\lambda-1} \phi^p$ is integrable at 0, which implies that there exists $\alpha > 0$ such that

$$\left. \begin{aligned} \phi(r) &= \alpha - \int_0^r \left(1 - \frac{s}{r}\right) s K(s) \phi^p(s) ds, \quad r > 0, \\ \phi'(r) &= -\frac{1}{r^2} \int_0^r s^2 K(s) \phi^p(s) ds, \quad r > 0. \end{aligned} \right\} \quad (4.2)$$

- d) There exists $\alpha > 0$ such that for $r \rightarrow 0$

$$\begin{aligned} \phi(r) &= \alpha - \frac{\alpha^p}{\lambda(\lambda+1)} r^\lambda + \frac{p\alpha^{2p-1}}{2\lambda^2(\lambda+1)(2\lambda+1)} r^{2\lambda} \\ &\quad + \frac{\lambda\alpha^p}{2(\lambda+2)(\lambda+3)} r^{\lambda+2} + O(r^{\lambda+\min\{4, 2\lambda\}}), \\ \phi'(r) &= -\frac{\alpha^p}{(\lambda+1)} r^{\lambda-1} + \frac{p\alpha^{2p-1}}{\lambda(\lambda+1)(2\lambda+1)} r^{2\lambda-1} \\ &\quad + \frac{\lambda\alpha^p}{2(\lambda+3)} r^{\lambda+1} + O(r^{\lambda-1+\min\{4, 2\lambda\}}). \end{aligned}$$

- e) There exists $\alpha > 0$ such that

$$\begin{aligned} u(t) &= \lambda + 1 - \frac{p\alpha^{p-1}}{2\lambda+1} e^{\lambda t} - \frac{\lambda(\lambda+1)}{(\lambda+3)} e^{2t} + O(e^{\min\{4, 2\lambda\}t}), \\ v(t) &= \frac{\alpha^{p-1}}{(\lambda+1)} e^{\lambda t} \left[1 - \frac{[p(\lambda+1) - (2\lambda+1)]\alpha^{p-1}}{\lambda(\lambda+1)(2\lambda+1)} e^{\lambda t} \right. \\ &\quad \left. - \frac{\lambda(\lambda+1)}{2(\lambda+3)} e^{2t} + O(e^{\min\{4, 2\lambda\}t}) \right]. \end{aligned}$$

f) $\varphi(t) \rightarrow P_3^- (t \rightarrow -\infty)$.

In this case

$$\phi(0) := \lim_{r \rightarrow 0} \phi(r) = \alpha,$$

and

$$\lim_{u \rightarrow q} \hat{v}'(u) = \hat{v}'(q) = \begin{cases} -\frac{2\lambda+1}{p(\lambda+1)}, & \text{if } \lambda < 2, \\ -\frac{5\alpha^{p-1}}{3(p\alpha^{p-1}+6)}, & \text{if } \lambda = 2, \\ 0, & \text{if } \lambda > 2. \end{cases}$$

Proof. The existence of E -solutions follows from the possibility to solve the initial-value problem locally with $\phi(0) = \alpha$; this follows from Lemma 2.1. If we let

$$L_{0\alpha}(r) = pK(r) |\alpha + \delta|^{p-1}$$

then

$$\frac{1}{r^{\lambda+1}} \int_0^r s^2 L_{0\alpha}(s) ds$$

is bounded near zero and the same lemma implies

$$r^{1-\lambda} \phi'(r) - \frac{1}{r^{\lambda+1}} \int_0^r s^2 K(s) \alpha^p ds \rightarrow 0 \quad (r \rightarrow 0).$$

L'Hospital's rule gives

$$\lim_{r \rightarrow 0} \frac{1}{r^{\lambda+1}} \int_0^r s^2 K(s) \alpha^p ds = \frac{\alpha^p}{\lambda+1}.$$

Hence a) implies the second equation of b), and the first follows by integration.

a) \Leftrightarrow c) \Leftrightarrow b) are straightforward (see also [24, p. 8]).

b) \Rightarrow d): We will stepwise improve the asymptotics. We begin with (4.1) and have

$$\phi(r) = \alpha + O(r^\lambda) \quad (r \rightarrow 0).$$

Using

$$(\alpha + x)^p = \alpha^p + p\alpha^{p-1}x + \frac{p(p-1)\alpha^{p-2}}{2}x^2 + O(x^3), \quad (x \rightarrow 0)$$

$$\frac{1}{(1+r^2)^{\lambda/2}} = 1 - \frac{\lambda}{2}r^2 + O(r^4), \quad (r \rightarrow 0)$$

and (4.2), we get

$$\begin{aligned} r^2 \phi'(r) &= - \int_0^r s^2 K(s) \phi^p(s) ds \\ &= - \int_0^r s^\lambda [1 + O(s^2)] [\alpha^p + O(s^\lambda)] ds \\ &= - \int_0^r s^\lambda [\alpha^p + O(s^{\min\{2, \lambda\}})] ds \\ &= - \frac{\alpha^p}{(\lambda+1)} r^{\lambda+1} + O(r^{\min\{\lambda+3, 2\lambda+1\}}), \end{aligned}$$

$$\left. \begin{aligned} \phi'(r) &= -\frac{\alpha^p}{(\lambda+1)}r^{\lambda-1} + O(r^{\min\{\lambda+1, 2\lambda-1\}}), \\ \phi(r) &= \alpha - \frac{\alpha^p}{\lambda(\lambda+1)}r^\lambda + O(r^{\min\{\lambda+2, 2\lambda\}}). \end{aligned} \right\} \quad (4.3)$$

Inserting (4.3) again into (4.2) we obtain

$$\begin{aligned} & r^2 \phi'(r) \\ &= -\int_0^r s^2 K(s) \phi^p(s) ds \\ &= -\int_0^r s^\lambda \left[1 - \frac{\lambda}{2}s^2 + O(s^4)\right] \left[\alpha^p + p\alpha^{p-1} \left(-\frac{\alpha^p}{\lambda(\lambda+1)}s^\lambda + O(s^{\min\{\lambda+2, 2\lambda\}})\right) + O(s^{2\lambda})\right] ds \\ &= -\int_0^r s^\lambda \left[\alpha^p - \frac{p\alpha^{2p-1}}{\lambda(\lambda+1)}s^\lambda - \frac{\lambda\alpha^p}{2}s^2 + O(s^{\min\{4, 2\lambda\}})\right] ds \\ &= -\frac{\alpha^p}{(\lambda+1)}r^{\lambda+1} + \frac{p\alpha^{2p-1}}{\lambda(\lambda+1)(2\lambda+1)}r^{2\lambda+1} + \frac{\lambda\alpha^p}{2(\lambda+3)}r^{\lambda+3} + O(r^{\min\{\lambda+5, 3\lambda+1\}}), \end{aligned}$$

and the equation for ϕ follows by integration (note that $\min\{4, \lambda + 2, 2\lambda\} = \min\{4, 2\lambda\}$).

d) \Rightarrow e), f) : we have

$$\begin{aligned} u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\phi(r)^p}{-\phi'(r)} \\ &= \frac{r^{\lambda-1} \alpha^p \left[1 - \frac{\alpha^{p-1}}{\lambda(\lambda+1)}r^\lambda + O(r^{\min\{\lambda+2, 2\lambda\}})\right]^p}{(1+r^2)^{\lambda/2} \frac{\alpha^p}{(\lambda+1)}r^{\lambda-1} \left[1 - \frac{p\alpha^{p-1}}{\lambda(2\lambda+1)}r^\lambda - \frac{\lambda(\lambda+1)}{2(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}})\right]} \\ &= \left[\lambda + 1 - \frac{p\alpha^{p-1}}{\lambda}r^\lambda + O(r^{\min\{\lambda+2, 2\lambda\}})\right] \left[1 - \frac{\lambda}{2}r^2 + O(r^4)\right] \\ &\quad \times \left[1 + \frac{p\alpha^{p-1}}{\lambda(2\lambda+1)}r^\lambda + \frac{\lambda(\lambda+1)}{2(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}})\right] \\ &= \lambda + 1 - \frac{p\alpha^{p-1}}{2\lambda+1}r^\lambda - \frac{\lambda(\lambda+1)}{(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}}), \quad r = e^t \end{aligned}$$

while we also have

$$\begin{aligned} v(t) &= r \frac{-\phi'(r)}{\phi(r)} = r \frac{\frac{\alpha^p}{(\lambda+1)}r^{\lambda-1} \left[1 - \frac{p\alpha^{p-1}}{\lambda(2\lambda+1)}r^\lambda - \frac{\lambda(\lambda+1)}{2(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}})\right]}{\alpha - \frac{\alpha^p}{\lambda(\lambda+1)}r^\lambda + O(r^{\min\{\lambda+2, 2\lambda\}})} \\ &= \frac{\alpha^{p-1}}{(\lambda+1)}r^\lambda \left[1 - \frac{p\alpha^{p-1}}{\lambda(2\lambda+1)}r^\lambda - \frac{\lambda(\lambda+1)}{2(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}})\right] \\ &\quad \times \left[1 + \frac{\alpha^{p-1}}{\lambda(\lambda+1)}r^\lambda + O(r^{\min\{\lambda+2, 2\lambda\}})\right] \\ &= \frac{\alpha^{p-1}}{(\lambda+1)}r^\lambda \left[1 - \frac{[p(\lambda+1) - (2\lambda+1)]\alpha^{p-1}}{\lambda(\lambda+1)(2\lambda+1)}r^\lambda - \frac{\lambda(\lambda+1)}{2(\lambda+3)}r^2 + O(r^{\min\{4, 2\lambda\}})\right], \end{aligned}$$

$r = e^t$ which is e), and f) follows.

f) \Rightarrow a): We specialize (2.4.1) to $\bar{u} := u - q$, $\bar{v} = v$ and have

$$\left. \begin{aligned} \dot{\bar{u}} &= -q\bar{u} - pqv - \bar{u}^2 - p\bar{u}v - \lambda(\bar{u} + q)g(t), & g(t) &:= \frac{e^{2t}}{1+e^{2t}} \\ \dot{\bar{v}} &= v(\lambda + \bar{u} + v). \end{aligned} \right\} \quad (4.4)$$

Let $\varepsilon(t) := \bar{u}(t) + v(t) \rightarrow 0$. Then, for some t_0

$$v(t) = v(t_0) e^{-\lambda t_0} \cdot e^{\lambda t + \int_{t_0}^t \varepsilon(s) ds}.$$

If $|\varepsilon(t)| \leq \delta < \lambda$ for $t \leq t_0$, then

$$v(t) \leq A_0 e^{(\lambda-\delta)t}, \quad A_0 := v(t_0) e^{-(\lambda-\delta)t_0}.$$

If ε is integrable at $-\infty$, then for some $A > 0$,

$$v(t) = A e^{\lambda t} \cdot e^{\int_{-\infty}^t \varepsilon(s) ds}, \quad t \leq t_0. \quad (4.5)$$

We may assume $\frac{q}{2} - \bar{u} \geq 0$. Then

$$\begin{aligned} \frac{1}{2} (\bar{u}^2)' &= -(q + \bar{u})\bar{u}^2 - p\bar{u}^2v - \lambda\bar{u}^2g(t) - q(pv + \lambda g(t))\bar{u} \\ &\leq -(q + \bar{u})\bar{u}^2 + \frac{q}{2} \left((pv + \lambda g(t))^2 + \bar{u}^2 \right) \\ &\leq q(p^2v^2 + \lambda^2g^2(t)), \\ \bar{u}^2(t) &\leq 2q \left(p^2 \int_{-\infty}^t v^2(s) ds + \lambda^2 \int_{-\infty}^t g^2(s) ds \right) \\ &\leq 2q \left(p^2 A_0^2 e^{2(\lambda-\delta)t} + \lambda^2 e^{4t} \right), \\ \bar{u}(t) &= O \left(e^{\min(2, \lambda-\delta)t} \right). \end{aligned} \quad (4.6)$$

We have shown that ε is integrable at $-\infty$. Hence with (4.5)

$$\begin{aligned} v(t) &= A e^{\lambda t} \left[1 + O \left(\int_{-\infty}^t (\bar{u}(s) + v(s)) ds \right) \right] \\ &= A e^{\lambda t} \left[1 + O \left(e^{\min(2, \lambda)t} \right) \right]. \end{aligned}$$

Together with (4.6) we get $\bar{u}(t) = O(e^{\min(2, \lambda)t})$,

$$\begin{aligned} \dot{\bar{u}} &= -(\lambda + 1)\bar{u} - \lambda(\lambda + 1)e^{2t} - p(\lambda + 1)Ae^{\lambda t} + O(e^{\min\{4, 2\lambda\}t}), \\ \left[e^{(\lambda+1)t} \bar{u} \right]' &= -\lambda(\lambda + 1)e^{(\lambda+3)t} - p(\lambda + 1)Ae^{(2\lambda+1)t} + O(e^{\min\{\lambda+5, 3\lambda+1\}t}), \\ \bar{u}(t) &= -\frac{\lambda(\lambda + 1)}{\lambda + 3} e^{2t} - \frac{p(\lambda + 1)}{2\lambda + 1} Ae^{\lambda t} + O(e^{\min\{4, 2\lambda\}t}), \\ u(t) &= \lambda + 1 - \frac{\lambda(\lambda + 1)}{\lambda + 3} e^{2t} - \frac{p(\lambda + 1)}{2\lambda + 1} Ae^{\lambda t} + O(e^{\min\{4, 2\lambda\}t}), \end{aligned}$$

and with (4.5).

$$v(t) = Ae^{\lambda t} \left[1 - \frac{\lambda(\lambda+1)}{2(\lambda+3)} e^{2t} - \frac{p(\lambda+1) - (2\lambda+1)}{\lambda(2\lambda+1)} Ae^{\lambda t} + O(e^{\min\{4,2\lambda\}t}) \right].$$

The asymptotics for u, v imply with (2.2.5)

$$\begin{aligned} \phi^{p-1}(r) &= (1+r^2)^{\lambda/2} r^{-\lambda} u(\ln r) v(\ln r) \\ &\rightarrow qA = \alpha^{p-1} \quad (r \rightarrow 0), \end{aligned}$$

This shows a).

We still have to show the last statement:

$$\begin{aligned} \hat{v}'(q) &= \lim_{t \rightarrow -\infty} \frac{v(t)}{u(t) - q} \\ &= \lim_{t \rightarrow -\infty} \frac{\frac{\alpha^{p-1}}{(\lambda+1)} e^{\lambda t} \left[1 - \frac{[p(\lambda+1) - (2\lambda+1)]\alpha^{p-1}}{\lambda(\lambda+1)(2\lambda+1)} e^{\lambda t} - \frac{\lambda(\lambda+1)}{2(\lambda+3)} e^{2t} + O(e^{\min\{4,2\lambda\}t}) \right]}{-\frac{p\alpha^{p-1}}{2\lambda+1} e^{\lambda t} - \frac{\lambda(\lambda+1)}{(\lambda+3)} e^{2t} + O(e^{\min\{4,2\lambda\}t})} \\ &= \begin{cases} -\frac{2\lambda+1}{p(\lambda+1)}, & \text{if } \lambda < 2, \\ -\frac{5\alpha^{p-1}}{3(p\alpha^{p-1}+6)}, & \text{if } \lambda = 2, \\ 0, & \text{if } \lambda > 2. \end{cases} \end{aligned}$$

On the other hand

$$\begin{aligned} \lim_{u \rightarrow q} \hat{v}'(u) &= \lim_{t \rightarrow -\infty} \frac{v(-1+u+v)}{u(q(t)-u-pv)} \\ &= \frac{\lambda}{\lambda+1} \lim_{t \rightarrow -\infty} \frac{v}{q(t)-u-pv} = \frac{\lambda}{\lambda+1} \lim_{t \rightarrow -\infty} \frac{1}{\frac{q(t)-u}{v} - p}, \end{aligned}$$

$$\begin{aligned} \frac{q(t)-u}{v} &= \frac{q - \lambda g(t) - u}{v} \\ &= \frac{\lambda+1 - \lambda e^{2t} - \left[\lambda+1 - \frac{p\alpha^{p-1}}{2\lambda+1} e^{\lambda t} - \frac{\lambda(\lambda+1)}{(\lambda+3)} e^{2t} \right] + O(e^{\min\{4,2\lambda\}t})}{\frac{\alpha^{p-1}}{(\lambda+1)} e^{\lambda t} [1 + O(e^{\min\{2,\lambda\}t})]} \end{aligned}$$

$$= \frac{\frac{p\alpha^{p-1}}{2\lambda+1} e^{\lambda t} - \frac{2\lambda}{\lambda+3} e^{2t} + O(e^{\min\{4,2\lambda\}t})}{\frac{\alpha^{p-1}}{(\lambda+1)} e^{\lambda t} [1 + O(e^{\min\{2,\lambda\}t})]} \rightarrow \begin{cases} \frac{p(\lambda+1)}{2\lambda+1}, & \text{if } \lambda < 2, \\ \frac{3(p\alpha^{p-1}-4)}{5\alpha^{p-1}}, & \text{if } \lambda = 2, \\ -\infty, & \text{if } \lambda > 2, \end{cases}$$

so that

$$\lim_{u \rightarrow q} \hat{v}'(u) = \frac{\lambda}{\lambda+1} \lim_{t \rightarrow -\infty} \frac{1}{\frac{q(t)-u}{v} - p} = \hat{v}'(q). \quad \square$$

5. M -solutions

In this section we let $p, q > 1$, and consider solutions ϕ on $(0, R)$, $0 < R \leq \infty$, and the corresponding φ on $(-\infty, T)$, $T := \ln R$. Our aim is the characterization and description of the M -solutions. We have to distinguish between the four different cases: $q > p$ (with the three subcases, $q > p + 1$, $q = p + 1$, $p < q < p + 1$), $q = p$, $q < p$ (but $q \neq \frac{p+1}{2}$), and $q = \frac{p+1}{2}$.

5.1. The case $q > p$

In this section we show that the M -solutions are characterized by the property that their $\varphi(t)$ tend to P_2 as $t \rightarrow -\infty$. Though the existence of such solutions will follow later, the asymptotical stability of P_2 for $t \rightarrow -\infty$ follows from specializing the general linearization (2.4.1) to P_2 :

$$\begin{pmatrix} u \\ \bar{v} \end{pmatrix}' = A \begin{pmatrix} u \\ \bar{v} \end{pmatrix} + \begin{pmatrix} -u^2 - p u \bar{v} \\ u \bar{v} + \bar{v}^2 \end{pmatrix} + \begin{pmatrix} -(q-1) \frac{e^{2t}}{1+e^{2t}} u \\ 0 \end{pmatrix}$$

with

$$A := \begin{pmatrix} q-p & 0 \\ 1 & 1 \end{pmatrix}, \quad \bar{v} = v - 1.$$

A has the two positive eigenvalues $\lambda_1 = q - p > 0$, $\lambda_2 = 1$. It is possible to extend the classical theorem for asymptotic stability to this case where a time-dependent perturbation with the estimate $c \cdot e^{2t} |u|$ appears. We omit the details. Hence any solution φ starting in a small neighborhood of P_2 with a sufficiently negative t_0 will converge to P_2 as $t \rightarrow -\infty$.

Theorem 5.1. *Let $q > p$. The following conditions are equivalent:*

- a) φ is an M -solution.
- b) There exists $c > 0$ such that

$$\left. \begin{aligned} \phi(r) &= \frac{c}{r} [1 + o(1)], \\ \phi'(r) &= -\frac{c}{r^2} [1 + o(1)], \quad (r \rightarrow 0). \end{aligned} \right\} \quad (5.1.1)$$

- c) There exists $c > 0$ such that

$$\begin{aligned} u(t) &= c^{p-1} e^{(q-p)t} [1 + o(1)], \\ v(t) - 1 &= o(1) \quad (t \rightarrow -\infty). \end{aligned}$$

- d) $\varphi(t) \rightarrow P_2$ ($t \rightarrow -\infty$).

In this case

$$r^2 \phi' = -c - \int_0^r s^2 K(s) \phi^p ds, \quad (5.1.2)$$

and c is uniquely determined.

Proof. For a) \rightarrow b) we extend arguments for $(\text{EF}_{p,q})$ [5, p.504] to our case. The function $z(t) := \phi\left(\frac{1}{t}\right)$ defined on $\left(\frac{1}{R}, \infty\right)$ is an unbounded solution of the equation

$$z''(t) + f_\lambda(t) z^p = 0, \text{ where } f_\lambda(t) := \frac{1}{t^2(1+t^2)^{\lambda/2}}. \quad (5.1.3)$$

We multiply (5.1.3) with z' and integrate by parts from t to $t_0 > \frac{1}{R}$:

$$\frac{1}{2} z'(t)^2 + f_\lambda(t) \frac{z^{p+1}(t)}{p+1} - \int_{t_0}^t f'_\lambda(s) \frac{z^{p+1}(s)}{p+1} ds = c_1 + f_\lambda(t_0) \frac{z^{p+1}(t_0)}{p+1}.$$

Because $f'_\lambda < 0$, z' is bounded, and there exists a constant $K \geq 0$ such that

$$z(t) \leq K \cdot f_\lambda(t)^{-\frac{1}{p+1}}.$$

Because for $t, \tau > t_0$ with (5.1.3)

$$z'(\tau) - z'(t) = - \int_t^\tau f_\sigma(s) z^p(s) ds,$$

and

$$f_\lambda(s) z^p(s) \leq K^p f_\lambda^{1-\frac{p}{p+1}} = K^p \left[t^2(1+t^2)^{\lambda/2} \right]^{-\frac{1}{p+1}}$$

is integrable ($-(2+\lambda)\frac{1}{p+1} < -1 \iff 1+q > p+1 \iff q > p$), the limit $\lim_{t \rightarrow \infty} z'(t) =: c$ exists. Because z is unbounded, $c \geq 0$. Assume $c = 0$. Then

$$z'(t) = \int_t^\infty f_\lambda(s) z^p(s) ds. \quad (5.1.4)$$

We define

$$r_0 := \inf \{ r \in \mathbb{R}; z'(t) = O(t^r) \}.$$

Because z' is bounded, we have $r_0 \leq 0$. For $r > r_0$ we have with (5.1.4)

$$z(t) = O(t^{r+1}) \quad \text{and} \quad z'(t) = O\left(t^{-2-\lambda+p(r+1)+1}\right),$$

hence $r_0 \leq -2 - \lambda + p(r_0 + 1) + 1 = p(r_0 + 1) - (\lambda + 1) \iff r_0 + q \leq pr_0 + p \iff q - p \leq (p-1)r_0 \implies r_0 > 0$, which is a contradiction. Hence $c > 0$ and

$$\begin{aligned} z'(t) &= c[1 + o(1)] & (t \rightarrow -\infty), \\ r^2 \cdot \phi'(r) &= -c[1 + o(1)] & (r \rightarrow 0), \end{aligned}$$

and b) follows.

b) \rightarrow c): With (2.2.4) and (5.1.1)

$$\begin{aligned} u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \cdot \frac{\frac{c^p}{r^p} [1 + o(1)]}{\frac{c}{r^2} [1 + o(1)]} = c^{p-1} r^{\lambda+1-p} [1 + o(1)] \longrightarrow 0, \\ v(t) &= r \frac{\frac{c}{r^2} [1 + o(1)]}{\frac{c}{r} [1 + o(1)]} = 1 + o(1), \quad r = e^t, \quad (t \rightarrow -\infty). \end{aligned}$$

c) \rightarrow d) is immediate.

d) \rightarrow a): ϕ is an M - or an E -solution. In the latter case, $\varphi(t) \rightarrow P_3^-$ by Theorem 4.1, which is a contradiction to d).

Finally, (5.1.3) follows by integration, because $r^\lambda \phi^p$ is integrable at 0. \square

In the sequel we represent the M -solutions ϕ in the form

$$\phi = S + \Theta, \quad (5.1.5)$$

where S is a singular term of the form $S = \frac{c}{r}P(r)$ with an elementary explicitly given function P of r with $P(r) = 1 + o(1)$ ($r \rightarrow 0$), while Θ is a (regular) solution of the initial value problem

$$\left. \begin{aligned} \frac{1}{r^2} (r^2 \Theta')' &= -K(r) (\Theta + S)^p - \frac{1}{r^2} (r^2 S')', \quad 0 < r < R \\ \Theta(0) &:= \lim_{r \rightarrow 0} \Theta(r) = \beta \in \mathbb{R}. \end{aligned} \right\} \quad (5.1.6)$$

In terms of P , we have

$$\begin{aligned} \frac{1}{r^2} (r^2 \Theta')' &= -K(r) \left(\frac{c}{r}P + \Theta \right)^p - \frac{c}{r}P'' \\ &= -K(r) \frac{c^p}{r^p} P^p \left(\left(1 + \frac{r}{cP} \Theta \right)^p - 1 \right) \\ &\quad + \left(-\frac{c}{r}P'' - c^p r^{\mu-3} \frac{P^p}{(1+r^2)^{\lambda/2}} \right) \\ &=: h_1(r, \Theta) + h_2(r). \end{aligned} \quad (5.1.7)$$

Because

$$K(r) \frac{c^p}{r^p} P^p(r) = O(r^{q-p-3}) \quad \text{and} \quad \left(1 + \frac{r}{cP} \Theta \right)^p - 1 = O(r) \quad (r \rightarrow 0)$$

we remark that $h_1(r, \Theta)$ satisfies the assumptions of Lemma 2.1. We are going to derive the different forms of P and the asymptotic expansions for Θ in the subcases $q > p + 1$, $q = p + 1$, $p < q < p + 1$ in the following three subsections.

5.1.1. The case $q > p + 1$

Theorem 5.2. *Let $q > p + 1$. Then*

a) *Every M -solution ϕ has the form $\phi = S + \Theta$, where*

$$S(r) = \frac{c}{r}$$

and Θ solves (5.1.6) and has the expansions

$$\begin{aligned} \Theta(r) &= \beta - \frac{c^p}{q-p} \left[\frac{1}{q-p-1} r^{q-p-1} + \frac{p\beta}{c(q-p+1)} r^{q-p} \right] + o(r^{q-p}) \\ \Theta'(r) &= -c^p \left[\frac{1}{q-p} r^{q-p-2} + \frac{p\beta}{c(q-p+1)} r^{q-p-1} \right] + o(r^{q-p-1}) \end{aligned}$$

for some uniquely defined constants $c > 0$, $\beta \in \mathbb{R}$.

b) Conversely, given any $c > 0$ and $\beta \in \mathbf{R}$, there exists a unique solution Θ of (5.1.6) with $S(r) := \frac{c}{r}$ and $\phi := S + \Theta$ (as long as ϕ is positive) is an M -solution.

In this case we have

$$\left. \begin{aligned} \Theta(r) &= \beta - \int_0^r \left(1 - \frac{s}{r}\right) sK(s) \left[\frac{c}{s} + \Theta(s)\right]^p ds, \quad 0 < r < R, \\ u(t) &= c^{p-1} e^{(q-p)t} \left[1 + \frac{p\beta e^t}{c} + O(e^{\min\{2, q-p\}t})\right], \\ v(t) &= 1 - \frac{\beta e^t}{c} + \frac{c^{p-1} e^{(q-p)t}}{q-(p+1)} + O(e^{(q+1-p)t}), \quad (t \rightarrow -\infty), \end{aligned} \right\} \quad (5.1.8)$$

and

$$\lim_{u \rightarrow 0^+} \hat{v}'(u) = \hat{v}'(0) = \begin{cases} \infty, & \text{if } \beta < 0, \\ \frac{1}{q-(p+1)}, & \text{if } \beta = 0, \\ -\infty, & \text{if } \beta > 0. \end{cases}$$

Proof. a) Let $c > 0$ be determined by Theorem 5.1. Then $\frac{1}{r^2} \int_0^r s^2 K(s) \phi^p ds = O(r^{\lambda-p-1})$ is integrable because $\lambda - p > 0 \Leftrightarrow q > p + 1$. Hence, we have from (5.1.2) with $r_0 \in (0, R)$

$$\begin{aligned} \phi(r) &= \frac{c}{r} - \frac{c}{r_0} + \phi(r_0) - \int_{r_0}^r \frac{d\sigma}{\sigma^2} \int_0^\sigma s^2 K(s) \phi^p(s) ds, \\ &= \frac{c}{r} + \beta - \int_{r_0}^r \frac{d\sigma}{\sigma^2} \int_0^\sigma s^2 K(s) \phi^p(s) ds \\ &= \frac{c}{r} + \Theta(r), \quad 0 < r < R. \end{aligned}$$

We define $S(r) := \frac{c}{r}$ and see that Θ satisfies (5.1.5) and (5.1.6) with $\frac{1}{r^2} (r^2 S')' = 0$. Hence (5.1.8) follows by partial integration. Any equation $\frac{c_1}{r} + \Theta_1(r) = \frac{c_2}{r} + \Theta_2(r)$ implies $c_1 = c_2$ and $\Theta_1(r) = \Theta_2(r)$, which shows the uniqueness of c, Θ and β . Now we remember $\Theta(r) = \beta + o(1)$ and have

$$\begin{aligned} \Theta'(r) &= -\frac{1}{r^2} \int_0^r s^2 K(s) \left[\frac{c}{s} + \Theta(s)\right]^p ds \\ &= -\frac{c^p}{r^2} \int_0^r s^{\lambda-p} (1+s^2)^{-\frac{\lambda}{2}} \left[1 + \frac{s\Theta(s)}{c}\right]^p ds \\ &= -\frac{c^p}{r^2} \int_0^r s^{\lambda-p} [1 + O(s^2)] \left[1 + \frac{p\beta}{c}s + o(s)\right] ds \\ &= -\frac{c^p}{r^2} \int_0^r s^{\lambda-p} \left[1 + \frac{p\beta}{c}s + o(s)\right] ds \\ &= -\frac{c^p}{r^2} \int_0^r \left[s^{\lambda-p} + \frac{p\beta}{c}s^{\lambda+1-p} + o(s^{\lambda+1-p})\right] ds \\ &= -c^p \left[\frac{1}{\lambda+1-p} r^{\lambda-1-p} + \frac{p\beta}{c(\lambda+2-p)} r^{\lambda-p}\right] + o(r^{\lambda-p}), \end{aligned}$$

and the expansion for Θ follows by integration. This shows a). b) Given c, β , we define $S(r) := \frac{c}{r}$. Lemma 2.1 tells us that (5.1.6) (where $(r^2 S')' = 0$) has a unique solution Θ , and $\phi := S + \Theta$ is an M -solution. Finally, using (2.2.4) and $\phi = S + \Theta$ we get

$$\begin{aligned}
u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\left[\frac{c}{r} + \beta - \frac{c^p}{\lambda+1-p} \left[\frac{1}{\lambda-p} r^{\lambda-p} + \frac{p\beta}{c(\lambda+2-p)} r^{\lambda+1-p} \right] + o(r^{\lambda+1-p}) \right]^p}{\frac{c}{r^2} + c^p \left[\frac{1}{\lambda+1-p} r^{\lambda-1-p} + \frac{p\beta}{c(\lambda+2-p)} r^{\lambda-p} \right] + o(r^{\lambda-p})} \\
&= \frac{c^{p-1} r^{\lambda+1-p}}{(1+r^2)^{\lambda/2}} \frac{\left[1 + \frac{\beta r}{c} + O(r^{\lambda+1-p}) \right]^p}{1 + O(r^{\lambda+1-p})} \\
&= \frac{c^{p-1} r^{\lambda+1-p}}{1 + O(r^2)} \frac{1 + p \left[\frac{\beta r}{c} + O(r^{\lambda+1-p}) \right] + O(r^2)}{1 + O(r^{\lambda+1-p})} \\
&= c^{p-1} r^{\lambda+1-p} \frac{1 + p \frac{\beta r}{c} + O(r^{\min\{2, \lambda+1-p\}})}{1 + O(r^{\min\{2, \lambda+1-p\}})} \\
&= c^{p-1} r^{\lambda+1-p} \left[1 + \frac{p\beta r}{c} + O(r^{\min\{2, \lambda+1-p\}}) \right]
\end{aligned}$$

and similarly

$$\begin{aligned}
v(t) &= r \frac{-\phi'(r)}{\phi(r)} \\
&= r \frac{\frac{c}{r^2} + c^p \left[\frac{1}{\lambda+1-p} r^{\lambda-1-p} + \frac{p\beta}{c(\lambda+2-p)} r^{\lambda-p} \right] + o(r^{\lambda-p})}{\frac{c}{r} + \beta - \frac{c^p}{\lambda+1-p} \left[\frac{1}{\lambda-p} r^{\lambda-p} + \frac{p\beta}{c(\lambda+2-p)} r^{\lambda+1-p} \right] + o(r^{\lambda+1-p})} \\
&= \frac{1 + c^{p-1} r^{\lambda+1-p} \left[\frac{1}{\lambda+1-p} + \frac{p\beta}{c(\lambda+2-p)} r \right] + o(r^{\lambda+2-p})}{1 + \frac{\beta r}{c} - \frac{c^{p-1} r^{\lambda+1-p}}{\lambda+1-p} \left[\frac{1}{\lambda-p} + \frac{p\beta}{c(\lambda+2-p)} r \right] + o(r^{\lambda+2-p})} \\
&= 1 - \frac{\beta r}{c} + \frac{c^{p-1} r^{\lambda+1-p}}{\lambda-p} + O(r^{\lambda+2-p}).
\end{aligned}$$

Hence

$$\begin{aligned} \hat{v}'(0) &= \lim_{t \rightarrow -\infty} \frac{v(t) - 1}{u(t)} \\ &= \lim_{t \rightarrow -\infty} \frac{-\frac{\beta e^t}{c} + \frac{c^{p-1} e^{(\lambda+1-p)t}}{\lambda-p} + O(e^{(\lambda+2-p)t})}{c^{p-1} e^{(\lambda+1-p)t} \left[1 + \frac{p\beta e^t}{c} + O(e^{\min\{2, \lambda+1-p\}t}) \right]} \\ &= \begin{cases} \infty, & \text{if } \beta < 0, \\ \frac{1}{\lambda-p}, & \text{if } \beta = 0, \\ -\infty, & \text{if } \beta > 0. \end{cases} \end{aligned}$$

and

$$\lim_{u \rightarrow 0} \hat{v}'(u) = \lim_{t \rightarrow -\infty} \frac{\dot{v}(t)}{\dot{u}(t)} = \lim_{t \rightarrow -\infty} \frac{v(-1+u+v)}{u(q(t)-u-pv)} = \frac{\hat{v}'(0) + 1}{q-p} = \hat{v}'(0). \quad \square$$

(5.1.9)

5.1.2. The case $q = p + 1$

Theorem 5.3. *Let $q = p + 1$, i.e. $\lambda = p$. Then:*

a) *Every M -solution has the form $\phi = S + \Theta$, where*

$$S(r) = \frac{c}{r} - c^p \ln r \quad (5.1.10)$$

and Θ solves (5.1.6) and has the expansions

$$\Theta(r) = \beta + \frac{p}{2} c^{2p-1} r \ln r - \frac{p}{4} c^{p-1} (3c^p + 2\beta) r + o(r)$$

$$\Theta'(r) = \frac{p}{2} c^{2p-1} \ln r - \frac{p}{4} c^{p-1} (3c^p + 2\beta) + o(1), \quad (r \rightarrow 0)$$

for some uniquely defined constants $c > 0$, $\beta \in \mathbb{R}$.

b) *Conversely, given any $c > 0$ and $\beta \in \mathbb{R}$, there exists a unique solution Θ of (5.1.6) with S given by (5.1.10), and $\phi := S + \Theta$ is an M -solution (as long as ϕ is positive).*

In this case Θ satisfies (5.1.7) with $P(r) = 1 - c^{p-1} r \ln r$, $P''(r) = -\frac{c^{p-1}}{r}$, so that

$$h_2(r) = \frac{c^p}{r^2} \left(1 - \frac{P^p(r)}{(1+r^2)^{p/2}} \right) = \frac{p}{2} c^p (1 + o(1)), \quad (5.1.11)$$

$$u(t) = c^{p-1} e^t [1 - p c^{p-1} t e^t + O(e^t)],$$

$$v(t) = 1 + c^{p-1} t e^t + O(e^t),$$

and

$$\lim_{u \rightarrow 0^+} \hat{v}'(u) = \hat{v}'(0) = -\infty. \quad (5.1.12)$$

Proof. a) Let $c > 0$ be given by Theorem 5.1, which also gives

$$\begin{aligned} s^2 K(s) \phi^p(s) &= \frac{s^\lambda}{(1+s^2)^{\lambda/2}} \frac{c^p}{s^p} \left(\frac{s}{c} \phi(s) \right)^p = c^p (1 + O(s^2)) (1 + o(1))^p \\ &= c^p (1 + o(1)), \\ \phi'(r) &= -\frac{c}{r^2} - \frac{1}{r^2} \int_0^r s^2 K(s) \phi^p(s) ds = -\frac{c}{r^2} - \frac{c^p}{r} + o\left(\frac{1}{r}\right), \\ \phi(r) &= \frac{c}{r} - c^p \ln r + o(\ln r) = \frac{c}{r} (1 - c^{p-1} r \ln r + o(r \ln r)), \\ s^2 K(s) \phi^p(s) &= c^p (1 - pc^{p-1} s \ln s + o(s \ln s)), \\ \phi'(r) &= -\frac{c}{r^2} - \frac{c^p}{r} + \frac{p}{2} c^{2p-1} \ln r + o(\ln r), \\ \phi(r) &= \frac{c}{r} - c^p \ln r + \frac{p}{2} c^{2p-1} r \ln r + o(r \ln r). \end{aligned}$$

We define $S(r) := \frac{c}{r} - c^p \ln r$ and $\Theta(r) := \phi(r) - S(r)$. Then Θ satisfies the differential equation in (5.1.6), and because $\Theta' = \phi' - S'$ is integrable, $\Theta(0) =: \beta$ exists. Again,

$$\begin{aligned} s^2 K(s) \phi^p(s) &= s^2 K(s) \frac{c^p}{s^p} \left(\frac{s}{c} \phi(s) \right)^p \\ &= c^p (1 + O(s^2)) \left(1 - c^{p-1} s \ln s + \frac{s}{c} (\beta + o(1)) \right)^p, \\ \phi'(r) &= -\frac{c}{r^2} - \frac{1}{r^2} \int_0^r s^2 K(s) \phi^p(s) ds \\ &= -\frac{c}{r^2} - \frac{c^p}{r^2} \int_0^r \left(1 - pc^{p-1} s \ln s + \frac{p\beta}{c} s + o(s) \right) ds \\ &= -\frac{c}{r^2} - \frac{c^p}{r} + \frac{p}{2} c^{2p-1} \ln r - \frac{p}{4} c^{2p-1} - \frac{p\beta}{2} c^{p-1} + o(1), \\ \Theta'(r) &= \frac{p}{2} c^{2p-1} \ln r - \frac{p}{4} c^{2p-1} - \frac{p\beta}{2} c^{p-1} + o(1). \end{aligned}$$

The expansion for Θ follows by integration and the uniqueness is shown as above. This shows a).

b) Given c, β , we define $S(r) := \frac{c}{r} - c^p \ln r = \frac{c}{r} P(r)$ and (5.1.11) follows. Together with the remark following (5.1.7) we see that the assumptions of Lemma 2.1 are satisfied, and (5.1.6) has a unique solution Θ . Then $\phi := S + \Theta$ is an M -solution.

Finally,

$$\begin{aligned} u(t) &= \frac{r^{\lambda-1} \phi^\lambda(r)}{(1+r^2)^{\lambda/2} - \phi'(r)} = \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \cdot \frac{\frac{c^p}{r^p} [1 - c^{p-1} r \ln r + O(r)]^p}{-\frac{c}{r^2} [1 - c^{p-1} r + \frac{p}{2} c^{p-2} r^2 \ln r + O(r^2 \ln r)]} \\ &= \frac{c^{p-1} r^{\lambda+1-p}}{(1+r^2)^{\lambda/2}} \frac{[1 - c^{p-1} r \ln r + O(r)]^p}{1 - c^{p-1} r + \frac{p}{2} c^{p-2} r^2 \ln r + O(r^2 \ln r)} \\ &= c^{p-1} r [1 - pc^{p-1} r \ln r + O(r)], \quad r = e^t \end{aligned}$$

and

$$\begin{aligned}
v(t) &= r \frac{-\phi'(r)}{\phi(r)} = r \frac{\frac{c}{r^2} + c^\lambda \left[\frac{1}{r} - \frac{\lambda c^{\lambda-1}}{2} \ln r + \frac{\lambda(c^\lambda + 2\beta)}{4c} \right] + o(1)}{\frac{c}{r} - c^\lambda \ln r + \beta r + \frac{p}{2} c^{2p-1} r \ln r + O(r)} \\
&= \frac{1 + c^{\lambda-1} \left[r - \frac{\lambda c^{\lambda-1}}{2} r^2 \ln r + \frac{\lambda(c^\lambda + 2\beta)}{4c} r^2 \right] + o(r^2)}{1 - c^{\lambda-1} r \ln r + \frac{\beta}{c} r + \frac{p}{2} c^{2p-2} r^2 \ln r + O(r^2)} \\
&= 1 + c^{\lambda-1} r \ln r + O(r),
\end{aligned}$$

hence

$$\begin{aligned}
\hat{v}'(0) &= \lim_{t \rightarrow -\infty} \frac{v(t) - 1}{u(t)} \\
&= \lim_{t \rightarrow -\infty} \frac{c^{\lambda-1} t e^t + O(e^t)}{c^{\lambda-1} e^t [1 - \lambda c^{\lambda-1} t e^t + O(e^t)]} \\
&= -\infty.
\end{aligned}$$

Because $\lim_{u \rightarrow 0} \hat{v}'(u) = \frac{\hat{v}'(0)+1}{q-p}$ as in (5.1.9) this shows (5.1.12). \square

5.1.3. The case $p < q < p+1$

Theorem 5.4. *Let $p < q < p+1$. Let $\mu := q-p \in (0, 1)$ and $k_0 \in \mathbb{N}$ such that $k_0 \mu < 1 \leq (k_0 + 1) \mu$. Then there exist constants $d_1 := \frac{1}{\mu(1-\mu)}$, d_2, \dots, d_{k_0+2} depending upon p and q with the following properties:*

a) *Every M -solution ϕ has the form $\phi = S + \Theta$, where*

$$S(r) = \left\{ \begin{array}{l} \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i \rho^i \right], \quad \rho := c^{p-1} r^\mu, \\ \quad \text{if } (k_0 + 1) \mu > 1, \\ \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i \rho^i + d_{k_0+1} c^{(k_0+1)(p-1)} r \ln r \right], \\ \quad \text{if } (k_0 + 1) \mu = 1, \end{array} \right\} \quad (5.1.13)$$

and Θ is a solution of (5.1.6) with the expansions

$$\Theta(r) = \left\{ \begin{array}{l} \beta + d_{k_0+1} c^{(k_0+1)(p-1)+1} r^{(k_0+1)\mu-1} + o(r^{(k_0+1)\mu-1}), \\ \quad \text{if } (k_0 + 1) \mu > 1, \\ \beta + d_{k_0+2} c^{(k_0+1)(p-1)+1} r^\mu \ln r + o(r^\mu \ln r), \\ \quad \text{if } (k_0 + 1) \mu = 1, \end{array} \right.$$

$$\Theta'(r) = \begin{cases} ((k_0+1)\mu - 1) d_{k_0+1} c^{(k_0+1)(p-1)+1} r^{(k_0+1)\mu-2} + o(r^{(k_0+1)\mu-2}), \\ \quad \text{if } (k_0+1)\mu > 1, \\ \mu d_{k_0+2} c^{(k_0+1)(p-1)+1} r^{\mu-1} \ln r + o(r^{\mu-1} \ln r), \\ \quad \text{if } (k_0+1)\mu = 1, \end{cases}$$

for uniquely determined constants $c > 0$, $\beta \in \mathbb{R}$.

- b) Conversely, given any $c > 0$ and $\beta \in \mathbb{R}$, there exists a unique solution Θ of (5.1.6) with S given by (5.1.13) and $\phi := S + \Theta$ is an M -solution (as long as ϕ is positive).

In this case Θ satisfies (5.1.7) with $P(r) = \frac{r}{c} S(r)$ and

$$h_2(r) = \begin{cases} O(r^{(k_0+1)\mu-3}), & \text{if } (k_0+1)\mu > 1, \\ O(r^{\mu-2} \ln r), & \text{if } (k_0+1)\mu = 1, \end{cases}$$

$$u(t) = c^{p-1} e^{\mu t} \left[1 + \frac{q-1}{\mu(1-\mu)} c^{p-1} e^{\mu t} \right] + o(e^{2\mu t}),$$

$$v(t) = 1 - \frac{1}{\mu} c^{p-1} e^{\mu t} + o(e^{\mu t}),$$

and

$$\lim_{u \rightarrow 0} \hat{v}'(u) = \hat{v}'(0) = -\frac{1}{1-\mu}.$$

Proof. a) Let c be determined by Theorem 5.1, which also gives

$$s^2 K(s) \phi^p(s) = c^p s^{\mu-1} (1 + o(1)),$$

$$\phi'(r) = -\frac{c}{r^2} - \frac{c^p}{\mu} r^{\mu-2} + o(r^{\mu-2}),$$

$$\begin{aligned} \phi(r) &= \frac{c}{r} - \frac{c^p}{\mu(\mu-1)} r^{\mu-1} + o(r^{\mu-1}) + A \\ &= \frac{c}{r} + \frac{c^p}{\mu(1-\mu)} r^{\mu-1} + o(r^{\mu-1}) = \frac{c}{r} [1 + d_1 c^{p-1} r^\mu + o(r^\mu)] \end{aligned}$$

(see the proof of a) of Theorem 5.3), because $\mu < 1$. More generally, assume

$$\phi(r) = \frac{c}{r} \left[1 + d_1 c^{p-1} r^\mu + \dots + d_k (c^{p-1} r^\mu)^k + o(r^{k\mu}) \right] \quad (5.1.14)$$

with constants $d_i = d_i(p, q)$, $i = 1, \dots, k$, where $k\mu < 1$. Let us introduce the notation

$$\Sigma := d_1 \rho + \dots + d_k \rho^k, \quad \sigma := \Sigma + o(\rho^k).$$

Then for $j \in \mathbb{N}$,

$$\Sigma^j = O(\rho^j), \quad \sigma^j = \sum_{l=0}^j \binom{j}{l} \Sigma^{j-l} O(\rho^{\ell k}) = \Sigma^j + o(\rho^k),$$

so that

$$\begin{aligned}
(1 + \sigma)^p &= 1 + p\sigma + \cdots + \frac{p(p-1)(p-k+1)}{k!} \sigma^k + O(\sigma^{k+1}) \\
&= 1 + p\Sigma + \cdots + \frac{p(p-1)(p-k+1)}{k!} \Sigma^k + O(\rho^{k+1}) + o(\rho^k) + O(\Sigma^{k+1}) \\
&= 1 + \tilde{d}_1 \rho + \cdots + \tilde{d}_k \rho^k + o(\rho^k)
\end{aligned}$$

with appropriate definition of $\tilde{d}_i = \tilde{d}_i(p, q)$, $i = 1, \dots, k$. Hence

$$\begin{aligned}
s^2 K(s) \phi^p(s) &= c^p s^{\mu-1} (1 + O(s^2)) (1 + \sigma)^p \\
&= c^p s^{\mu-1} \left[1 + \sum_{i=1}^k \tilde{d}_i c^{i(p-1)} s^{i\mu} + o(s^{k\mu}) \right] \\
&= c \left[c^{p-1} s^{\mu-1} + \sum_{i=1}^k \tilde{d}_i c^{(i+1)(p-1)} s^{(i+1)\mu-1} + o(s^{(k+1)\mu-1}) \right], \\
\int_0^r s^2 K(s) \phi^p(s) ds &= c \left[c^{p-1} \frac{r^\mu}{\mu} + \sum_{i=1}^k \tilde{d}_i c^{(i+1)(p-1)} \frac{r^{(i+1)\mu}}{(i+1)\mu} + o(s^{(k+1)\mu}) \right] \\
\phi'(r) &= -\frac{c}{r^2} - c \left[\frac{1}{\mu} c^{p-1} r^{\mu-2} + \sum_{i=1}^k \frac{\tilde{d}_i}{(i+1)\mu} c^{(i+1)(p-1)} r^{(i+1)\mu-2} + o(r^{(k+1)\mu-2}) \right] \\
&= -\frac{c}{r^2} - c \left[\frac{1}{\mu} c^{p-1} r^{\mu-2} + \sum_{i=2}^k \frac{\tilde{d}_{i-1}}{i\mu} c^{i(p-1)} r^{i\mu-2} \right. \\
&\quad \left. + \frac{\tilde{d}_k}{(k+1)\mu} c^{(k+1)(p-1)} r^{(k+1)\mu-2} + o(r^{(k+1)\mu-2}) \right].
\end{aligned}$$

The first k terms in the bracket are not integrable at 0 because $k\mu < 1$. Integrating from $r_0 > 0$ to r we get, with some $\beta \in \mathbb{R}$,

$$\begin{aligned}
\phi(r) &= \frac{c}{r} + c \left[\frac{1}{\mu(1-\mu)} c^{p-1} r^{\mu-1} + \sum_{i=2}^k \frac{\tilde{d}_{i-1}}{i\mu(1-i\mu)} c^{i(p-1)} r^{i\mu-1} \right] \\
&\quad + \beta - c \frac{\tilde{d}_k}{(k+1)\mu} c^{(k+1)(p-1)} \int_{r_0}^r s^{(k+1)\mu-2} ds + o(r^{(k+1)\mu-1}).
\end{aligned}$$

If $(k+1)\mu < 1$, we find $d_i = \frac{\tilde{d}_{i-1}}{i\mu(1-i\mu)}$, $i = 2, \dots, k$ and

$$\phi(r) = \frac{c}{r} \left[1 + d_1 \rho + \cdots + d_k \rho^k + d_{k+1} \rho^{k+1} + o(r^{(k+1)\mu}) \right],$$

where $d_{k+1} := \frac{\tilde{d}_k}{(k+1)\mu(1-(k+1)\mu)}$, that is, (5.1.14) holds for $k+1$. By induction, we can repeat the process until the biggest k with $k\mu < 1$, which is k_0 . Hence

(5.1.14) also holds for k_0 . If $(k_0 + 1)\mu > 1$, we have for $k = k_0$

$$\phi(r) = \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i \rho^i \right] + \beta + d_{k_0+1} c^{(k_0+1)(p-1)+1} r^{(k_0+1)\mu} + o\left(r^{(k_0+1)\mu}\right).$$

The singular term S and the regular term Θ can now be read off. In the case $(k_0 + 1)\mu = 1$, we let $D = -\tilde{d}_{k_0} c^{(k_0+1)(p-1)+1}$ and we have

$$\begin{aligned} \phi(r) &= \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i s^i \right] + \beta - \tilde{d}_{k_0} c^{(k_0+1)(p-1)+1} (\ln r - \ln r_0) + o(\ln r) \\ &= \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i s^i + D \cdot r \ln r + o(r \ln r) \right]. \end{aligned}$$

We need a further iteration. We let

$$\Sigma := d_1 \rho + \cdots + d_{k_0} \rho^{k_0}, \quad L(r) := D r \ln r, \quad \sigma := \Sigma + L + o(r \ln r).$$

Then for $j \geq 2$,

$$\begin{aligned} \sigma^j &= \Sigma^j + \sum_{i=1}^{j-1} \binom{j}{i} \Sigma^{j-i} (L + o(r \ln r))^i + (L + o(r \ln r))^j \\ &= \Sigma^j + \sum_{i=1}^{j-1} O(\rho^{j-i}) L^i (1 + o(1)) + L^j (1 + o(1)) = \Sigma^j + o(r \ln r), \end{aligned}$$

so that

$$\begin{aligned} (1 + \sigma)^p &= 1 + p\sigma + \frac{p(p-1)}{2} \sigma^2 + \cdots + \frac{p(p-1) \cdots (p-k_0+1)}{k!} \sigma^{k_0} + O(\sigma^{k_0+1}) \\ &= 1 + p\Sigma + \frac{p(p-1)}{2} \Sigma^2 + \cdots + \frac{p(p-1) \cdots (p-k_0+1)}{k!} \Sigma^{k_0} + pL \\ &\quad + O(\Sigma^{k_0+1}) + o(r \ln r) \\ &= 1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i + O(\rho^{k_0+1}) + pL + O(\rho^{k_0+1}) + o(r \ln r) \\ &= 1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i + pL + o(r \ln r) \end{aligned}$$

for suitably chosen $\tilde{d}_i = \tilde{d}_i(p, q)$, $i = 1, \dots, k_0$; note that $\rho^{k_0+1} = c^{(k_0+1)(p-1)}r$. Hence as above,

$$\begin{aligned}
s^2 K(s) \phi^p(s) &= c^p s^{\mu-1} \left[1 + \sum_{i=1}^{k_0} \tilde{d}_i c^{i(p-1)} s^{i\mu} + pDs \ln s + o(s \ln s) \right] \\
&= c \left[c^{p-1} s^{\mu-1} + \sum_{i=1}^{k_0} \tilde{d}_i c^{(i+1)(p-1)} s^{(i+1)\mu-1} + pDs^\mu \ln s + o(s^\mu \ln s) \right], \\
\phi'(r) &= -\frac{c}{r^2} - c \left[\frac{1}{\mu} c^{p-1} r^{\mu-2} + \sum_{i=2}^{k_0} \frac{\tilde{d}_{i-1}}{i\mu} c^{i(p-1)} r^{i\mu-2} + \tilde{d}_{k_0} c^{(k_0+1)(p-1)} r^{-1} \right. \\
&\quad \left. - \frac{pD}{\mu+1} r^{\mu-1} \ln r + o(r^{\mu-1} \ln r) \right] \\
\phi(r) &= \frac{c}{r} + c \left[\frac{1}{\mu(1-\mu)} c^{p-1} r^{\mu-1} + \sum_{i=2}^{k_0} \frac{\tilde{d}_{i-1}}{i\mu(1-i\mu)} c^{i(p-1)} r^{i\mu-1} + d_{k_0+1} c^{(k_0+1)(p-1)} \ln r \right] \\
&\quad + \beta + \frac{pD}{\mu(\mu+1)} c^{(k_0+1)(p-1)+1} r^\mu \ln r + o(r^\mu \ln r) \\
&= \frac{c}{r} \left[1 + \sum_{i=1}^{k_0} d_i \rho^i + d_{k_0+1} c^{(k_0+1)(p-1)} r \ln r \right] \\
&\quad + \beta + d_{k_0+2} c^{(k_0+1)(p-1)+1} r^\mu \ln r + o(r^\mu \ln r),
\end{aligned}$$

where $d_{k_0+1} := \tilde{d}_{k_0}$, $d_{k_0+2} = \frac{-p \tilde{d}_{k_0}}{\mu(\mu+1)}$, and S and Θ can be read off again. This shows a).

b) With the constants given we define S as in a) and $P(r) := \frac{r}{c} S(r)$. Let $(k_0+1)\mu = 1$. Using the relation between the d_i 's and the \tilde{d}_i 's one easily shows

$$P''(r) = -c^{p-1} r^{\mu+2} \left[1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i \right].$$

Furthermore, we remember

$$P^p(r) = (1 + \Sigma + L)^p = 1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i + pL + o(r \ln r).$$

Hence with (5.1.7),

$$\begin{aligned}
h_2(r) &= -\frac{c}{r} P''(r) - c^p \frac{r^{\mu-3}}{(1+r^2)^{\lambda/2}} P^p(r) \\
&= -\frac{c}{r} P''(r) - c^p r^{\mu-3} P^p(r) + c^p r^{\mu-3} P^p(r) \left(1 - (1+r^2)^{-\lambda/2} \right) \\
&= c^p r^{\mu-3} \left[1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i - \left(1 + \sum_{i=1}^{k_0} \tilde{d}_i \rho^i + pL + o(r \ln r) \right) \right] + O(r^{\mu-1}) \\
&= r^{\mu-3} \cdot O(r \ln r) = O(r^{\mu-2} \ln r).
\end{aligned}$$

The case $(k_0 + 1)\mu > 1$ is similar. Hence h_2 satisfies the assumptions of Lemma 2.1, and together with the remark following (5.1.7) we see that there exists a unique solution Θ of (5.1.6).

Finally,

$$\begin{aligned}
u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\phi(r)^p}{-\phi'(r)} \\
&= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\left[\frac{c}{r} \left(1 + \frac{1}{\mu(1-\mu)} c^{p-1} r^\mu + \dots + c^{k_0(p-1)} d_{k_0} r^{k_0\mu} + o(r^{k_0\mu}) \right) \right]^p}{\frac{c}{r^2} \left[1 + \frac{1}{\mu} c^{p-1} r^\mu + \dots + \frac{d_{k_0}}{(k_0+1)\mu} (c^{p-1} r^\mu)^{k_0+1} + o(r^{\mu(k_0+1)}) \right]} \\
&= \frac{c^{p-1} r^{q-p}}{(1+r^2)^{\lambda/2}} \frac{\left[1 + \frac{1}{\mu(1-\mu)} c^{p-1} r^\mu + \dots + c^{k_0(p-1)} d_{k_0} r^{k_0\mu} + o(r^{k_0\mu}) \right]^p}{1 + \frac{1}{\mu} c^{p-1} r^\mu + \dots + \frac{d_{k_0}}{(k_0+1)\mu} (c^{p-1} r^\mu)^{k_0+1} + o(r^{\mu(k_0+1)})} \\
&= c^{p-1} r^{q-p} \left[1 + \frac{q-1}{\mu(1-\mu)} c^{p-1} r^\mu + o(r^\mu) \right],
\end{aligned}$$

and

$$\begin{aligned}
v(t) &= r \frac{-\phi'(r)}{\phi(r)} = r \frac{\frac{c}{r^2} \left[1 + \frac{1}{\mu} c^{p-1} r^\mu + \dots + \frac{d_{k_0}}{(k_0+1)\mu} (c^{p-1} r^\mu)^{k_0+1} + o(r^{\mu(k_0+1)}) \right]}{\frac{c}{r} \left(1 + \frac{1}{\mu(1-\mu)} c^{p-1} r^\mu + \dots + c^{k_0(p-1)} d_{k_0} r^{k_0\mu} + o(r^{k_0\mu}) \right)} \\
&= \frac{1 + \frac{1}{\mu} c^{p-1} r^\mu + \dots + \frac{d_{k_0}}{(k_0+1)\mu} (c^{p-1} r^\mu)^{k_0+1} + o(r^{\mu(k_0+1)})}{1 + \frac{1}{\mu(1-\mu)} c^{p-1} r^\mu + \dots + c^{k_0(p-1)} d_{k_0} r^{k_0\mu} + o(r^{k_0\mu})} \\
&= 1 - \frac{1}{1-\mu} c^{p-1} r^\mu + o(r^\mu).
\end{aligned}$$

Hence, with (5.1.9),

$$\begin{aligned}
\hat{v}'(0) &= \lim_{t \rightarrow -\infty} \frac{v(t) - 1}{u(t)} = \lim_{t \rightarrow -\infty} \frac{-\frac{1}{1-\mu} c^{p-1} e^{\mu t} + o(e^{\mu t})}{c^{p-1} e^{\mu t} \left[1 + \frac{p+\mu-1}{\mu(1-\mu)} c^{p-1} e^{\mu t} + o(e^{\mu t}) \right]} \\
&= -\frac{1}{1-\mu} = \lim_{u \rightarrow 0} \hat{v}'(\mu). \quad \square
\end{aligned}$$

5.2. The case $q = p$

There exist solutions φ for which $\varphi(t) \rightarrow P_2$ ($t \rightarrow -\infty$). In fact, because $q > \frac{p+1}{2}$, by Lemma 2.3, if $\varphi(t_0) \in G_-$, $\varphi(t) \in G_-$ for $t \leq t_0$. By Thieme's theorem and the absence of Cases \textcircled{B} and \textcircled{C} for the limit system (EFS $_{p,q}$) [4, case $n = m + 3$], φ must converge to a stationary point. It cannot converge to P_1 by Lemma 2.4, and $P_3^- \notin \bar{G}_-$ because $q > \frac{p+1}{2}$. Hence φ converges to P_2 .

Before proving the main theorem of this section we first establish the following lemmas.

Lemma 5.5. *Consider the differential equation $y' = r(t)y + b(t)$ on $(-\infty, t_0]$. If $r(t) \rightarrow r_0 > 0$, $b(t) \rightarrow b_0 \in \mathbb{R}$ ($t \rightarrow -\infty$), then*

$$y(t) \rightarrow -\frac{b_0}{r_0} \quad (t \rightarrow -\infty).$$

Proof. We let $R(t) = \int_{t_0}^t r(s) ds$ and have

$$\begin{aligned} y(t) &= e^{R(t)} \left[y(t_0) - \int_t^{t_0} b(s) e^{-R(s)} ds \right] \\ &= e^{R(t)} \left[y(t_0) - \int_t^{t_0} \left(-r(s) \frac{b_0}{r_0} + b(s) \right) e^{-R(s)} ds \right] - \frac{b_0}{r_0} (1 - e^{R(t)}), \quad t \leq t_0. \end{aligned}$$

Given $\varepsilon > 0$, there exists $t_1 < t_0$ such that for $s \leq t_1$

$$\left| -r(s) \frac{b_0}{r_0} + b(s) \right| < \frac{\varepsilon}{2} r_0 < \varepsilon r(s).$$

For $t \leq t_1$ we write $\int_t^{t_0} = \int_t^{t_1} + \int_{t_1}^{t_0}$ and then have

$$\begin{aligned} \left| \int_t^{t_1} \left(-r(s) \frac{b_0}{r_0} + b(s) \right) e^{-R(s)} ds \right| &< \varepsilon \int_t^{t_1} r(s) e^{-R(s)} ds \\ &= \varepsilon \left(e^{-R(t)} - e^{-R(t_1)} \right) < \varepsilon e^{-R(t)}, \end{aligned}$$

which shows $y(t) \rightarrow -\frac{b_0}{r_0}$ because $e^{R(t)} \rightarrow 0$. \square

Lemma 5.6. *Consider the differential equation $y' = r(t)y + b(t)$ on $(-\infty, t_0]$. If $r(t) = O((-t)^{-1-\delta})$ for some $\delta > 0$, $b(t) \rightarrow b_0 \neq 0$ as $t \rightarrow -\infty$, then*

$$\frac{y(t)}{t} \rightarrow b_0 \quad \text{as } t \rightarrow -\infty.$$

Proof. We let $R(t) := \int_{t_0}^t r(s) ds = A + B(t)$, where $A := \int_{t_0}^{-\infty} r(s) ds$ and $B(t) := \int_{-\infty}^t r(s) ds = O\left((-t)^{-\delta}\right)$. Then

$$\begin{aligned} \frac{y(t)}{t} &= \frac{e^{R(t)}y(t_0)}{t} + \frac{\int_{t_0}^t b(s)e^{R(t)-R(s)} ds}{t} \\ &= \frac{e^{R(t)}y(t_0)}{t} + e^{B(t)} \frac{\int_{t_0}^t b(s)e^{-B(s)} ds}{t}. \end{aligned}$$

It is clear that $\frac{e^{R(t)}y(t_0)}{t} \rightarrow 0$, $e^{B(t)} \rightarrow 1$, and $\int_{t_0}^t b(s)e^{-B(s)} ds \rightarrow -\text{sign}(b_0)\infty$. Applying L'Hospital's rule then gives

$$\lim_{t \rightarrow -\infty} \frac{y(t)}{t} = \lim_{t \rightarrow -\infty} b(t)e^{-B(t)} = b_0. \quad \square$$

Remark 2. The above two lemmas can be easily extended to the case $t \rightarrow \infty$. For example, in Lemma 5.5. we have to require $r_0 < 0$ in this case.

Theorem 5.7. *Let $q = p$. The following conditions are equivalent:*

- a) ϕ is an M -solution.
- b) We have

$$\begin{aligned} \phi(r) &= \frac{1}{(-\ln r^\lambda)^{1/\lambda}} \cdot \frac{1}{r} (1 + o(1)), \\ \phi'(r) &= -\frac{1}{(-\ln r^\lambda)^{1/\lambda}} \cdot \frac{1}{r^2} (1 + o(1)) \quad (r \rightarrow 0). \end{aligned}$$

- c) We have

$$\begin{aligned} u(t) &= -\frac{1}{\lambda} \cdot \frac{1}{t} (1 + o(1)), \\ v(t) - 1 &= \frac{1}{\lambda} \cdot \frac{1}{t} (1 + o(1)) \quad (t \rightarrow -\infty). \end{aligned}$$

- d) $\varphi(t) \rightarrow P_2$ ($t \rightarrow -\infty$).

In this case

$$\lim_{u \rightarrow 0^+} \hat{v}'(u) = \hat{v}'(0) = -1. \quad (5.2.1)$$

Proof. a) \rightarrow d): By Theorem 3.1 on F -solutions, we conclude that $C^-(\varphi)$ is bounded. By Thieme's theorem and the absence of Cases \textcircled{B} and \textcircled{C} , φ must converge to a stationary point of $(EF_{p,q})$ [4, case $n = m + 3$]. The convergence to P_1 and P_3^- is not possible by Lemma 2.4 and Theorem 4.1 on E -solutions. Hence $\varphi(t) \rightarrow P_2$.

d) \rightarrow c): We define $y := \frac{v-1}{u}$. Then

$$\begin{aligned} \dot{y} &= \frac{u\dot{v} - (v-1)\dot{u}}{u^2} = \frac{v(-1+u+v)}{u} - \frac{v-1}{u}(q(t) - u - pv) \\ &= vy + v - q(t)y + v - 1 + pvy \\ &= ((p+1)v - q(t))y + 2v - 1, \end{aligned} \quad (5.2.2)$$

where

$$((p+1)v - q(t)) \rightarrow 1, \quad 2v - 1 \rightarrow 1 \quad (t \rightarrow -\infty).$$

This implies that $y(t) \rightarrow -1$ as $t \rightarrow -\infty$ by Lemma 5.5. Now

$$\begin{aligned} \left(\frac{1}{u}\right)' &= -\frac{\dot{u}}{u^2} = -\frac{u(q(t) - u - pv)}{u^2} = -q(t)\frac{1}{u} + 1 + p\frac{v-1}{u} + p\frac{1}{u} \\ &= (q - q(t))\frac{1}{u} + qy + 1 =: g(t)\frac{1}{u} + qy + 1, \quad \text{where} \\ g(t) &= \lambda\frac{e^{2t}}{1+e^{2t}} = O(e^{2t}), \quad qy + 1 \rightarrow -q + 1 = -\lambda \quad (t \rightarrow -\infty), \end{aligned}$$

which implies that $\frac{1}{t} \cdot \frac{1}{u(t)} \rightarrow -q + 1 = -\lambda$ as $t \rightarrow -\infty$ by Lemma 5.6. Hence

$$\lim_{t \rightarrow -\infty} t(v(t) - 1) = \lim_{t \rightarrow -\infty} t u(t) y(t) = \frac{1}{\lambda}.$$

This shows c).

c) \rightarrow b): c) means

$$u(\ln r) = -\frac{1}{\lambda} \cdot \frac{1}{\ln r} (1 + o(1)) = -\frac{1}{\ln r^\lambda} (1 + o(1)),$$

Hence with (2.2.5)

$$\begin{aligned} \phi(r) &= \left((1+r^2)^{\lambda/2} \cdot r^{-\lambda} \cdot u(\ln r) \cdot v(\ln r) \right)^{\frac{1}{\lambda}} \quad (\lambda = p-1) \\ &= \frac{1}{r} \cdot \left(-\frac{1}{\ln r^\lambda} (1 + o(1)) \right)^{1/\lambda} (1 + o(1)) \\ &= \frac{1}{(-\ln r^\lambda)^{1/\lambda}} \cdot \frac{1}{r} (1 + o(1)), \\ \phi'(r) &= -\frac{1}{r} \phi(r) (1 + o(1)). \end{aligned}$$

This shows b).

b) \rightarrow a) is trivial.

Finally, (5.2.2) means

$$\dot{y} = r(t)y + b(t), \quad \text{with } r(t) = (p+1)v - q(t), \quad b(t) = 2v - 1,$$

where c) implies

$$r(t) = 1 + \frac{\lambda+2}{\lambda t} [1 + o(1)], \quad b(t) = 1 + \frac{2}{\lambda t} [1 + o(1)].$$

Let us define $y_1(t) := t(y(t) + 1)$. Then

$$\begin{aligned} \dot{y}_1 &= \left[r(t) + \frac{1}{t} \right] y_1 + t[b(t) - r(t)], \\ r(t) + \frac{1}{t} &\rightarrow 1, \quad t[b(t) - r(t)] \rightarrow -1, \end{aligned}$$

so that again by Lemma 5.5 we have $y_1 \rightarrow 1$ as $t \rightarrow -\infty$. Hence by c)

$$\begin{aligned} \lim_{u \rightarrow 0^+} \hat{v}'(u) &= \lim_{t \rightarrow -\infty} \frac{v(-1 + u + v)}{u(q(t) - u - pv)} = \lim_{t \rightarrow -\infty} \frac{v\left(\frac{v-1}{u} + 1\right)}{q(t) - p - u - p(v-1)} \\ &= \lim_{t \rightarrow -\infty} \frac{y_1}{t(q(t) - p) - tu - pt(v-1)} \\ &= 1 \cdot \frac{1}{0 + \frac{1}{\lambda} - p\frac{1}{\lambda}} = -1 = -\hat{v}'(0). \end{aligned}$$

□

5.3. The case $q < p$, $q \neq \frac{p+1}{2}$

For $q < p$ there are no solutions φ which converge to P_2 for $t \rightarrow -\infty$. In fact, if $\frac{q}{p} < 1$, then $P_2 \in S_-(t)$ for all t (see Section 2.6). Hence $\dot{u}(t) < 0$ for sufficiently small t , which implies $u(t) \not\rightarrow 0$. On the other hand, there are always φ converging to P_4 . It is easy to see this in the case $\frac{p+1}{2} < q$. In fact, $q < p$ implies $P_4 \in \bar{G}_-$ and $\frac{p+1}{2} < q$ implies $P_3^- \notin \bar{G}_-$, so that all solutions starting in G_- at some $t_0 < T^*$ stay in G_- by Lemma 2.3 and necessarily converge to P_4 by Thieme's theorem. For $a \neq 0$ (2.4.1) does not have 0 as a solution.

The picture is not so clear for $q < \frac{p+1}{2}$, when $\operatorname{Re} \lambda_i < 0$, $i = 1, 2$. Let us convert $(\text{MS}_{p,q})$ into an autonomous system by adding the additional variable $W(t) := e^{2t}$. By (2.2.6),

$$\dot{q}(t) = q - (q-1)W(t) + (q-1)\frac{W^2(t)}{1+W(t)},$$

so that $t \mapsto (u(t), v(t), W(t))$ satisfies the 3-dimensional system

$$\left. \begin{aligned} \dot{U} &= U(q - U - pV) - (q-1)UW + (q-1)\frac{UW^2}{1+W}, \\ \dot{V} &= V(-1 + U + V), \\ \dot{W} &= 2W. \end{aligned} \right\} \quad (\overline{\text{MS}}_p, \lambda)$$

Its linearization at the stationary point $\bar{P}_4 := (u_4, v_4, 0)$ is, with $\bar{U} := U - u_4$, $\bar{V} := V - v_4$,

$$\begin{pmatrix} \dot{\bar{U}} \\ \dot{\bar{V}} \\ \dot{W} \end{pmatrix} = \bar{A} \begin{pmatrix} \bar{U} \\ \bar{V} \\ W \end{pmatrix} + \begin{pmatrix} -\bar{U}^2 - p\bar{U}\bar{V} - (q-1)\bar{U}W \\ \bar{U}\bar{V} - \bar{V}^2 \\ 0 \end{pmatrix} + \begin{pmatrix} (q-1)\left(\bar{U} + u_4\frac{W^2}{1+W}\right) \\ 0 \\ 0 \end{pmatrix}$$

$$\text{with } \bar{A} = \left(\begin{array}{c|c} A & -(q-1)u_4 \\ \hline 0 & 0 \\ \hline 0 & 2 \end{array} \right).$$

Because $\operatorname{Re} \lambda_i < 0$, $i = 1, 2$ and $\lambda_3 = 2$ (see Section 2.4), \bar{P}_4 is a hyperbolic point for $(\overline{\text{MS}}_{p,\lambda})$. In the case $\frac{p+1}{2} < q$, \bar{P}_4 has a 3-dimensional stable manifold for $t \rightarrow -\infty$. If $q < \frac{p+1}{2}$, the stable manifold for $t \rightarrow -\infty$ has dimension one, whose tangent vector is given by an eigenvector to λ_3 , which is

$$\xi_3 = \left(- (q-1) \frac{(2-v_4)(1-v_4)}{(3-v_4)(2-v_4) + p(1-v_4)v_4}, \right. \\ \left. - (q-1) \frac{(1-v_4)v_4}{(3-v_4)(2-v_4) + p(1-v_4)v_4}, 1 \right).$$

Hence there exists a solution $\psi = (U, V, W)$ converging to \bar{P}_4 and for which $W(t_0) > 0$ for some $t_0 \in \mathbb{R}$. It follows that $W(t) = W(t_0) e^{2(t-t_0)}$, and there exists t_1 with $W(t_1) = 1$, so that $W(t_0) \cdot e^{2(t_1-t_0)} = 1$. Then $t \mapsto (U(t+t_1), V(t+t_1), W(t+t_1) = e^{2t})$ is also a solution of $(\overline{\text{MS}}_{p,\lambda})$ converging to \bar{P}_4 , and $\varphi(t) := (U(t+t_1), V(t+t_1))$ is a solution of $(\text{MS}_{p,\lambda})$ converging to P_4 .

The following theorem shows that P_4 takes the role of P_2 for $q < p$, $q \neq \frac{p+1}{2}$.

Theorem 5.8. *Let $q < p$, $q \neq \frac{p+1}{2}$. The following conditions are equivalent:*

- a) ϕ is an M -solution.
- b) We have

$$\phi(r) = c_{p,\lambda} r^{-\frac{\lambda}{p-1}} (1 + o(1)), \\ \phi'(r) = -c_{p,\lambda} \frac{\lambda}{p-1} r^{-\frac{\lambda}{p-1}-1} (1 + o(1)), \quad c_{p,\lambda} := \left(\frac{(p-\lambda-1)\lambda}{(p-1)^2} \right)^{\frac{1}{p-1}}.$$

- c) We have

$$u(t) = \frac{p-\lambda-1}{p-1} (1 + o(1)), \\ v(t) = \frac{\lambda}{p-1} (1 + o(1)) \quad (t \rightarrow -\infty).$$

- d) $\varphi(t) \rightarrow P_4$.

Proof. a) \rightarrow d): Similarly as in a) \rightarrow d) of the proof of Theorem 5.7.

d) \leftrightarrow c) and b) \rightarrow a) are trivial.

c) \rightarrow b): c) implies with (2.2.2)

$$\begin{aligned}\phi(r) &= \left((1+r^2)^{\lambda/2} r^{-\lambda} \frac{p-q}{p-1} (1+o(1)) \frac{q-1}{p-1} (1+o(1)) \right)^{\frac{1}{p-1}} \\ &= c_{p,\lambda} r^{-\frac{\lambda}{p-1}} (1+o(1)), \\ \phi'(r) &= -\frac{1}{r} \frac{q-1}{p-1} (1+o(1)) \cdot \phi(r),\end{aligned}$$

which is b). \square

5.4. The case $q = \frac{p+1}{2}$

We will use the functional

$$H(u, v) = (q - u - qv) u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} = -G(u, v) u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}}$$

on $\mathbb{R}_0^+ \times \mathbb{R}_0^+$. We have $H > 0$ on G_- (the open triangle $P_1 P_2 P_3^-$), $H = 0$ on the boundary ∂G_- of G_- , $H < 0$ on G_+ , and the maximum value is

$$H(P_4) = H\left(\frac{1}{2}, \frac{1}{2}\right) = \frac{\lambda}{2^{2+2/\lambda}} =: c_\lambda.$$

H is constant along the solutions of $(\text{EFS}_{p, \frac{p+1}{2}})$ [4, p. 510, case $n = 3m+5$] in G_- . These are periodic orbits around P_4 whose graph is $H^{-1}(c)$, $c \in (0, c_\lambda)$. $H^{-1}(0)$ is the union of the stationary points P_1, P_2, P_3^- cyclically chained by orbits of $(\text{EFS}_{p, \frac{p+1}{2}})$ [4, p. 509].

To investigate $(MS_{p, \frac{p+1}{2}})$, we need the following two lemmas.

Lemma 5.9. *Let φ be a solution of $(MS_{p, \frac{p+1}{2}})$ on (T_0, T) . Then*

$$\frac{d}{dt} H(\varphi(t)) = q \cdot u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \frac{e^{2t}}{1+e^{2t}} W(u, v), \quad u = u(t), \quad v = v(t). \quad (5.4.1)$$

Consequently, if B is a bounded subset of $\mathbb{R}^+ \times \mathbb{R}^+$,

$$M_B := \sup_{(u,v) \in B} q u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} |W(u, v)|$$

and $\varphi(t) \in B$ on $[t_1, t_2] \subset (T_0, T)$ then

$$|H(\varphi(t_2)) - H(\varphi(t_1))| \leq \frac{1}{2} M_B (e^{2t_2} - e^{2t_1}). \quad (5.4.2)$$

Proof. We have, with $q(t) = q - g(t)$, $g(t) = \lambda \frac{e^{2t}}{1+e^{2t}}$,

$$\begin{aligned}
\frac{d}{dt}H(\varphi(t)) &= (q - u - qv) u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} + (q - u - qv) \frac{1}{\lambda} u^{\frac{1}{\lambda}-1} \dot{u} v^{1+\frac{1}{\lambda}} \\
&\quad + (q - u - qv) u^{\frac{1}{\lambda}} \left(1 + \frac{1}{\lambda}\right) v^{\frac{1}{\lambda}} \cdot \dot{v} \\
&= u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \left(-u(q(t) - u - pv) - qv(-1 + u + v) \right. \\
&\quad \left. + \frac{1}{\lambda} (q - u - qv)(q(t) - u - pv) \right. \\
&\quad \left. + (q - u - qv) \left(1 + \frac{1}{\lambda}\right) (-1 + u + v) \right) \\
&= u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \left((q(t) - u - pv) \left(-u + \frac{1}{\lambda} (q - u - qv) \right) \right. \\
&\quad \left. + (-1 + u + v) \left(-qv + \left(1 + \frac{1}{\lambda}\right) (q - u - qv) \right) \right) \\
&= u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \left((q(t) - u - pv) \left(-\left(1 + \frac{1}{\lambda}\right) u + \frac{q}{\lambda} - \frac{q}{\lambda} v \right) \right. \\
&\quad \left. + (-1 + u + v) \left(\left(1 + \frac{1}{\lambda}\right) q - \left(1 + \frac{1}{\lambda}\right) u \right. \right. \\
&\quad \quad \left. \left. - \left(q + \frac{q^2}{\lambda} \right) v \right) \right) \\
&= u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \left((q - g(t) - u - pv) \frac{q}{\lambda} (1 - u - v) \right. \\
&\quad \left. + (-1 + u + v) \left(\frac{q^2}{\lambda} - \frac{q}{\lambda} u - \left(1 + \frac{q}{\lambda}\right) v \right) \right) \\
&= u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \left(\frac{q}{\lambda} g(t) (-1 + u + v) \right. \\
&\quad \left. + (-1 + u + v) \cdot 0 \right) \text{ (because } p - \lambda - q = 0), \\
&= \frac{q}{\lambda} u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} g(t) W(u, v) = q \cdot u^{\frac{1}{\lambda}} v^{1+\frac{1}{\lambda}} \frac{e^{2t}}{1+e^{2t}} W(u, v).
\end{aligned}$$

The last assertion follows immediately. \square

In the sequel we will apply the name F -, E -, M -solution also to the corresponding φ .

Lemma 5.10. *Let φ be a solution of $(MS_{p, \frac{p+1}{2}})$ on (T_0, T) . Then φ is an F -solution iff there exists t_0 with $\varphi(t_0) \in G_+ \cup G_0$. In this case $\varphi(t) \in G_+$ for $t < t_0$.*

Proof. The necessity of the condition follows from Theorem 3.1.d). Conversely, if $\varphi(t_0) \in G_+$, then $\varphi(t) \in G_+$ for $t \leq t_0$ by Lemma 2.3. Because $G_+ \subset W_+$, we have $\frac{d}{dt}H(\varphi(t)) > 0$ by (5.4.1), which implies

$$H(\varphi(t)) < H(\varphi(t_0)) =: c_0 \text{ for } t < t_0.$$

If φ were not an F -solution, by Theorem 3.1 $L^-(\varphi)$ would be a bounded set lying outside $H^{-1}((c_0, c_\lambda])$, which is an open neighborhood of $G_- \cup G_0$, because $c_0 < 0$. But this contradicts Thieme's theorem. Hence φ is an F -solution. If $\varphi(t_0) \in G_0$, then $\varphi(t) \in G_+$ for $t < t_0$ by Lemma 2.3 and the conclusion follows again. \square

Lemma 5.10 implies that the F -solutions, which are all in G_+ for small t by Theorem 3.1 either stay in G_+ or (see Lemma 9.5) enter G_- and stay in G_- for larger t by Lemma 2.3. It also implies that the E - and M -solutions are exactly those which totally lie in G_- . For the E -solutions φ , we always have $H(\varphi(-\infty)) := \lim_{t \rightarrow -\infty} H(\varphi(t)) = 0$ by Theorem 4.1.

The M -solutions are the object of the following theorem. For $c \in (0, c_\lambda)$, let $\Gamma_c = H^{-1}(c)$; for a set $K \in \mathbb{R}^2, r > 0$ let $B_r(K) := \cup_{z \in K} B_r(z)$ (the open ball about z with radius r), and for any $z \in \mathbb{R}^2 \setminus \{P_4\}$ let $R(z)$ be the ray $\{P_4 + s(z - P_4), s \geq 0\}$. And let $\Omega_c := \{(u, v) : H(u, v) > c\}$, the region inside Γ_c .

Theorem 5.11. Let $q = \frac{p+1}{2}$.

- a) If φ is an M -solution, then φ lies in G_- and $H(\varphi(-\infty)) \in [0, c_\lambda]$ exists. Conversely we have:
b) $\forall c \in (0, c_\lambda), \forall \varepsilon < \frac{1}{2} \min((c, c_\lambda - c)) \exists \tau \in \mathbb{R} \forall t_0 \leq \tau, \forall z \in \Gamma_c \exists z_0 \in B_\varepsilon(z) \cap R(z) :$

$$H(\varphi(-\infty, t_0, z_0)) = c.$$

- c) For $c = c_\lambda$ there exists a unique M -solution φ such that $H(\varphi(-\infty)) = c_\lambda$.

Proof. a) follows from the preceding observations and (5.4.2). As for b), for $(u, v) \in G_-$, let us consider the solution $\varphi(\cdot, t_0)$ of the initial value problem $\varphi(t_0) = (u, v)$ with $t_0 \in \mathbb{R}$. Obviously,

$$\mathbb{R} = A \cup B \cup C$$

is a partition of \mathbb{R} into the subsets

- $A := \{t_0 \in \mathbb{R}; \varphi(\cdot, t_0) \text{ is an } M\text{-solution with } H(\varphi(-\infty, t_0)) > 0\},$
 $B := \{t_0 \in \mathbb{R}; \varphi(\cdot, t_0) \text{ is an } F\text{-solution}\},$
 $C := \{t_0 \in \mathbb{R}; \varphi(\cdot, t_0) \text{ is a solution in } G_- \text{ with } H(\varphi(-\infty, t_0)) = 0\}.$

The set A contains the interval $(-\infty, s_0)$, where s_0 is defined by $M_{G_-} e^{2s_0} = H(u, v)$. In fact, if $t_0 < s_0$, then $M_{G_-} e^{2t_0} < H(u, v)$, and, as long as $\varphi(s) \in G_-$ on $(t, t_0]$, we have with (5.4.2)

$$H(u, v) - H(\varphi(t, t_0)) \leq \frac{1}{2} M_{G_-} (e^{2t_0} - e^{2t}) < \frac{1}{2} H(u, v),$$

which implies

$$H(\varphi(t, t_0)) > \frac{1}{2} H(u, v) > 0, \quad \varphi(t, t_0) \in G_- \text{ for all } t \leq t_0,$$

so that $(-\infty, s_0) \subset A$ and $(-\infty, s_0) \cap B = (-\infty, s_0) \cap C = \emptyset$. Now we vary $(u, v) \in G_-$ and consider $s_0 = s_0(u, v) := \frac{1}{2} \ln \frac{H(u, v)}{M_{G_-}}$ and $A = A(u, v)$.

Let U be an open set in G_- with closure $\bar{U} \subset G_-$ and define

$$\sigma_U := \min_{(u, v) \in \bar{U}} s_0(u, v) > -\infty.$$

Given $\varepsilon > 0$, it follows from (5.4.2) with $t_1 = -\infty$ and $t_2 := t$ that

$$\begin{aligned} \exists s_\varepsilon < \sigma_U, \quad \forall t \leq s_\varepsilon, \quad \forall (u, v) \in U, \quad \forall t_0 \in A(u, v): \\ |H(\varphi(-\infty, t_0, u, v)) - H(\varphi(t, t_0, u, v))| < \varepsilon. \end{aligned} \quad (5.4.3)$$

Now we show that i) the set

$$\Omega := \{(t_0, u, v); (u, v) \in G_-, t_0 \in A(u, v)\}$$

is open in $\mathbb{R} \times \mathbb{R}_+^2$ and ii) the mapping

$$\mathcal{H}: \Omega \rightarrow (0, c_\lambda], \quad (t_0, u, v) \mapsto H(\varphi(-\infty, t_0, u, v))$$

is continuous. For i) let $(t_1, u_1, v_1) \in \Omega$. Because $t_1 < A(u_1, v_1)$, the set $\{\varphi(t, t_1, u_1, v_1), t \leq t_1\}$ has a positive distance from $H^{-1}(0)$ and there exists an open set U with closure $\bar{U} \subset G_-$ such that $\{\varphi(t, t_1, u, v), t \leq t_1\} \subset U$. Now let $\varepsilon > 0$, $s_\varepsilon < \sigma_U$ be chosen as above, $\delta := \sigma_U - s_\varepsilon > 0$, and $s_\varepsilon^* := \min(s_\varepsilon - \delta, t_1 - \delta)$. Then for all $(u, v) \in U$ we have $\sigma_\varepsilon^* \in A(u, v)$ (because $s_\varepsilon^* < s_\varepsilon < \sigma_U \leq \sigma_0(u, v)$) and for all $t \leq s_\varepsilon^* (< s_\varepsilon)$ we have with (5.4.3)

$$|H(\varphi(-\infty, \sigma_\varepsilon^*, u, v)) - H(\varphi(t, \sigma_\varepsilon^*, u, v))| < \varepsilon. \quad (5.4.4)$$

The continuous dependence of the initial conditions implies the existence of a neighborhood of (t_1, u_1, v_1) which has the form $(t_1 - \delta_1, t_1 + \delta_1) \times V$, $\delta_1 < \delta$, $\bar{V} \subset U$, such that for all $t_2 \in (t_1 - \delta_1, t_1 + \delta_1)$ and $(u, v) \in V$ we have $s_\varepsilon^* < t_2$ and

$$\varphi(s_\varepsilon^*, t_2, u, v) \in U.$$

This implies $s_\varepsilon^* \in A(\varphi(s_\varepsilon^*, t_2, u, v))$ as above and $\varphi(\cdot, t_2, u, v) = \varphi(\cdot, \sigma_\varepsilon^*, \varphi(s_\varepsilon^*, t_2, u, v))$ is an M -solution with H -limit > 0 , that is, $t_2 \in A(u, v)$. This proves i). Furthermore, for the same (t_2, u, v) we have by (5.4.4),

$$|H(\varphi(-\infty, t_2, u, v)) - H(\varphi(s_\varepsilon^*, t_2, u, v))| < \varepsilon,$$

in particular,

$$|H(\varphi(-\infty, t_1, u_1, v_1)) - H(\varphi(s_\varepsilon^*, t_1, u_1, v_1))| < \varepsilon,$$

and δ_1 and V can be assumed to be small enough for having

$$|H(\varphi(s_\varepsilon^*, t_2, u, v)) - H(\varphi(s_\varepsilon^*, t_1, u_1, v_1))| < \varepsilon.$$

This shows ii) also.

To prove b), let c and ε be given. Let $B = \overline{\Omega}_{c - \frac{1}{2} \min((c, c_\lambda - c)}}$. It is easy to see that there exists $\tau \leq \sigma_B$ and $\delta < \frac{\varepsilon}{4}$ such that for all $\forall t_0 \leq \tau, t \leq t_0, z, z' \in B$ with $|z - z'| < \delta$ we have

$$|H(\varphi(t, t_0, z_0)) - H(z')| < \frac{\varepsilon}{4}$$

(one uses Lemma 5.9 and the continuity of \mathcal{H}). This implies:

$$\forall z \in B_\delta(\Gamma_{d_i}) \forall t \leq t_0 : \varphi(t, t_0, z_0) \in B_{\frac{\varepsilon}{4}}(\Gamma_{d_i}) \quad i = 1, 2$$

where $d_1 := c - \frac{\varepsilon}{2}$ and $d_2 := c + \frac{\varepsilon}{2}$, in particular we have

$$\begin{aligned} \forall z \in B_\delta(\Gamma_{d_1}) : \mathcal{H}(t_0, z) &\leq d_1 + \frac{\varepsilon}{4} < c, \\ \forall z \in B_\delta(\Gamma_{d_2}) : \mathcal{H}(t_0, z) &\geq d_2 - \frac{\varepsilon}{4} > c. \end{aligned}$$

Now, if $z \in \Gamma_c$ and $\{z_i\} := \Gamma_{d_i} \cap R(z)$, $i = 1, 2$, then $\mathcal{H}(t_0, z_1) < c < \mathcal{H}(t_0, z_2)$, and the continuity of \mathcal{H} implies the existence of some $z_0 \in R(z)$ between z_1 and z_2 (so that $z_0 \in B_\varepsilon(z)$) with $\mathcal{H}(t_0, z_0) = c$.

To prove c), we first derive an a-priori estimate for such a φ . We have $\varphi(t) \in G_-$ for all t , and with (5.4.2)

$$|H(u(t), v(t)) - H(P_4)| \leq \frac{1}{2} M_{G_-} e^{2t}. \quad (5.4.5)$$

By Taylor's formula, for all (u, v) in a neighborhood of P_4 one has

$$\begin{aligned} H(u, v) - H(P_4) &= \frac{p+1}{(p-1)2^{\frac{p+2}{p-1}}} \left(u - \frac{1}{2}\right)^2 + 2 \left(u - \frac{1}{2}\right) \left(v - \frac{1}{2}\right) \\ &\quad + p \left(v - \frac{1}{2}\right)^2 + o\left(\left|u - \frac{1}{2}, v - \frac{1}{2}\right|^2\right), \end{aligned}$$

which implies

$$\left|u - \frac{1}{2}\right|^2 + \left|v - \frac{1}{2}\right|^2 \leq C_1 |H(u, v) - H(P_4)| \quad (5.4.6)$$

for some $C_1 > 0$ in a neighborhood of $(\frac{1}{2}, \frac{1}{2})$. The decomposition (2.4.1) for P_4 is

$$\dot{\psi}(t) = A\psi(t) + g_1(\psi(t)) + g_2(t, \psi(t))$$

with

$$A = \begin{pmatrix} -\frac{1}{2} & -\frac{p}{2} \\ \frac{1}{2} & \frac{1}{2} \end{pmatrix}, \quad \lambda_{1,2} = \pm \frac{1}{2} \sqrt{p-1} \cdot i,$$

$$g_1(z) := \begin{pmatrix} -\bar{u}^2 - p\bar{u}\bar{v} \\ \bar{u}\bar{v} + \bar{v}^2 \end{pmatrix}, \quad g_2(t, z) := \begin{pmatrix} -(q-1) \left(\bar{u} + \frac{1}{2} \right) \frac{e^{2t}}{1+e^{2t}} \\ 0 \end{pmatrix},$$

(see Section 2.4). Hence for small $t, \sigma \in \mathbb{R}$,

$$\psi(t) = e^{(t-\sigma)A} \psi(\sigma) + \int_{\sigma}^t e^{(t-s)A} (g_1(\psi(s)) + g_2(s, \psi(s))) ds. \quad (5.4.7)$$

We want to let $\sigma \rightarrow -\infty$. To this end, we note that it follows from (5.4.5) and (5.4.6) that

$$|\psi(t)| = O(e^t), \quad t \rightarrow -\infty. \quad (5.4.8)$$

Furthermore, $s \mapsto e^{sA}$ is bounded on \mathbb{R} , and (5.4.7) and (5.4.8) imply for $\sigma \rightarrow -\infty$

$$\psi(t) = \int_{-\infty}^t e^{(t-s)A} (g_1(\psi(s)) + g_2(s, \psi(s))) ds. \quad (5.4.9)$$

Now we are going to construct a $\varphi(t) \rightarrow P_4$, that is, a solution $\psi(t) \rightarrow 0$ of (5.4.9). The estimate (5.4.8) suggests to consider the Banach space

$$X_S := \{ \psi \in C(-\infty, S], \quad t \mapsto e^{-t} |\psi(t)| \text{ is bounded on } (-\infty, S] \}$$

equipped with the norm

$$\|\psi\| := \sup_{t \leq S} e^{-t} |\psi(t)|, \quad |z| := |z_1| + |z_2|, \quad z = (z_1, z_2) \in \mathbb{R}^2.$$

If we denote by $F\psi(t)$ the right hand side of (5.4.9), the definitions of g_1, g_2 and (5.4.8) give with constants C_2, C_3, C_4 ,

$$\begin{aligned} \|F\psi\| &\leq \sup_{t \in S} e^{-t} \int_{-\infty}^t C_2 \left(|\psi(s)|^2 + |\psi(s)| + 1 \right) e^{2s} ds \\ &\leq C_2 \sup_{t \in S} e^{-t} \int_{-\infty}^t \left((\|\psi\|^2 + 1) e^{2s} + \|\psi\| e^{3s} \right) ds \\ &\leq C_2 \left((\|\psi\|^2 + 1) e^S + \|\psi\| e^{2S} \right), \\ \|F\psi_1 - F\psi_2\| &\leq \sup_{t \in S} e^{-t} \int_{-\infty}^t C_3 (|\bar{u}_1^2 - \bar{u}_2^2| \\ &\quad + |\bar{u}_1 \bar{v}_1 - \bar{u}_2 \bar{v}_2| + |\bar{v}_1^2 - \bar{v}_2^2| + |\bar{u}_1 - \bar{u}_2| e^{2s}) ds \\ &\leq C_4 \sup_{t \in S} e^{-t} \int_{-\infty}^t (\|\psi_1\| + \|\psi_2\| + e^S) \|\psi_1 - \psi_2\| e^{2s} ds \\ &\leq C_4 (\|\psi_1\| + \|\psi_2\| + e^S) e^S \|\psi_1 - \psi_2\|. \end{aligned}$$

For any $r > 0$, $S \in \mathbb{R}$ which satisfy

$$C_2 (r^2 + 1 + re^S) e^S \leq r, \quad C_4 (2r + e^S) e^S < 1 \quad (5.4.10)$$

F is a contraction in the closed ball in X_S of radius r about the origin. The unique fixed point ψ is a solution of (5.4.9), hence $\varphi := \psi + (\frac{1}{2}, \frac{1}{2})$ solves $(MS_{p, \frac{p+1}{2}})$ and $\psi(t) \rightarrow P_4$. Now let φ_1, φ_2 be two solutions with $\varphi_1(t) \varphi_2(t) \rightarrow P_4$. Then $\psi_i := \varphi_i - (\frac{1}{2}, \frac{1}{2})$ satisfies (5.4.9) and there exist $r_i > 0$, $S_i \in \mathbb{R}$ with $\|\psi_i\| \leq r_i$ in X_{S_i} , $i = 1, 2$. There exists $S \leq \min(S_1, S_2)$ which satisfies (5.4.10) for $r := \max(r_1, r_2)$ and $\|\psi_i\| \leq r$ in X_S . Hence $\psi_1 = \psi_2$ and $\varphi_1 = \varphi_2$. \square

Remark 3. We do not know if there exists an M -solution φ with $L^-(\varphi) = \partial G_-$ (\odot of Thieme's theorem) and if for all $(u, v) \in G_-$ there exists t_0 such that $\varphi(\cdot, t_0, u, v)$ is an E -solution.

6. Solutions with $R < \infty$

In this section, we consider solutions ϕ on (R_0, R) , $0 \leq R_0 < R \leq \infty$ and corresponding φ on (T_0, T) , $-\infty \leq T_0 < T \leq +\infty$, $T_0 = \ln R_0$, $T = \ln R$. The following theorem implies the existence of ϕ with $R < \infty$, that is, $\lim_{r \uparrow R} \phi(r) = 0$ for finite $R > 0$.

Theorem 6.1. *The following conditions are equivalent:*

- a) $R < \infty$.
- b) $T < \infty$.
- c) $C^+(\varphi)$ is unbounded as $t \uparrow T$.
- d) $\exists t_0 \in (T_0, T)$ with $v(t_0) > 1$.
- e) $v(t) \uparrow \infty$ as $t \uparrow T$.

In this case, $u(t) \downarrow 0$ as $t \uparrow T$, and $L^+(\varphi) = \emptyset$ (we write $\varphi \rightarrow \emptyset$).

Proof. a) \leftrightarrow b) follows from the definitions of R and T , and e) \rightarrow d) is clear.

b) \rightarrow c): If $C^+(\varphi)$ were bounded, the closure of $C^+(\varphi)$ would be a compact subset of $\mathbb{R}_0^+ \times \mathbb{R}_0^+$. But because $T < \infty$, it cannot contain a point of the axes, it is a subset of $\mathbb{R}^+ \times \mathbb{R}^+$. This implies that φ can be extended beyond T , which is a contradiction.

c) \rightarrow d): If d) were not true, u is unbounded as $t \rightarrow -\infty$. We have

$$\dot{u} \leq u(q(t) - u) < u(q - u).$$

The solution of the initial-value problem $\dot{z} = z(q - z)$, $z(t_0) = u(t_0) =: u_0$ for some $t_0 \in (T_0, T)$ is

$$z(t) = \frac{q}{1 + ce^{-qt}}, \quad c := \frac{q - u_0}{u_0} e^{qt_0}$$

on the interval $\left(-\frac{1}{q} \ln\left(-\frac{1}{c}\right), \infty\right)$ for $u_0 > q$ and on \mathbb{R} for $u_0 \leq q$, and $u(t) \leq z(t)$ for $t_0 \leq t < T$. Because $z(t) \rightarrow q$, u must be bounded, which is a contradiction.

d) \rightarrow b), e): We have $\dot{v} > v(-1+v)$. The solution z of the initial-value problem $\dot{z} = z(-1+z)$, $z(t_0) = v(t_0) =: v_0$ is

$$z(t) = \frac{1}{1 - ce^{ct}}, \quad c := \frac{v_0 - 1}{v_0} e^{t_0} > 0$$

on the interval $(-\infty, \ln \frac{1}{c})$, and $z(t) < v(t)$ for those $t \geq t_0$ where z and v exist. This implies $T < \infty$ and necessarily $v(t) \uparrow \infty$ as $t \uparrow T$.

Now we show $u(t) \downarrow 0$ as $t \uparrow T$. Without loss of generality we can assume $v_0 > \max\left\{1, \frac{q}{p}\right\}$, which implies $\varphi(t) \in W_+ \cap S_-(t)$ on (t_0, T) . Hence the function $\hat{v}(u) := v(t(u))$ exists for (L, u_0) , $L := \lim_{t \uparrow T} u(t)$. \hat{v} satisfies the differential equation

$$\begin{aligned} \hat{v}'(u) &= \frac{\dot{v}(t(u))}{\dot{u}(t(u))} = \frac{v(t(u))(-1+u+v(t(u)))}{u(q(t(u))-u-v(t(u)))} \\ &= \frac{\hat{v}(u)(-1+u+\hat{v}(u))}{u(q(t(u))-u-\hat{v}(u))}, \quad L \leq u \leq u_0. \end{aligned} \tag{6.1}$$

We consider the function $y(u) := c\left(\frac{u_0}{u}\right)^{\frac{p+1}{p}}$ for some $c > \max\left\{v_0, \frac{(q-1)p+q}{p^2}\right\}$. Then we have for $L \leq u \leq u_0$:

1.

$$y(u) \geq y(u_0) = c \geq v_0 > \max\left\{1, \frac{q}{p}\right\}, \tag{6.2}$$

$$\text{hence} \quad q - u - p \cdot y(u) < 0 \quad \text{and} \quad -1 + u + y(u) > 0, \tag{6.3}$$

2.

$$\begin{aligned} p^2 y(u) &> ((q-1)p+q)\left(\frac{u_0}{u}\right)^{\frac{p+1}{p}} > (q-1)p+q \\ &\implies -qp - q + pu + py(u) + p^2 y(u) > -p + pu + py(u) \\ &\implies -(p+1)(q-u-py(u)) > p(-1+u+y(u)) \\ &\implies -\frac{p+1}{p} < \frac{-1+u+y(u)}{q-u-py(u)} \quad \text{with (6.3)} \\ &\implies y'(u) = -\frac{p+1}{p} \frac{y(u)}{u} < \frac{y(u)(-1+u+y(u))}{u(1-u-y(u))}. \end{aligned} \tag{6.4}$$

It follows from (6.1), (6.4) and (6.2) that

$$\hat{v}(u) < y(u), \quad u \leq u_0,$$

hence $u(t)^{\frac{p+1}{p}} < cu_0^{\frac{p+1}{p}} v(t)^{-1} \rightarrow 0$ ($t \uparrow T$) with c). \square

Corollary 6.2. *The following conditions are equivalent:*

- a) $R = \infty$.
- b) $T = \infty$.
- c) $C^+(\varphi)$ is bounded as $t \rightarrow \infty$.
- d) $v \leq 1$.
- e) φ converges to a fixed point of $(EFS_{p,q})$, that is, either $\varphi(t) \rightarrow P_2$ or $\varphi(t) \rightarrow P_3^+$ for $t \rightarrow \infty$.

Proof. The equivalence a) \leftrightarrow b) \leftrightarrow c) \leftrightarrow d) follows from Theorem 6.1.

c) \rightarrow e) follows from Thieme's theorem and the absence of the situations \textcircled{B} and \textcircled{C} at $(EFS_{p,1})$ (see Section 2.5).

e) \rightarrow c) is obvious. \square

Definition. The solutions ϕ of $(M_{p,\lambda})$ with $R = \infty$ for which $\varphi(t) \rightarrow P_2$, will be called P_2 -solutions; those for which $\varphi(t) \rightarrow P_3^+$, will be called P_3^+ -solutions. In the literature, these solutions are called solutions of "fast" and "slow" decay, respectively.

We will investigate the existence and the properties of the P_2 - and P_3^+ -solutions in the next two sections.

7. P_2 -solutions (solutions of fast decay)

In this section we consider $p, q > 1$ and solutions ϕ on (R_0, ∞) and the corresponding φ on (T_0, ∞) , $T_0 := \ln R_0$.

We convince ourselves that there are solutions φ tending to P_2 as $t \rightarrow +\infty$. As in Section 5.3, we convert $(MS_{p,\lambda})$ into an autonomous system by adding the additional variable $W(t) := e^{-2t}$. By (2.2.7),

$$q(t) = 1 + (q - 1)W(t) - (q - 1)\frac{W^2(t)}{1 + W(t)},$$

so that $t \mapsto (u(t), v(t), W(t))$ satisfies the 3-dimensional system

$$\left. \begin{aligned} \dot{U} &= U(1 - U - pV) + (q - 1)UW - (q - 1)\frac{UW^2}{1 + W}, \\ \dot{V} &= V(-1 + U + V), \\ \dot{W} &= -2W. \end{aligned} \right\} \quad (\overline{MS}_{p,\lambda})$$

Its linearization at the stationary point $\bar{P}_2 := (1, 0, 0)$ with $\bar{V} = V - 1$, is

$$\begin{pmatrix} U \\ \bar{V} \\ W \end{pmatrix}' = \bar{A} \begin{pmatrix} U \\ \bar{V} \\ W \end{pmatrix} + \begin{pmatrix} -U^2 - pU\bar{V} + (q - 1)UW - (q - 1)\frac{W^2}{1 + W} \\ U\bar{V} + \bar{V}^2 \\ 0 \end{pmatrix}$$

with $\bar{A} = \begin{pmatrix} -(p - 1) & 0 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$. The eigenvalues are $\lambda_1 = -(p - 1)$, $\lambda_2 = 1$,

$\lambda_3 = -2$, and the corresponding eigenvectors $\bar{\xi}_1 = (p, -1, 0)$, $\bar{\xi}_2 = (0, 1, 0)$, $\bar{\xi}_3 = (0, 0, 1)$.

The plane through $\bar{\xi}_1$ and $\bar{\xi}_3$ is tangential to the two-dimensional stable manifold in \bar{P}_2 . Consequently, there is a solution $\psi = (U, V, W)$ of $(\overline{\text{MP}}_{p,\lambda})$ tending to \bar{P}_2 for $t \rightarrow +\infty$ such that there exists $t_0 \in \mathbb{R}$ with $\psi(t) \in \mathbb{R}^+ \times \mathbb{R}^+ \times \mathbb{R}^+$ for $t \geq t_0$. We have $W(t) = W(t_0)e^{-2(t-t_0)}$. We define t_1 by $W(t_0)e^{-2(t_1-t_0)} = 1$. Then $t \mapsto \psi(t+t_1)$ is also a solution in $\mathbb{R}^+ \times \mathbb{R}^+ \times \mathbb{R}^+$ for $t \geq t_0 - t_1$ and $W(t+t_1) = W(t_0)e^{-2(t+t_1-t_0)} = e^{-2t}$, so that $\varphi(t) := (U(t+t_1), V(t+t_1)) \in \mathbb{R}^+ \times \mathbb{R}^+$ for $t \geq t_0 - t_1$ solves $(\text{MS}_{p,\lambda})$ and converges to P_2 as $t \rightarrow \infty$.

We first state the following lemma.

Lemma 7.1. *Let $\alpha, \gamma > 0$; $K, L, M \geq 0$ with $\beta := \frac{L}{\alpha} + \frac{M}{\gamma} < 1$. Let $u : [\tau, \infty) \rightarrow [0, \infty)$ be continuous and bounded, and assume for $t \geq \tau$*

$$u(t) \leq Ke^{-\alpha t} + L \int_{\tau}^t e^{-\alpha(t-s)} u(s) ds + M \int_{\tau}^{\infty} e^{-\gamma s} u(t+s) ds.$$

Then for $t \geq \tau$

$$u(t) \leq K^* e^{-(\frac{L}{1-\beta} - \alpha)t}, \quad K^* := \frac{K}{1-\beta} e^{-\frac{L}{1-\beta}\tau}.$$

Sketch of proof: Define $v(t) := \sup_{s \geq t} u(s)$ and $z(t) := e^{\alpha t} v(t)$. Then

$$z(t) \leq \frac{K}{1-\beta} + \frac{L}{1-\beta} \int_{\tau}^t z(s) ds,$$

and Gronwall's inequality proves the assertion. \square

The following theorem characterizes the P_2 -solutions.

Theorem 7.2. *The following conditions are equivalent:*

- a) ϕ is a P_2 -solution, that is, $\varphi(t) \rightarrow P_2$ ($t \rightarrow \infty$).
- b) There exists $c > 0$ such that

$$\phi(r) = \frac{c}{r} (1 + o(1)), \quad \phi'(r) = -\frac{c}{r^2} (1 + o(1)) \quad (r \rightarrow \infty).$$

- c) There exist $c_1 > 0$ and $c_2 \in \mathbb{R}$ such that

$$\begin{aligned} u(t) &= e^{-(p-1)t} c_1 (1 + o(1)), \\ v(t) - 1 &= e^{-(p-1)t} c_2 (1 + o(1)) \quad (t \rightarrow \infty). \end{aligned}$$

In this case, $c_1 = c^{p-1}$, $c_2 = -\frac{1}{p} c^{p-1}$, and $\lim_{u \rightarrow 0} \hat{v}'(u) = \hat{v}'(0) = -\frac{1}{p}$.

Proof. a) \rightarrow b), c): The linearization at P_2 is (see Section 2.5 and (2.2.7))

$$\begin{pmatrix} \dot{u} \\ \dot{\bar{v}} \end{pmatrix} = A \begin{pmatrix} u \\ \bar{v} \end{pmatrix} + \begin{pmatrix} -u^2 - pu\bar{v} \\ u\bar{v} + \bar{v}^2 \end{pmatrix} + \begin{pmatrix} (q-1)u \frac{e^{-2t}}{1+e^{-2t}} \\ 0 \end{pmatrix}$$

or

$$\dot{\psi}(t) = A\psi(t) + g_1(\psi(t)) + g_2(t, \psi(t)) \quad (7.1)$$

where $\psi = (u, \bar{v})$, $\bar{v} = v - 1$, with

$$A := \begin{pmatrix} 1-p & 0 \\ 1 & 1 \end{pmatrix},$$

$$g_1(z) = \begin{pmatrix} -u^2 - pu\bar{v} \\ u\bar{v} + \bar{v}^2 \end{pmatrix}, \quad g_2(t, z) = \begin{pmatrix} (q-1)u \frac{e^{-2t}}{1+e^{-2t}} \\ 0 \end{pmatrix},$$

$z = (u, \bar{v})$.

We remember $\lambda_1 = 1 - p < 0$, $\lambda_2 = 1$, $\xi_1 = (p, -1)$, $\xi_2 = (0, 1)$ from Section 2.5. The stable manifold of the linear part of $(EFS_{p,1})$ is given by $S = \mathbb{R}\xi_1$, the unstable manifold by $U = \mathbb{R}\xi_2$. Let

$$Q_S : \mathbb{R}^2 \rightarrow S,$$

$$Q_U : \mathbb{R}^2 \rightarrow U,$$

be the canonical projections. There exists $a > 0$ so that $|Q_S x| \leq a|x|$, $|Q_U x| \leq a|x|$. Furthermore,

$$Q_S e^{tA} = e^{tA} Q_S = e^{\lambda_1 t} Q_S, \quad Q_U e^{tA} = e^{tA} Q_U = e^{\lambda_2 t} Q_U.$$

The solution ψ of (7.1) satisfies for $\sigma, t \geq \tau > T_0$

$$\psi(t) = e^{(t-\sigma)A} \psi(\sigma) + \int_{\sigma}^t e^{(t-s)A} g_1(\psi(s)) ds + \int_{\sigma}^t e^{(t-s)A} g_2(s, \psi(s)) ds$$

and from a) we have $\psi(t) \rightarrow 0$ ($t \rightarrow \infty$). We apply Q_S for $\sigma = \tau$ and Q_U and get

$$\begin{aligned} \psi(t) &= e^{\lambda_1(t-\tau)} Q_S \psi(\tau) + \int_{\tau}^t e^{\lambda_1(t-s)} Q_S g_1(\psi(s)) ds \\ &\quad + \int_{\tau}^t e^{\lambda_1(t-s)A} Q_S g_2(s, \psi(s)) ds + e^{t-\sigma} Q_U \psi(\sigma) \\ &\quad + \int_{\sigma}^t e^{t-s} Q_U g_1(\psi(s)) ds + \int_{\sigma}^t e^{t-s} Q_U g_2(s, \psi(s)) ds. \end{aligned} \quad (7.2)$$

Let $\delta > 0$ be such that

$$\frac{a(p+q)p}{p-1} \cdot \delta < 1, \quad \frac{a(p+q)\delta}{1 - \frac{a(p+q)}{p-1} \cdot \delta} < \min\left(\frac{p-1}{2}, 2\right) \quad (7.3)$$

and let τ be so large that

$$|\psi(t)| \leq \delta, \quad \frac{e^{-2t}}{1+e^{-2t}} \leq \delta \quad \text{for } t \geq \tau. \quad (7.4)$$

The estimates

$$\begin{aligned} |Q_U \psi(\sigma)| &\leq a \cdot \delta, & |Q_U g_1(\psi(s))| &\leq a(p+1)\delta^2, \\ & & |Q_U g_2(\psi(s))| &\leq a(q-1)e^{-2s} \end{aligned}$$

show that we can let $\sigma \rightarrow +\infty$ in (7.2) and we get

$$\begin{aligned} \psi(t) &= e^{\lambda_1(t-\tau)} Q_S \psi(\tau) + \int_{\tau}^t e^{\lambda_1(t-s)} Q_S g_1(\psi(s)) ds + \int_{\tau}^t e^{\lambda_1(t-s)} Q_S g_2(s, \psi(s)) ds \\ &\quad - \int_t^{\infty} e^{t-s} Q_U g_1(\psi(s)) ds - \int_t^{\infty} e^{t-s} Q_U g_2(s, \psi(s)) ds \end{aligned} \quad (7.5)$$

The next step is to deduce an estimate for ψ . We have for $Q = Q_U$ or $Q = Q_S$, with (7.4),

$$\left. \begin{aligned} |Qg_1(\psi(s))| &\leq a(p+1)\delta|\psi(s)|, \\ |Qg_2(s, \psi(s))| &\leq a(q-1)\frac{e^{-2t}}{1+e^{-2t}}|\psi(s)| \leq a(q-1)\delta|\psi(s)|, \end{aligned} \right\} \quad (7.6)$$

so that (7.5) yields

$$\begin{aligned} |\psi(t)| &\leq e^{\lambda_1 t} e^{-\lambda_1 \tau} |Q_S \psi(\tau)| + a(p+q)\delta \int_{\tau}^t e^{\lambda_1(t-s)} |\psi(s)| ds \\ &\quad + a(p+q)\delta \int_t^{\infty} e^{(t-s)} |\psi(s)| ds \\ &=: K e^{-\alpha t} + L \int_{\tau}^t e^{-\alpha(t-s)} |\psi(s)| ds + M \int_0^{\infty} e^{-\gamma s} |\psi(s+t)| ds \end{aligned}$$

with $\alpha := -\lambda_1 = p-1 > 0$, $\gamma := 1$. We may apply Lemma 7.1, because

$$\beta := \frac{L}{\alpha} + \frac{M}{\gamma} = a(p+q) \left(\frac{1}{p-1} + 1 \right) \delta = \frac{a(p+q)p}{p-1} \cdot \delta < 1, \quad (7.7)$$

and we get

$$|\psi(t)| \leq K^* e^{\left(\frac{L}{1-\beta} - \alpha\right)t}, \quad t \geq \tau. \quad (7.8)$$

Multiplying (7.5) by $e^{-\lambda_1 t}$, we get

$$\begin{aligned} e^{-\lambda_1 t} \psi(t) &= e^{-\lambda_1 \tau} Q_S \psi(\tau) + \int_{\tau}^t e^{-\lambda_1 s} Q_S g_1(\psi(s)) ds \\ &\quad + \int_{\tau}^t e^{-\lambda_1 s} Q_S g_2(s, \psi(s)) ds - e^{-\lambda_1 t} \int_t^{\infty} e^{t-s} Q_U g_1(\psi(s)) ds \\ &\quad - e^{-\lambda_1 t} \int_t^{\infty} e^{t-s} Q_U g_2(s, \psi(s)) ds \\ &=: e^{-\lambda_1 t} Q_S \psi(\tau) + \textcircled{\text{I}} + \textcircled{\text{II}} + \textcircled{\text{III}} + \textcircled{\text{IV}}. \end{aligned}$$

We want to show that the limits of all terms exist. Let us define $\varepsilon := \frac{L}{1-\beta}$. From (7.3) and (7.7) we have

$$\varepsilon < \frac{p-1}{2} \quad \text{and} \quad \varepsilon < 2. \quad (7.9)$$

(I): The estimates

$$e^{-\lambda_1 s} |Q_S g_1(\psi(s))| \leq e^{-\lambda_1 s} a(p+1) K^{*2} e^{2(\varepsilon-\alpha)s}$$

(with (7.8)) and $-\lambda_1 + 2\varepsilon - 2\alpha = p - 1 + 2\varepsilon - 2(p - 1) = 2\varepsilon - (p - 1) < 0$ (with (7.9)) show that the integrand is integrable on $[\tau, \infty)$.

(II): The estimates

$$\begin{aligned} e^{-\lambda_1 s} |Q_S g_2(s, \psi(s))| &\leq e^{-\lambda_1 s} a(q-1) e^{-2s} |\psi(s)| \\ &\leq a(q-1) K^* e^{-\lambda_1 s} e^{-2s} e^{(\varepsilon-\alpha)s} \end{aligned}$$

(with (7.6) and (7.8)) and $-\lambda_1 - 2 + \varepsilon - \alpha = -2 + \varepsilon < 0$ (with (7.9)) show that the integrand is also integrable on $[\tau, \infty)$.

(III):

$$\begin{aligned} e^{-\lambda_1 t} \int_t^\infty e^{t-s} |Q_U g_1(\psi(s))| ds &\leq e^{pt} \int_t^\infty e^{-s} a(p+1) |\psi(s)|^2 ds \\ &\leq e^{pt} \int_t^\infty e^{-s} a(p+1) K^{*2} e^{2(\varepsilon-\alpha)s} ds \\ &< \frac{a(p+1) K^{*2}}{-(-1+2(\varepsilon-\alpha))} e^{pt} e^{(-1+2(\varepsilon-\alpha))t} \end{aligned}$$

(with (7.8)) and $p - 1 + 2\varepsilon - 2\alpha = p - 1 + 2\varepsilon - 2(p - 1) = 2(\varepsilon - \frac{p-1}{2}) < 0$ (with (7.9)) show that (III) $\rightarrow 0$.

(IV):

$$\begin{aligned} e^{-\lambda_1 t} \int_t^\infty e^{t-s} |Q_U g_2(s, \psi(s))| ds &\leq e^{pt} \int_t^\infty e^{-s} a(q-1) e^{-2s} |\psi(s)| ds \\ &\leq a(q-1) K^* e^{pt} \int_t^\infty e^{-s} \cdot e^{-2s} e^{(\varepsilon-\alpha)s} ds \\ &< \frac{a(q-1) K^*}{-(-1+2+\varepsilon-\alpha)} e^{pt} e^{(-1-2+\varepsilon-\alpha)t} \end{aligned}$$

(with (7.8)) and $p - 3 + \varepsilon - \alpha = p - 3 + \varepsilon + 1 - p = \varepsilon - 2 < 0$ (with (7.9)) show that (IV) $\rightarrow 0$.

Hence there exist $c_1, c_2 \in \mathbb{R}$ such that

$$\begin{aligned} e^{(p-1)t} u(t) &\rightarrow c_1, & e^{(p-1)t} (v(t) - 1) &\rightarrow c_2 & (t \rightarrow \infty), \\ r^{p-1} u(\ln r) &\rightarrow c_1, & r^{p-1} (v(\ln r) - 1) &\rightarrow c_2 & (r \rightarrow \infty). \end{aligned}$$

With (2.2.5),

$$\begin{aligned} r^{p-1} \phi^{p-1}(r) &= r^{p-1} \frac{(1+r^2)^{\lambda/2}}{r^\lambda} u(\ln r) v(\ln r) \rightarrow c_1 \geq 0, \\ r\phi(r) &\rightarrow c_1^{\frac{1}{p-1}}, \quad -r^2 \cdot \phi'(r) = r\phi(r) \cdot v(\ln r) \rightarrow c_1^{\frac{1}{p-1}}. \end{aligned}$$

Integration of $(M_{p,\lambda})$ gives

$$-r^2\phi'(r) = -r_0^2\phi'(r_0) + \int_{r_0}^r \frac{s^\lambda}{(1+s^2)^{\lambda/2}} \phi(s)^p ds > -r_0^2\phi'(r_0) > 0,$$

which implies that $c_1 > 0$. Hence a) \rightarrow b), c) with $c := c_1^{\frac{1}{p-1}}$.

b) \rightarrow a): With (2.2.4),

$$\begin{aligned} u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\phi^p(r)}{-\phi'(r)} = \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\frac{c^p}{r^p}(1+o(1))}{-\frac{c}{r^2}(1+o(1))} \\ &= -c^{p-1} \frac{r^\lambda}{(1+r^2)^{\lambda/2}} \frac{r^2}{r \cdot r^p} (1+o(1)) = -\frac{c^{p-1}}{r^{p-1}} (1+o(1)) \rightarrow 0, \\ v(t) &= r \frac{-\phi'(r)}{\phi(r)} = \frac{-r^2\phi'(r)}{r^2\phi(r)} = \frac{c(1+o(1))}{c(1+o(1))} \rightarrow 1. \end{aligned}$$

c) \rightarrow a) is obvious.

We need to show the remaining assertions. For $q(t) < p$, we have $P_2 \in \overline{S_-(t)}$

(see Section 2.6). Hence there exist the inverse function $t = t(u)$ and the function $\hat{v}(u) = v(t(u))$. Furthermore,

$$\begin{aligned} \hat{v}'(u(t)) &= \frac{\dot{v}(t)}{\dot{u}(t)} = \frac{v(t)(-1+u(t)+v(t))}{u(t)(q(t)-u(t)-pv(t))} \\ &= \frac{v(t)}{q(t)-u(t)-pv(t)} \left(\frac{e^{-(p-1)t}(v(t)-1)}{e^{-(p-1)t}u(t)} + 1 \right) \\ &\rightarrow \frac{1}{1-p} \left(\frac{c_2}{c_1} + 1 \right). \end{aligned}$$

This shows that the limit

$$\lim_{u \rightarrow 0} \hat{v}'(u) = \lim_{t \rightarrow \infty} \hat{v}'(u(t)) = \frac{1}{1-p} \left(\frac{c_2}{c_1} + 1 \right) =: L$$

exists, and, by de l'Hospital's rule,

$$\frac{c_2}{c_1} = \lim_{t \rightarrow \infty} \frac{v(t)-1}{u(t)} = L.$$

From this we get the equation

$$L = \frac{1}{1-p} (L+1),$$

which gives $L = -\frac{1}{p}$ and $c_2 = -\frac{1}{p} \cdot c_1 = -\frac{1}{p} c^{p-1}$. \square

8. P_3^+ -solutions (solutions of slow decay)

The general assumption of this section is again the one of Section 7.

The following theorem characterizes the P_3^+ -solutions, from which their existence will follow.

Theorem 8.1. *Let $p, q > 1$. The following conditions are equivalent:*

- a) ϕ is a P_3^+ -solution.
- b) We have

$$\begin{aligned}\phi(r) &= \left(\frac{1}{(p-1)\ln r} \right)^{\frac{1}{p-1}} (1 + o(1)), \\ \phi'(r) &= - \left(\frac{1}{(p-1)\ln r} \right)^{\frac{p}{p-1}} \cdot \frac{1}{r} (1 + o(1)) \quad (r \rightarrow \infty).\end{aligned}$$

- c) We have

$$\begin{aligned}u(t) &= 1 - \frac{p}{p-1} \frac{1}{t} (1 + o(1)), \\ v(t) &= \frac{1}{p-1} \frac{1}{t} (1 + o(1)) \quad (t \rightarrow \infty).\end{aligned}$$

- d) There exists $t_0 > T^*$ with $\varphi(t_0) \in G_-$.
- e) There exists $t_1 > T^*$ with $\varphi(t_1) \in W_-$.

In this case

$$\lim_{u \rightarrow 1} \hat{v}'(u) = \hat{v}'(1) = -\frac{1}{p}.$$

Proof. a) \rightarrow c): Let $y(t) := \frac{1-u(t)}{v(t)}$ and $g(t) := (q-1) \frac{e^{-2t}}{1+e^{-2t}}$. Then

$$\begin{aligned}\dot{y} &= \frac{-v\dot{u} - (1-u)\dot{v}}{v^2} = -\frac{1}{v}u(1+g(t) - u - pv) - y(-1+u+v) \\ &= -uy + pu - \frac{u}{v}g(t) + y - uy + u - 1 \\ &= (-2u+1)y + (p+1)u - 1 - \frac{u}{v}g(t) =: r(t)y + b(t).\end{aligned}$$

By a), $r(t) \rightarrow -1$. To determine the limit of b , we note

$$\left(\frac{1}{v} \right)' = -\frac{\dot{v}}{v^2} = \frac{(1-u-v)}{v} \leq \frac{1}{v}, \text{ hence } \frac{1}{v(t)} \leq Ke^t,$$

for some constant $K > 0$. Hence $b(t) \rightarrow p$. By Lemma 5.5 and Remark 2, $y(t) \rightarrow p$. Therefore $\left(\frac{1}{v} \right)' = y - 1$ implies

$$\begin{aligned}v(t) &= \frac{1}{t \frac{1}{t \frac{1}{v(t_0)}} + \frac{1}{t} \int_{t_0}^t (y(s) - 1) ds} = \frac{1}{t} \frac{1}{p-1} (1 + o(1)), \\ 1 - u(t) &= v(t)y(t) = \frac{1}{t} \frac{p}{p-1} (1 + o(1)),\end{aligned}$$

which is c).

c)→b) follows from (2.2.5).

b) implies a weaker form of c), namely

$$\begin{aligned}
u(t) &= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \frac{\phi^p(r)}{-\phi'(r)} \\
&= \frac{r^{\lambda-1}}{(1+r^2)^{\lambda/2}} \left(\frac{1}{(p-1)\ln r} \right)^{\frac{p}{p-1}} ((p-1)\ln r)^{\frac{p}{p-1}} r(1+o(1)) \\
&= 1+o(1), \\
v(t) &= r \frac{-\phi'(r)}{\phi(r)} = r \left(\frac{1}{(p-1)\ln r} \right)^{\frac{p}{p-1}} \cdot \frac{1}{r} ((p-1)\ln r)^{\frac{1}{p-1}} (1+o(1)) \\
&= \frac{1}{p-1} \frac{1}{t} (1+o(1)),
\end{aligned}$$

which already implies a).

d) → a): By Lemma 2.3, $\varphi(t) \in G_-$ for $t \geq t_0$ and by Thieme's theorem, φ tends to P_2 or P_3^+ . However, any solution tending to P_2 must have slope $\hat{v}'(0) = -\frac{1}{p} > -\frac{2}{p+1}$ (the slope of ∂G_-) by Theorem 7.2, which implies a).

a) → e): follows from the definition of G_- .

e) → d) is trivial because $W_- \subset G_-$.

Finally, let us consider $y_1 := t(y-p)$. We get

$$\dot{y}_1 = \frac{1}{t}y_1 + t(ry+b) = \left(r + \frac{1}{t}\right)y_1 + t(pr+b).$$

Because $r + \frac{1}{t} = (-2u+1) + \frac{1}{t} \rightarrow 1$ and

$$\begin{aligned}
t(pr+b) &= t(-2pu+p+(p+1)u-1 + \frac{pu}{v}g) \\
&= t(p-1)(1-u) + t\frac{pu}{v}g = p(1+o(1)) + O(t^2e^{-2t}) = p(1+o(1)),
\end{aligned}$$

we have $y_1(t) \rightarrow b$ by Lemma 5.5 and Remark 2. Hence

$$\begin{aligned}
\lim_{u \rightarrow 1} \hat{v}(u) &= \lim_{t \rightarrow \infty} \frac{v(-1+u+v)}{u(1+g(t)-u-pv)} = \lim_{t \rightarrow \infty} \frac{\frac{-1+u}{v} + 1}{u \left(\frac{g(t)}{v^2} + \frac{\frac{1-u}{v}-p}{tv} \right)} \\
&= \lim_{t \rightarrow \infty} \frac{-y+1}{u \left(\frac{g(t)}{v^2} + \frac{y_1}{tv} \right)} = \frac{-p+1}{1(0+(p-1)p)} = -\frac{1}{p}. \quad \square
\end{aligned}$$

9. Global behaviour of the solutions and their radius and mass

The aim of this section is to investigate the question which F -, E -, M -solutions have finite or infinite radius and mass. We give numerous initial

conditions in $\mathbb{R}^+ \times \mathbb{R}^+$, for which the corresponding solutions have the desired properties. These initial conditions can be transformed back to conditions for $\phi(r_0), \phi'(r_0)$ with the equations (2.2.5). The total mass of the stellar dynamic model induced by a solution ϕ of $(M_{p,\lambda})$ is given by

$$\begin{aligned} M &= \bar{c} \int_{R_-}^R \frac{r^\lambda}{(1+r^2)^{\lambda/2}} \phi(r)^p dr = -\bar{c} \int_{R_-}^R (r^2 \phi'(r))' dr \\ &= \bar{c} \lim_{r \downarrow R_-} r^2 \phi'(r) + \bar{c} \lim_{r \uparrow R} r^2 (-\phi'(r)) =: L_- + L, \end{aligned}$$

with \bar{c} as in (1.4).

Let us remark immediately that the limit L_- is always finite—this follows for the F -solutions from Section 3, for the E -solutions from Theorem 4.1, and for the M -solutions from Theorems 5.1, 5.7, 5.8 (note that $-\frac{\lambda}{p+1} - 1 + 2 > 0 \Leftrightarrow q < p$) and 5.11 (note that $r^2 \phi' \rightarrow 0$ for $r \rightarrow 0$ by (2.2.5)). Hence $M < \infty \Leftrightarrow L < \infty$.

The following theorem characterizes the solutions ϕ with finite and infinite radius and mass by their images φ in $\mathbb{R}^+ \times \mathbb{R}^+$.

Theorem 9.1. *We have the following three equivalences:*

- a) $R < \infty \Leftrightarrow \varphi \rightarrow \emptyset$. In this case $M < \infty$.
- b) $R = \infty, M < \infty \Leftrightarrow \phi$ is a P_2 -solution, that is, $\varphi \rightarrow P_2$.
- c) $R = \infty, M = \infty \Leftrightarrow \phi$ is a P_3^+ -solution, that is, $\varphi \rightarrow P_3^+$.

Proof. We have $R < \infty \Rightarrow M < \infty \Leftrightarrow L < \infty$. Hence a) follows from Theorem 6.1. Parts b) and c) follow from Theorems 7.2 and 8.1 in connection with Corollary 6.2. \square

9.1. The case $q \leq \frac{p+1}{2}$

Theorem 9.2. *Let $q \leq \frac{p+1}{2}$. If φ is any solution for which $\varphi(t_0) \in G_-$, then $\varphi(t) \rightarrow P_3^+$ ($t \rightarrow \infty$). Consequently, all M - and E -solutions, and all F -solutions for which $\varphi(t_0) \in G_-$ for some t_0 , have $R = \infty, M = \infty$.*

Proof. The first assertion follows from Theorem 8.1. The condition $\varphi(t_0) \in G_-$ is satisfied for M -solutions because $\varphi(t) \rightarrow P_4$ ($t \rightarrow -\infty$) by Theorem 5.8, and for E -solutions by Theorem 4.1 in the case $q < \frac{p+1}{2}$, and by the remarks following Lemma 5.10 in the case $q = \frac{p+1}{2}$. \square

Remark 4. The part on E -solutions is known (see Section 9.4).

9.2. *M-solutions for $q > \frac{p+1}{2}$*

We are going to show that in this case the three possibilities a), b), c) of Theorem 9.1 all occur. To this end, it is useful to remark that for any solution we have for $q > \frac{p+1}{2}$:

(i) If there exists $t_0 < T^*$ with $\varphi(t_0) \in G_- \cup G_0$, then φ corresponds to an *M-solution*.

(ii) If there exists $t_1 \in \mathbb{R}$ with $v(t_1) > 1$, then $\varphi \rightarrow \emptyset$ ($t \rightarrow \infty$).

(iii) If there exists $t_2 > T^*$ with $\varphi(t_2) \in G_-$, then $\varphi \rightarrow P_3^+$ ($t \rightarrow \infty$).

In fact, (i) implies with Lemma 2.3 that $\varphi(t_0) \in G_-$ for all $t \leq t_0$. Because $L^-(\varphi)$ is bounded, by Thieme's theorem, $\varphi \rightarrow P_i$ or P_3^- ($t \rightarrow -\infty$) where $i = 2$ in the case $q \geq p$ and $i = 4$ in the case $\frac{p+1}{2} < q < p$. Convergence to P_3^- is not possible because $q > \frac{p+1}{2}$, and φ is an *M-solution* by Theorems 5.1, 5.7 or 5.8. (ii) follows from Theorem 6.1 and (iii) from Theorem 8.1.

Furthermore, we need the following lemma.

Lemma 9.3. *Let $q > \frac{p+1}{2}$. Then there exists $\varepsilon^* > 0$, $T_1 < T^*$ such that for all $(u_0, v_0) \in G_0$ with $|(u_0, v_0) - P_2| < \varepsilon$, $0 < \varepsilon < \varepsilon^*$ and for all $t_0 \leq T_1$*

there exists $t_1 > t_0$ such that $v(t_1, t_0, u_0, v_0) > 1$.

Proof. For $0 < \varepsilon < 1$ we have $(u_0, v_0) = (\varepsilon \frac{p+1}{2}, 1 - \varepsilon) \in G_0$. We let

$$\begin{aligned} x(t) &:= \frac{1}{\varepsilon} u(t, t_0, \varepsilon \frac{p+1}{2}, 1 - \varepsilon) \\ y(t) &:= \frac{1}{\varepsilon} \left[1 - v(t, t_0, \varepsilon \frac{p+1}{2}, 1 - \varepsilon) \right], \end{aligned}$$

and we are going to specify t_1, t_0 . The functions x, y satisfy

$$\begin{aligned} \dot{x} &= x [q(t) - \varepsilon x - p(1 - \varepsilon y)], \\ \dot{y} &= -(1 - \varepsilon y)(x - y). \end{aligned}$$

We consider the limit case $\varepsilon \rightarrow 0$ and arrive at

$$\begin{aligned} \dot{x} &= x [q(t) - p] = x(q - p - g(t)), \\ \dot{y} &= -x + y, \text{ where } g(t) = \lambda \frac{e^{2t}}{1 + e^{2t}}, \end{aligned}$$

so that for $t \geq t_0$

$$\begin{aligned} x(t) &= \frac{p+1}{2} e^{(q-p)(t-t_0)} \left(\frac{1+e^{2t_0}}{1+e^{2t}} \right)^{\frac{\lambda}{2}}, \\ y(t) &= e^{t-t_0} \left(1 - \int_{t_0}^t x(s) e^{-(s-t_0)} ds \right) \\ &= e^{t-t_0} \left(1 - \int_0^{t-t_0} x(s+t_0) e^{-s} ds \right) \\ &= e^{t-t_0} \left(1 - \frac{p+1}{2} \int_0^{t-t_0} e^{(q-p-1)s} \left(\frac{1+e^{2t_0}}{1+e^{2s}e^{2t_0}} \right)^{\frac{\lambda}{2}} ds \right). \end{aligned}$$

If we fix $\tau := t - t_0 > 1$ and let $t_0 \rightarrow -\infty$, the above integral converges to

$$\int_0^\tau e^{(q-p-1)s} ds.$$

Hence there exists $T_1 < T^*$ such that for $t_0 \leq T_1$

$$\int_0^\tau e^{(q-p-1)s} \left(\frac{1+e^{2t_0}}{1+e^{2s}e^{2t_0}} \right)^{\frac{\lambda}{2}} ds > \int_0^{\tau-1} e^{(q-p-1)s} ds.$$

On the other hand we have

$$\frac{p+1}{2} \int_0^\infty e^{(q-p-1)s} ds = \begin{cases} \infty & \text{for } q \geq p+1, \\ \frac{1}{p+1-q} > 1 & \text{for } \frac{p+1}{2} < q < p+1. \end{cases}$$

Therefore we may assume that τ had been chosen so large that

$$\frac{p+1}{2} \int_0^{\tau-1} e^{(q-p-1)s} ds > 1.$$

Then for $t_1 := \tau + t_0$ we have $y(t_1) < 0$ and, by continuity, $v(t_1) > 1$ for all sufficiently small $\varepsilon > 0$. \square

Theorem 9.4. Let $q > \frac{p+1}{2}$.

(a) If $(u_0, v_0) \in G_0$ is close to P_2 and $t_0 < T_1$ as in Lemma 9.3, then the M -solution $\varphi(\cdot, t_0, u_0, v_0)$ has $R < \infty$ and $M < \infty$.

(b) If $(u_1, v_1) \in G_- \cup G_0$, then the M -solution $\varphi(\cdot, T^*, u_1, v_1)$ has $R = \infty$ and $M = \infty$.

(c) For $t_0, (u_0, v_0)$ from (a) and (u_1, v_1) from (b), there exists some $s \in (0, 1)$ such that for $(t_s, u_s, v_s) = (1-s)(t_0, u_0, v_0) + s(T^*, u_1, v_1)$, $\varphi(\cdot, t_s, u_s, v_s)$ is an M -solution with $R = \infty$ and $M < \infty$.

Proof. We apply Theorem 9.1. (a) follows from the preceding remarks (i) and (ii).

(b): Lemma 2.3 shows that $\varphi(t, T^*, u_1, v_1) \in G_-$ for $t \neq T^*$, and remarks (i) and (iii) apply.

(c): We let $Q(s) := (t_s, u_s, v_s) = (1-s)(t_0, u_0, v_0) + s(T^*, u_1, v_1) : s \in [0, 1] \mapsto [t_0, T^*] \times G_- \cup G_0$ and define the sets

$$\begin{aligned} A &:= \{s \in [0, 1]; \varphi(\cdot, Q(s)) \text{ satisfies } (i) \text{ and } (ii)\}, \\ B &:= \{s \in [0, 1]; \varphi(\cdot, Q(s)) \text{ satisfies } (i) \text{ and } (iii)\}. \end{aligned}$$

We have $0 \in A, 1 \in B$. Because φ depends continuously on the initial data, A and B are disjoint, non-empty and relatively open in $[0, 1]$. Because $[0, 1]$ is connected, there is some $s_0 \in (0, 1)$ which is not in $A \cup B$. On the other hand, $Q(s_0) \in (t_0, T^*) \times G_- \cup G_0$ and therefore $\varphi(\cdot, Q(s_0))$ satisfies (i) but not (ii) and (iii), and hence necessarily tends to P_2 . This proves (c). \square

9.3. F -solutions

Our first result in this section is that for all $p, q > 1$ the three possibilities a), b), c) of Theorem 9.1 all occur. We need the following lemma.

Lemma 9.5. *Let $p, q > 1$. There exists $\eta > 0$ such that for all $v_0 \in (0, \eta)$ and $t_0 > T^*$ there exists $t_1 > t_0$ such that $\varphi(t_1, t_0, q, v_0) \in G_-$.*

Proof. If $q < \frac{p+1}{2}$, $(q, v_0) \in G_-$ for small v_0 and the assertion is obvious for all t_0 . Therefore we may assume $q \geq \frac{p+1}{2}$. We define $q_0 := q(t_0) < \frac{p+1}{2}$. We consider (EF_{p, q_0}) and the F -solution $\bar{\varphi}_2$ with $\bar{\varphi}_2(t) \rightarrow P_2$ (see [4, Theorem 5.7, case $n > 3m+5$]). We may assume $\bar{u}_2(t_0) = q$ ((EF_{p, q_0}) is autonomous) and define $\bar{v}_2(t_0) =: \eta$. Let $v_0 < v_1 < \eta$. From [4, Theorem 5.7] we deduce that the solution $\bar{\varphi}(\cdot, t_0, q, v_1)$ of (EF_{p, q_0}) stays below $\bar{\varphi}_2$ and there exist $\bar{t}_2 > \bar{t}_1 > t_0$ such that

$$\bar{\varphi}(\bar{t}_1) \in G_0 \text{ and } \bar{\varphi}(t) \in S(p, q_0) \text{ on } [t_0, \bar{t}_2].$$

Hence $\hat{v} = \hat{v}(u)$ exists on $[\bar{u}_2(\bar{t}_2), q]$ with graph $\hat{v} \subset S(p, q_0)_-$ and satisfies

$$\hat{v}'(u) = \frac{\hat{v}(u)(-1 + u + \hat{v}(u))}{u(q_0 - u - p\hat{v}(u))} =: \bar{k}(u, \hat{v}(u)).$$

Now consider the F -solution $\varphi(\cdot, t_0, q_0, v_0)$ of $\text{MS}_{p, \lambda}$. We want to show: there exists $t_1 > t_0$ with $\varphi(t_1) \in G_0$ (the conclusion then follows from Lemma 2.3). We may assume that \hat{v} also exists on $[\bar{u}_2(\bar{t}_2), q]$ and graph $\hat{v} \subset S(p, q_0)_-$, because otherwise φ would have crossed $S(p, q_0)_0$ and hence G_0 sometime before (see Section 2.6) and we are done. \hat{v} satisfies

$$\hat{v}'(u) = \frac{\hat{v}(u)(-1 + u + \hat{v}(u))}{u(q(t(u)) - u - p\hat{v}(u))} =: k(u, \hat{v}(u)).$$

Because $t(u) > t_0$ for $u < q_0$, we have $q(t(u)) < q(t_0) = q_0$ and

$$k(u, v) > \bar{k}(u, v) \text{ in } S(p, q_0)_-.$$

Hence $\hat{v}(u) < \bar{v}(u)$ on $[\bar{u}_2(\bar{t}_2), q]$, and because $\bar{\varphi}$ crosses G_0 , φ crosses G_0 as well. \square

Theorem 9.6. *Let $q, p > 1$. Then:*

(a) *For (u_0, v_0) with $u_0 \geq q, v_0 \geq 1$ and all t_0 the F -solution $\varphi(\cdot, t_0, u_0, v_0)$ tends to \emptyset and has $R < \infty$ and $M < \infty$.*

(b) *For $v_0 < \eta$ as in Lemma 9.5 and for $t_0 > T^*$ the F -solution $\varphi(\cdot, t_0, q, v_0)$ tends to P_3^+ and has $R = \infty$ and $M = \infty$.*

(c) *There exists $v_0 \in [\eta, 1]$ such that for $t_0 > T^*$ the F -solution $\varphi(\cdot, t_0, q, v_0)$ tends to P_2 and has $R < \infty$ and $M = \infty$.*

Proof. (a) and (b) follow from Theorem 3.1, Lemma 9.5 and Theorem 8.1.

As for (c), the sets

$$A := \{v_0 \in (0, 1] : \exists t > t_0 \text{ with } v(t, t_0, q, v_0) > 1\},$$

$$B := \{v_0 \in (0, 1] : \exists t > t_0 \text{ with } \varphi(t, t_0, q, v_0) \in G_-\}$$

are disjoint, non-empty and relatively open in $(0, 1]$. Because $(0, 1]$ is connected, the complement of $A \cup B$ in $(0, 1]$ is non-empty and contains v_0 for which $\varphi(\cdot, t_0, q, v_0)$ necessarily tends to P_2 . \square

Part b) of the following theorem is inspired by the work of W.-M. Ni and S. Yotsutani [24]. For $R_0 > 0$ and $\alpha > 0$ we denote by $\phi(\cdot, R_0, \alpha)$ the F -solution with $\phi(R_0) = \alpha$ and $\phi'(R_0) = 0$.

Theorem 9.7. *For every $R_0 > 0$, there exist constants $\alpha_0 = \alpha_0(R_0, p, q)$ and $\alpha_1 = \alpha_1(R_0, p, q)$ such that*

a) *$\phi(\cdot, R_0, \alpha)$ has $R < \infty$ for $\alpha \geq \alpha_0$.*

b) *$\phi(\cdot, R_0, \alpha)$ has $R = \infty$ for $\alpha \leq \alpha_1$.*

Proof. a) The function $z(t) := \phi(\frac{1}{t}, R_0, \alpha)$ defined on $(\frac{1}{R}, t_0)$, $t_0 := \frac{1}{R_0}$ is a solution of

$$z''(t) + f_\lambda(t)z^p(t) = 0, \text{ where } f_\lambda(t) := \frac{1}{t^2(1+t^2)^{\frac{\lambda}{2}}}. \quad (9.1)$$

It is sufficient to prove: there exists $t_1 \in (\frac{1}{R}, t_0]$ such that

$$z'(t_1) \geq \frac{z(t_1)}{t_1}, \quad (9.2)$$

i.e., the tangent line passing through $(t_1, z(t_1))$ has a zero in $[0, t_1)$. Because $z'' < 0$, this implies that z has a zero in $(0, t_0)$ and $R < \infty$ follows.

If (9.2) were not true, we would have

$$\frac{z'(t)}{z(t)} < \frac{1}{t} \text{ and hence } \frac{z(t)}{t} > \frac{\alpha}{t_0} \text{ in } (\frac{1}{R}, t_0) \quad (9.3)$$

by integration. Integrating (9.1) once, we obtain that

$$\begin{aligned} z'(t) &= \int_t^{t_0} \frac{1}{s^2(1+s^2)^{\frac{\lambda}{2}}} z^p(s) ds \\ &\geq \frac{\alpha^p}{t_0^p} \int_t^{t_0} s^p f_\lambda(s) ds. \end{aligned}$$

Now there exists $t_1 \in (0, t_0)$ such that $\alpha_0 := \left[\inf_{t \in (0, t_0)} \frac{t_0^p}{t \int_t^{t_0} s^p f_\lambda(s) ds} \right]^{\frac{1}{p-1}} = \left[\frac{t_0^p}{t_1 \int_{t_1}^{t_0} s^p f_\lambda(s) ds} \right]^{\frac{1}{p-1}}$. Then, if $\alpha > \alpha_0$, we would have $z'(t_1) > \frac{z(t_1)}{t_1}$, which is a contradiction. This proves a).

b) We assume $R < \infty$, i.e. $\phi(R) = 0$. By the Pohozaev-type identity [24, Theorem 1, p5.] we have for $n = 3$ that

$$\begin{aligned} \frac{1}{2} R^3 \phi'^2(R) &= \int_{R_0}^R \left[\frac{5-p}{2(p+1)} K(r) + \frac{1}{p+1} r K'(r) \right] \phi^{p+1}(r) r^2 dr \\ &\quad + \frac{R_0^3 K(R_0)}{p+1} \alpha^{p+1} \\ &= \alpha^{p+1} \left\{ \frac{R_0^3 K(R_0)}{p+1} - \int_{R_0}^R G(r) \left[\frac{\phi(r)}{\alpha} \right]^{p+1} dr \right\}, \end{aligned} \quad (9.4)$$

where

$$G(r) = \frac{1}{2(p+1)} \left[p-1 + \frac{p-2\lambda-1}{r^2} \right] \frac{r^{\lambda+2}}{(1+r^2)^{\frac{\lambda}{2}+1}},$$

so that

$$G(r) \rightarrow \frac{p-1}{2(p+1)} =: G_\infty > 0 \text{ as } r \rightarrow \infty. \quad (9.5)$$

Let $R_1 \geq R_0$ such that $G(r) \geq \frac{G_\infty}{2}$ if $r \geq R_1$ and let $M := \sup_{r \geq 0} |G(r)| < \infty$. Since $\phi(r) \leq \alpha$, we have

$$\begin{aligned} r^2 \phi'(r) &= - \int_{R_0}^r s^2 K(s) \phi^p(s) ds \geq -\alpha^p \int_{R_0}^r s^2 K(s) ds, \\ \phi'(r) &\geq -\frac{\alpha^p}{r^2} \int_{R_0}^r s^2 K(s) ds \end{aligned}$$

and hence

$$\phi(r) \geq \alpha - \alpha^p \int_{R_0}^r \frac{1}{t^2} \int_{R_0}^t s^2 K(s) ds dt$$

or

$$\frac{\phi(r)}{\alpha} \geq 1 - \alpha^{p-1} g(r) \text{ for } R_0 < r < R, \quad (9.6)$$

where $g(r) := \int_{R_0}^r \frac{1}{t^2} \int_{R_0}^t s^2 K(s) ds dt$. We now fix

$$R_2 := R_1 + \frac{2^{p+2}}{G_\infty} \left[\frac{R_0^3 K(R_0)}{p+1} + M(R_1 - R_0) \right].$$

Since $g \in C([R_0, \infty))$ there exists $\alpha_1 > 0$ such that for all $\alpha \in (0, \alpha_1]$

$$1 - \alpha^{p-1} g(r) \geq \frac{1}{2}, \text{ on } [R_0, R_2], \quad (9.7)$$

which in particular implies that $R > R_2$, $\alpha \in (0, \alpha_1]$. From (9.4)-(9.7) we obtain for $\alpha \leq \alpha_1$

$$\begin{aligned} \frac{1}{2}R^3\phi'^2(R) &< \alpha^{p+1} \left\{ \frac{R_0^3 K(R_0)}{p+1} + \int_{R_0}^{R_1} M [1]^{p+1} dr - \int_{R_1}^{R_2} \frac{G_\infty}{2} \left[\frac{1}{2}\right]^{p+1} dr \right\} \\ &= \alpha^{p+1} \left\{ \frac{R_0^3 K(R_0)}{p+1} + M(R_1 - R_0) - \frac{G_\infty}{2^{p+2}}(R_2 - R_1) \right\} \\ &= 0 \end{aligned}$$

by the choice of R_2 . This is a contradiction and $R = \infty$ follows. \square

9.4. Known result on E -solutions

For the sake of completeness, we collect the known results on the E -solutions which are relevant in our context. For $\alpha > 0$ we denote by $\phi(\cdot, \alpha)$ the solution ϕ of $(M_{p,\lambda})$ with $\phi(0) = \alpha$.

Theorem 9.8. *a) For $q > \frac{p+1}{2}$ there exists some $\alpha_0 \in (0, \infty)$, such that*

- 1.) *for $\alpha > \alpha_0$, $\phi(\cdot, \alpha)$ has $R < \infty$ and $M < \infty$,*
- 2.) *$\phi(\cdot, \alpha_0)$ has $R = \infty$ and $M < \infty$,*
- 3.) *for $\alpha \in (0, \alpha_0)$ $\phi(\cdot, \alpha)$ has $R = \infty$ and $M = \infty$.*

b) For $q \leq \frac{p+1}{2}$, $\phi(\cdot, \alpha)$ has $R = \infty$, $M = \infty$ for all $\alpha > 0$.

A number of authors contributed to this theorem in connection with a study of conjectures of Matukuma's in the case $\lambda = 2$ [24, p. 2]. For example, W.-M. Ni and S. Yotsutani proved a) 1.) for large $\alpha > 0$, 3.) for all $p, q > 1$ for small α , and that for all $\alpha > 0$, $\phi(\cdot, \alpha)$ has $R = \infty$ if $q \leq \frac{p+1}{2}$ and $M = \infty$ if $q < \frac{p+1}{2}$ [24, Theorems 2, 5 and 6]; see Y. Li and W.-M. Ni for results related to a) 3.) for $\lambda \geq 2$ [16, Theorem 3.20].

The existence of E -solutions with $R = \infty$ and $M < \infty$ for $q > \frac{p+1}{2}$ was proven by Y. Li and W.-M. Ni [17, Theorem 2.1]; their uniqueness by E. Yanagida [36]. Y. Li and W.-M. Ni also proved the precise asymptotics of the E -solutions at $r = \infty$ [16, Theorem 2.41], while Y. Li presented the asymptotic expansions of E -solutions with $R = \infty$ and $M = \infty$ for all $p, q > 1$ in [14]. The result in its above final form was given by N. Kawano, E. Yanagida and S. Yotsutani [12, Theorem 5.2].

Remark 5. We do not know whether separation results analogous to Theorem 9.8 hold for F - and M -solutions.

Appendix

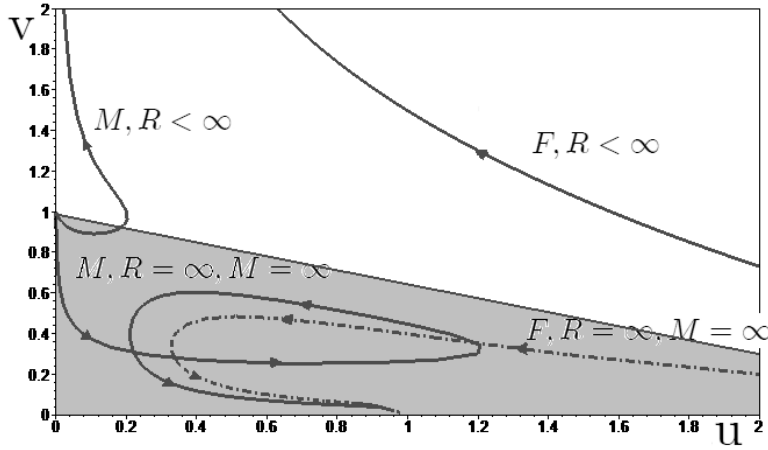


Fig. 1. Case $q > p$ ($q = 3, p = 2$). Illustration of Lemma 2.3 and Theorems 3.1, 9.6 (a), (b), 5.1 and 8.1. Note that the shaded region is the G^- (definition before Lemma 2.3) for Figure 1-4.

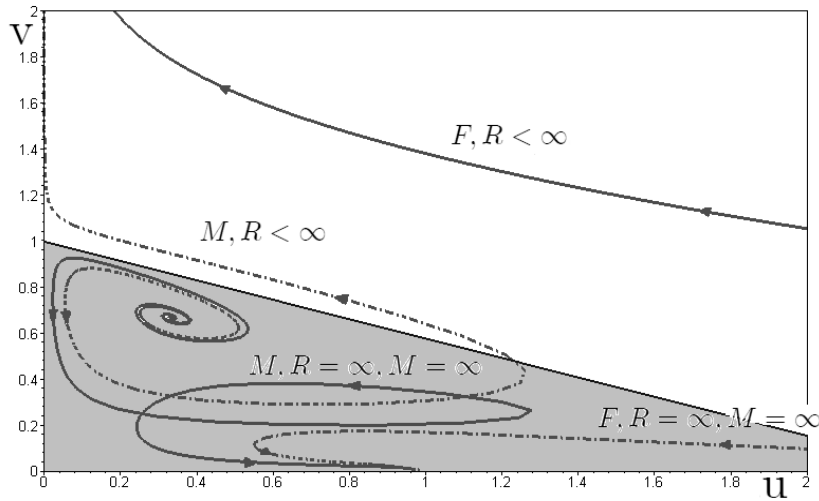


Fig. 2. Case $\frac{p+1}{2} < q < p$ ($q = 3, p = 4$). Illustration of Lemma 2.3 and Theorems 3.1, 9.6 (a), (b), 5.8 and 9.4 (a), (b).

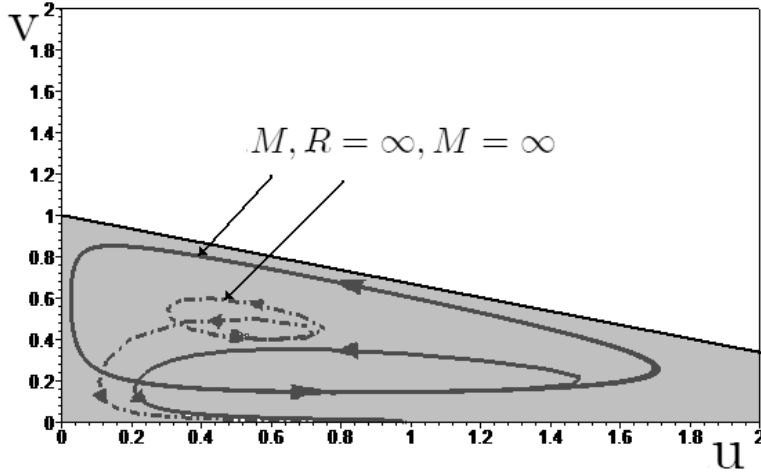


Fig. 3. Case $q = \frac{p+1}{2}$ ($q = 3, p = 5$). Illustration of Lemma 2.3 and Theorem 5.11.

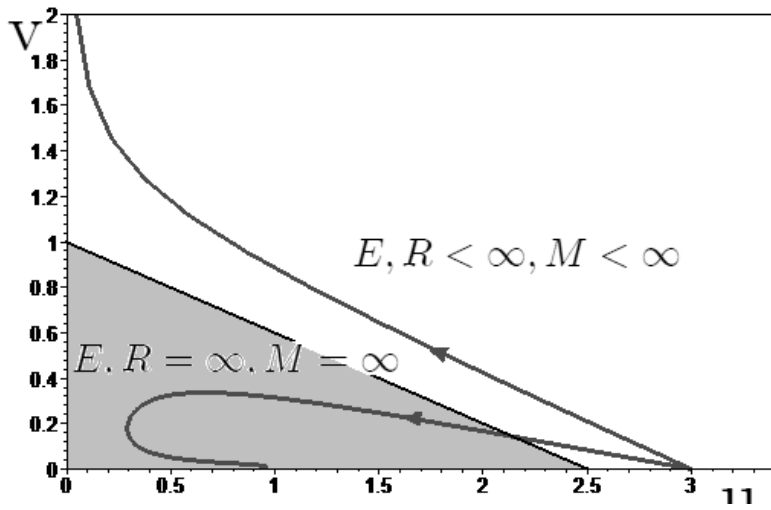


Fig. 4. Case $q > \frac{p+1}{2}$ ($q = 3, p = 4$). Illustration of Lemma 2.3 and Theorems 4.1, 8.1, and 9.8.

Acknowledgements. The authors are grateful to the Mathematisches Forschungs-Institut, Oberwolfach, Germany, and its Director, Prof. H.M. Greuel, for the grant of a stay of two weeks in Oberwolfach in November, 2006 within the program “Research in Pairs”. The first author wants to acknowledge an encouraging correspondence with Prof. Zhien Ma, Xi’an Jiatong University, Xi’an, China. His former student Wolfgang Prenninger did the first steps in extending the paper [4] to the Matukuma equation. He wishes to thank the PDE research group at the Hunan Normal University, Changsha, China for their hospitality during his visit in July/August 2007 and the Deutsche Forschungsgemeinschaft for the travel grant (GZ: BA747/16-1). The second author was supported in part by University of Iowa International Programs, University of Iowa Office of the Vice President for Research, Stanley-UI Support Organization and Hunan Normal University Xiao-Xiang Fund and is grateful for the support.

References

1. Batt, J.: Steady state solutions of the relativistic Vlasov-Poisson system. *Proceedings of the Fifth Marcel Grossmann Meeting on General Relativity, Part B* (Perth, 1988), World Sci. Publ., Teaneck, NJ, 1235–1247 (1989)
2. Batt, J., Faltenbacher, W., Horst, E.: Stationary spherically symmetric models in stellar dynamics. *Arch. Rational Mech. Anal.* **93**, 159–183 (1986)
3. Batt, J., Morrison, P. J., Rein, G.: Linear stability of stationary solutions of the Vlasov-Poisson system in three dimensions. *Arch. Rational Mech. Anal.* **130**, no. 2, 163–182 (1995)
4. Batt, J., Pfaffelmoser, K.: On the radius continuity of the models of polytropic gas spheres which correspond to the positive solutions of the generalized Emden-Fowler equation. *Math. Methods Appl. Sci.* **10**, 499–516 (1988)
5. Bellman, R.: *Stability Theory of Differential Equations*. McGraw-Hill, New York-Toronto-London, 1953
6. Camm, G.L.: Self-gravitating star system II. *Monthly Notices Royal Astronomical Society* **112**, 155–176 (1952)
7. Emden, R.: *Gaskugeln*. Leipzig und Berlin: Teubner 1907
8. Guo, Y.: Variational method for stable polytropic galaxies. *Arch. Rational Mech. Anal.* **150**, 209–224 (1999)
9. Guo, Y., Rein, G.: Existence and stability of Camm type steady states in galactic dynamics. *Indiana Univ. Math. J.* **48**, 1237–1255 (1999)
10. Guo, Y., Rein, G.: Stable steady states in stellar dynamics. *Arch. Rational Mech. Anal.* **147**, 225–243 (1999)
11. Heinzle, J.M., Rendall, A.D., Uggla, C.: Theory of Newtonian self-gravitating stationary spherically symmetric systems. *Math. Proc. Cambridge Philos. Soc.* **140**, 177–192 (2006)
12. Kawano, N., Yanagida, E., Yotsutani, S.: Structure theorems for positive radial solutions to $\Delta u + K(|x|)u^p = 0$ in \mathbf{R}^n . *Funkcial. Ekvac.* **36**, no. 3, 557–579 (1993)
13. Kwong, M.K., Li, Y.: Uniqueness of radial solutions of semilinear elliptic equations. *Trans. Amer. Math. Soc.* **333**, 339–364 (1992)
14. Li, Y.: Asymptotic behavior of positive solutions of equation $\Delta u + K(x)u^p = 0$ in \mathbf{R}^n . *J. Differential Equations* **95**, 304–330 (1992)
15. Li, Y.: On the positive solutions of the Matukuma equation. *Duke Math. J.* **70**, 575–589 (1993)
16. Li, Y., Ni, W.-M.: On conformal scalar curvature equations in \mathbf{R}^n . *Duke Math. J.* **57**, 895–924 (1988)

17. Li, Y., Ni, W.-M.: On the existence and symmetry properties of finite total mass solutions of the Matukuma equation, the Eddington equation and their generalizations. *Arch. Rational Mech. Anal.* **108**, 175–194 (1989)
18. Li, Y., Ni, W.-M.: On the asymptotic behavior and radial symmetry of positive solutions of semilinear elliptic equations in \mathbb{R}^n , Part I. Asymptotic Behavior, *Arch. Rational Mech. Anal.* **118**, 195–222 (1992)
19. Li, Y., Ni, W.-M.: On the asymptotic behavior and radial symmetry of positive solutions of semilinear elliptic equations in \mathbb{R}^n , Part II. Radial Symmetry, *Arch. Rational Mech. Anal.* **118**, 223–244 (1992)
20. Li, Y., Santanilla, J.: Existence and nonexistence of positive singular solutions for semilinear elliptic problems with applications. *Differential Integral Equations* **8**, 1369–1383 (1995)
21. Liu, Y., Li, Y., Deng, Y.B.: Separation property of solutions for a semilinear elliptic equation. *J. Differential Equations* **20**, 381–406 (2000)
22. Matukuma, T.: Sur la dynamique des amas globulaires stellaires. *Proc. Imp. Acad.* **6**, 133–136 (1930)
23. Ni, W.-M.: Uniqueness, Nonuniqueness and related questions of nonlinear elliptic and parabolic equations. *Proc. Symp. Pure Appl. Math.* **45**, Amer. Math. Soc., 229–241 (1986)
24. Ni, W.-M., Yotsutani, S.: Semilinear elliptic equations of Matukuma-type and related topics. *Japan J. Appl. Math.* **5**, 1–32 (1988)
25. Rein, G.: Static shells for the Vlasov-Poisson and Vlasov-Einstein systems. *Indiana Univ. Math. J.* **48**, 335–346 (1999)
26. Rein, G.: Flat steady states in stellar dynamics - Existence and stability. *Comm. Math. Phys.* **205**, 229–247 (1999)
27. Rein, G.: Stationary and static stellar dynamic models with axial symmetry. *Nonlinear analysis; Theory, Methods and Applications* **41**, 313–344 (2000)
28. Rein, G.: Static solutions of the spherically symmetric Vlasov - Einstein system. *Math. Proc. Camb. Phil. Soc.* **115**, 559–570 (1994)
29. Rein, G.: Static shells for the Vlasov-Poisson and Vlasov-Einstein systems. *Indiana University Math. J.* **48**, 335–346 (1999)
30. Rein, G.: Collisionless Kinetic Equations from Astrophysics - The Vlasov - Poisson System in: Handbook of Differential Equations, Evolutionary Equations, Vol. 3, Chapter 5, 383–476 (C.M. Dafermos and E. Feireisl ed.), Elsevier 2007.
31. Rein, G., Rendall, A.D.: Smooth static solutions of the spherically symmetric Vlasov - Einstein system. *Ann. de l'Institut H. Poincaré, Physique Théorique* **59**, 383–397 (1993)
32. Rein, G., Rendall, A.D.: Compact support of spherically symmetric equilibria in non-relativistic and relativistic galactic dynamics. *Math. Proc. Camb. Phil. Soc.* **128**, 363–380 (2000)
33. Rein, G., Rendall, A.D., Schaeffer, J.: A regularity theorem for the spherically symmetric Vlasov-Einstein system. *Comm. Math. Phys.* **168**, 467–478 (1995)
34. Thieme, H.R.: Asymptotically autonomous differential equations in the plane. *Rocky Mountain J. Math.* **24**, 351–380 (1994)
35. Yanagida, E.: Uniqueness of positive radial solutions of $\Delta u + g(r)u + h(r)u^p = 0$ in \mathbb{R}^n . *Arch. Rational Mech. Anal.* **115**, 257–274 (1991)
36. Yanagida, E.: Structure of positive radial solutions of Matukuma's equation. *Japan J. Indust. Appl. Math.* **8**, no. 1, 165–173 (1991)
37. Yanagida, E., Yotsutani, S.: Global Structure of Positive Solutions to Equations of Matukuma Type. *Arch. Rational Mech. Anal.* **134**, 199–226 (1996)
38. Yanagida, E.: Extinction and blowup of positive radial solutions for a semilinear elliptic equation. *Nonlinear Anal.* **39**, 365–377 (2000)
39. Wong, J.S.W.: On the generalized Emden-Fowler equation. *SIAM Rev.* **17**, 339–360 (1975)

Jürgen Batt
Mathematisches Institut
der Universität München
Theresienstr. 39,
D 80333 München, Germany
e-mail: batt@mathematik.uni-muenchen.de

Yi Li
Department of Mathematics
University of Iowa
Iowa City, Iowa 52242, USA
and
Department of Mathematics
Xian Jiaotong University
Xian, Shannxi, China
e-mail: yi-li@uiowa.edu

Keywords and phrases: Matukuma equation, radial solutions, semilinear elliptic equations, Emden-Fowler equation, asymptotically autonomous systems, Volterra systems, Vlasov-Poisson system, stationary stellar dynamic models